DustPedia: Multiwavelength photometry and imagery of 875 nearby galaxies in 42 ultraviolet-microwave bands

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ABSTRACT

Aims. The DustPedia project is capitalising on the legacy of the Herschel Space Observatory, using cutting-edge modelling techniques to study dust in the 875 DustPedia galaxies – representing the vast majority of extended galaxies within 3000 km s⁻¹ that were observed by Herschel. This work requires a database of multiwavelength imagery and photometry that greatly exceeds the scope (in terms of wavelength coverage and number of galaxies) of any previous local-Universal survey.

Methods. We constructed a database containing our own custom Herschel reductions, along with standardised archival observations from GALEX, SDSS, DSS, 2MASS, WISE, Spitzer, and Planck. Using these data, we performed consistent aperture-matched photometry, which we combined with external supplementary photometry from IRAS and Planck.

Results. We present our multiwavelength imagery and photometry across 42 UV-microwave bands for the 875 DustPedia galaxies. Our aperture-matched photometry, combined with the external supplementary photometry, represents a total of 21 857 photometric measurements. A typical DustPedia galaxy has multiwavelength photometry spanning 25 bands. We also present the Comprehensive & Adaptable Aperture Photometry Routine (CAAPR), the pipeline we developed to carry out our aperture-matched photometry. CAAPR is designed to produce consistent photometry for the enormous range of galaxy and observation types in our data. In particular, CAAPR is able to determine robust cross-compatible uncertainties, thanks to a novel method for reliably extrapolating the aperture noise for observations that cover a very limited amount of background. Our rich database of imagery and photometry is being made available to the community.

Key words. galaxies: photometry – galaxies: general – techniques: photometric – dust, extinction – surveys – catalogs

1. Introduction

Over the past 10–15 yr dramatic progress has been made in the study of cosmic dust as a window into the nature and evolution of galaxies. This advancement has been primarily driven by the wealth of excellent data provided by far-infrared (FIR) and submillimetre (submm) observatories such as Herschel (Pilbratt et al. 2010), Planck (Planck Collaboration I 2011), Spitzer (Werner et al. 2004), the James Clerk Maxwell Telescope (JCMT), and more recently the Atacama Large Millimetre/submillimetre Array (ALMA). Herschel in particular was especially well suited for the study of nearby galaxies; its rapid mapping abilities enabled it to observe a sizeable portion of the galaxies in the local Universe, with a combination of sensitivity, resolution, and broad wavelength coverage that remains unmatched.

The DustPedia project (Davies et al. 2017) is working towards a definitive understanding of dust in the local Universe, by capitalising on the legacy of Herschel. The DustPedia project is funded by the European Union, as a European Research Council (ERC) 7th Framework Program (FP7) project (PI Jon Davies, proposal 606824): http://dustpedia.com/

Photometry data tables are only available at the CDS via anonymous ftp to cdsarc.u-strasbg.fr (130.79.128.5) or via http://cdsarc.u-strasbg.fr/viz-bin/qcat?J/A+A/609/A37
sample consists of every galaxy that was observed by Her- schel that lies within a velocity of 3000 km s\(^{-1}\) (correspond- ing to 41 Mpc distance, assuming \(H_0 = 73.24 \text{ km s}^{-1} \text{ Mpc}^{-1}\); Riess et al. 2016), and has D25 > 1’ (D25 being the major axis isophote at which optical surface brightness falls beneath 25 mag arcsec\(^{-2}\)); these criteria were evaluated using the HyperLEDA database\(^2\),\(^3\) (Makarov et al. 2014). Additionally, the DustPedia sample only includes galaxies that have a detected stellar component; WISE observations at 3.4 \(\mu\)m are the deep- est all-sky data sensitive to the stellar component of galaxies, and hence provide the most consistent way of implementing this stellar detection requirement. Therefore, the DustPedia sample only includes galaxies brighter than the WISE 3.4 \(\mu\)m all-sky 5\(\sigma\) sensitivity limit of 19.91 \(m_{AB}\) (although this requirement has little impact upon the final sample, as <1% of candidate HyperLEDA galaxies fail to meet the brightness limit). For full sample details see Davies et al. (2017). Note that the DustPedia sample excludes the Local Group galaxies of Andromeda, Triangulum, and the Magellanic Clouds, as working with such exceptionally-extended systems would entail fundamentally different data pro- cessing and analysis.

DustPedia is combining cutting-edge methods for studying dust: physically-motivated dust modelling with The Heterogeneous Evolution Model for Interstellar SOLIDS (THEMIS; Jones et al. 2016; Jones et al. 2017; Ysard et al. 2016); hierarchichal Bayesian spectral energy distribution (SED) fitting with Hierarchical Bayesian Inference for dust Emission (HerBIE; Galliano et al., in prep.); and 3_dimensional radiative transfer modelling and fitting with Stellar Kinematics Including Radiative ‘Transfer (SKIRT; Baes et al. 2011; Camps & Baes 2015) and FitSKIRT (De Geyer et al. 2013).

The study of nearby galaxies now involves increasingly-extensive multiwavelength datasets, and DustPedia is no excep- tion to this. The tools used in modern extragalactic astronomy require that data from across broad swaths of the spectrum can be directly compared, despite the very different natures of the data employed. This is especially true for the study of dust, where radiative transfer and SED modelling tools – such as HerBIE, SKIRT, MAGPHYS (da Cunha et al. 2008), MOCASSIN (Ercolano et al. 2003, 2005), CIGALE (Burgarella et al. 2005; Noll et al. 2009), and GRASIL (Silva et al. 1998) – critically re- quire that the flux densities being used represent self-consistent measurements. For example, in energy-balance SED-fitting, it is vital that the stellar population sampled by ultraviolet (UV), optical, and near-infrared (NIR) data points is the same stellar component; WISE observations at 3.4 \(\mu\)m all-sky sensitivity limit of 19.91 \(m_{AB}\) (although this requirement has little impact upon the final sample, as <1% of candidate HyperLEDA galaxies fail to meet the brightness limit). For full sample details see Davies et al. (2017). Note that the DustPedia sample excludes the Local Group galaxies of Andromeda, Triangulum, and the Magellanic Clouds, as working with such exceptionally-extended systems would entail fundamentally different data pro- cessing and analysis.

Moreover, in the realm of multiwavelength photometry, it is often overlooked that measuring flux densities in a consistent and accurate manner is invariably much easier than measuring the uncertainties on those flux densities in a similarly consistent and accurate manner. However, with the growing prevalence of Bayesian techniques in astronomy (see review in Loredo 2013, and references therein), it has never been more important that photometric uncertainties be robust and directly-comparable.

Therefore, in order to best exploit DustPedia’s unique com- bination of advanced tools for the study of dust in galaxies, the cornerstone of the project is the DustPedia database\(^4\) – which contains standardised multiwavelength imagery of all 875 galaxies in our sample\(^5\), including both custom reductions and archival data, spanning over five orders of magnitude in wave- length from UV to microwave. Furthermore, the DustPedia database contains the results of consistent multiwavelength aperture-matched photometry, conducted using the Compre- hensive & Adaptable Aperture Photometry Routine (CAAPR) – along with supplementary photometry for additional instru- ments for which aperture-matched photometry was impractical. In total, we provide data for 42 UV-microwave bands. These data are already being employed in various accepted and in- preparation DustPedia works, including Casasola et al. (2017), Liou et al. (in prep.), Cassara et al. (in prep.), Moenko et al. (in prep.), Evans et al. (in prep.), Verstroken et al. (in prep.), and Nersesian et al. (in prep.). Ultimately, the DustPedia database will also contain the results of the SED fitting and radia- tive transfer modelling performed on these data, for each galaxy.

In this paper, we present the imagery and photometry that form the centrepiece of the DustPedia database. The multiwave- length imagery – both our custom Herschel reductions, and the standardised UV-microwave ancillary data – are presented in Sect. 2. The functionality of the CAAPR pipeline, which we developed to produce our consistent multiwavelength aperture-matched photometry, is described in Sect. 3. The flagging we carry out for our photometry is explained in Sect. 4. Supple- mentary photometry from additional observatories is laid out in Sect. 5. In Sect. 6 we detail the various tests we carried out to validate the quality of our photometry. The contents and format of the DustPedia data products are specified in Sect. 7, including information on distance measures for each galaxy.

For the sake of brevity and readability, we refer to “flux densities” as “fluxes” throughout the rest of this work. We adopt a Hubble constant of \(H_0 = 73.24 \text{ km s}^{-1} \text{ Mpc}^{-1}\) (Riess et al. 2016).

2. Multiwavelength imagery

For the DustPedia galaxies, we produced dedicated reductions of the Herschel-SPIRE and Herschel-PACS maps, along with standardised preparations of archive imagery for the ancillary data.

All maps were produced in Flexible Image Transport System (FITS; Wells et al. 1981; Hanisch et al. 2001) format, with standardised file names and headers. The file name indicates the target galaxy, facility, and band in question, taking the form: [galaxy]_[facility]_[band]_fits; in cases where an error map was also available, it was stored in a separate FITS file, with a file name of the form: [galaxy]_[facility]_[band]_Error.fits. As an example, the Herschel-SPIRE 250 µm map of DustPedia galaxy NGC0891 has the file name: NGC0891_SPIRE250.fits. All galaxies are referred to by the name under which they are listed in the HyperLEDA database.

The FITS headers contain the standard World Coordinate System (WCS; Greisen & Calabretta 2002) fields, along with fields that provide a range of additional information: target galaxy, telescope, filter effective wavelength, filter name, instru- ment, map units, and the name of the database the original data was acquired from (if this is not the same for all observations in a given band). If there is an error map available, then the header includes the WCS World Coordinate System (WCS; Greisen & Calabretta 2002) fields, along with fields that provide a range of additional information: target galaxy, telescope, filter effective wavelength, filter name, instru- ment, map units, and the name of the database the original data was acquired from (if this is not the same for all observations in a given band). If there is an error map available, then the header

\(^2\) http://leda.univ-lyon1.fr/

\(^3\) The HyperLEDA database is continually updated to reflect newly-incorporated data; therefore note that the DustPedia sample is derived from HyperLEDA queries performed in January 2015.

\(^4\) http://dustpedia.astro.noa.gr

\(^5\) Subsequent to the publication of Davies et al. (2017), DustPedia source SDDS164054.5+535383.8 was found to be a shredded portion of NGC 5474; so whereas Davies et al. (2017) refer to the DustPe- dia sample containing 876 galaxies, we here only deal with the revised sample of 875.
2.1. Herschel imagery

Full details of our Herschel data reduction process are provided in Davies et al. (2017); here we provide only a brief summary, for the sake of completeness. For both Spectral and Photometric Imaging REceiver (SPIRE; Griffin et al. 2010) and Photodetector Array Camera and Spectrometer (PACS; Poglitsch et al. 2010) data, the Herschel Science Archive (HSA)\(^6\) was queried to find all Herschel photometer observations that provided coverage of each DustPedia galaxy.

2.1.1. Herschel-SPIRE

In the first stage of our SPIRE reduction, raw Level-0 data (acquired from the HSA) was processed up to Level-1 (calibrated pointed timelines) using HIPE v13\(^7\). The Level-1 data was processed up to Level-2 (final maps) using the Bright Galaxy Adaptive Element (BRIGADE; Smith 2012) pipeline. Observations that overlapped were combined into single contiguous maps\(^8\). The final timelines were refined with the HIPE v13 desstriper; the improved processing provided by BrIGADE, and the fact that we combine all overlapping observations, allows the desstriper to operate more effectively than it would otherwise. Final SPIRE maps were produced using the HIPE v13 naïve mappemaker, with pixel sizes of 6, 8, and 12 arcseconds at 250, 350, and 500 \(\mu\text{m}\) respectively (\(\approx 1/3\) each band’s FWHM). The photometric calibration of our maps used the beam area values provided in HIPE v13; specifically, 469.7, 831.7, and 1793.5 arcsec\(^2\) at 250, 350, and 500 \(\mu\text{m}\). If a target galaxy was located in a reduced map larger than one degree in RA or Dec, then a cutout of 1° × 1°, centred on the target, was produced as the final map; if the reduced map was smaller than one degree, then it was used in its entirety as the final map. In total, 844 (96\%) of the DustPedia galaxies have SPIRE coverage. SPIRE maps are calibrated assuming a \(v^{-1}\) reference spectrum.

2.1.2. Herschel-PACS

For our PACS reduction, Level-0 raw data was processed up to Level-1 timelines using HIPE v13. The Level-1 data was then processed into Level-2 maps using the standalone SCANMORPHOUS v24 pipeline (Roussel 2012, 2013). The SCANMORPHOUS processing took account of whether the PACS data in question was taken in scan map, mini map, or parallel mode. Observations taken in different observing modes were not combined; when data from more than one mode was available, the superior set of observations was identified, based upon depth, coverage sky area, and scan speed. In the case of some target galaxies, the PACS data was found to be unusable. The final PACS maps have pixel sizes of 2, 3, and 4 arcseconds at 70, 100, and 160 \(\mu\text{m}\) respectively (\(\approx 1/3\) the FWHM, in keeping with the SPIRE maps). For data from very large PACS fields, the sheer quantity of data makes combining all available data computationally prohibitive – instead, we only make use of the deep data taken by HeViCS (the Herschel Virgo Cluster Survey; Davies et al. 2010).

\(\text{Fig. 1. Illustration of the spectral coverage provided by the DustPedia database, showing filter response functions of all bands for which we present data. As can be seen, the data we employ effectively provides complete sampling of over five orders of magnitude in wavelength. Response functions of the bands for which we present both imagery and aperture-matched photometry are traced with solid lines. Bands for which we present supplementary external photometry (see Sect. 5) are traced with dashed lines. Bands for which we present imagery only (see Sect. 2.2.3) are traced with dotted lines.}\)
Table 1. Details of each band for which we have data.

<table>
<thead>
<tr>
<th>Facility</th>
<th>Effective wavelength</th>
<th>Band name</th>
<th>Imagery present</th>
<th>Photometry present</th>
<th>Pixel width</th>
<th>Resolution FWHM</th>
<th>Calibration uncertainty</th>
</tr>
</thead>
<tbody>
<tr>
<td>GALEX</td>
<td>153 nm</td>
<td>FUV</td>
<td>797</td>
<td>794</td>
<td>3.2</td>
<td>4.3</td>
<td>4.5 (°)</td>
</tr>
<tr>
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<td>NUV</td>
<td>832</td>
<td>830</td>
<td>3.2</td>
<td>5.3</td>
<td>2.7 (°)</td>
</tr>
<tr>
<td>SDSS</td>
<td>353 nm</td>
<td>a</td>
<td>655</td>
<td>655</td>
<td>0.45</td>
<td>1.3</td>
<td>1.3 (°)</td>
</tr>
<tr>
<td>SDSS</td>
<td>475 nm</td>
<td>g</td>
<td>655</td>
<td>655</td>
<td>0.45</td>
<td>1.3</td>
<td>0.8 (°)</td>
</tr>
<tr>
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<td>622 nm</td>
<td>r</td>
<td>655</td>
<td>655</td>
<td>0.45</td>
<td>1.3</td>
<td>0.8 (°)</td>
</tr>
<tr>
<td>SDSS</td>
<td>763 nm</td>
<td>i</td>
<td>655</td>
<td>655</td>
<td>0.45</td>
<td>1.3</td>
<td>0.7 (°)</td>
</tr>
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<td>z</td>
<td>655</td>
<td>655</td>
<td>0.45</td>
<td>1.3</td>
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<td>DSS1</td>
<td>450 nm</td>
<td>B</td>
<td>794</td>
<td>–</td>
<td>1–1.7</td>
<td>1.9 (1.5–3.0)</td>
<td>–</td>
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<tr>
<td>DSS1</td>
<td>660 nm</td>
<td>R</td>
<td>794</td>
<td>–</td>
<td>1–1.7</td>
<td>1.9 (1.5–3.0)</td>
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</tr>
<tr>
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<td>B</td>
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<td>–</td>
<td>1–1.7</td>
<td>1.9 (1.5–3.0)</td>
<td>–</td>
</tr>
<tr>
<td>DSS2</td>
<td>660 nm</td>
<td>R</td>
<td>861</td>
<td>–</td>
<td>1–1.7</td>
<td>1.9 (1.5–3.0)</td>
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<td>J</td>
<td>875</td>
<td>875</td>
<td>1</td>
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<td>H</td>
<td>875</td>
<td>875</td>
<td>1</td>
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<td>1.9 (°)</td>
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<td>875</td>
<td>1</td>
<td>2.0</td>
<td>1.9 (°)</td>
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<td>875</td>
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<tr>
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<td>6.4</td>
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<td>WISE</td>
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<td>875</td>
<td>875</td>
<td>1.375</td>
<td>6.5</td>
<td>4.6 (°)</td>
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<tr>
<td>WISE</td>
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<td>875</td>
<td>1.375</td>
<td>12</td>
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<tr>
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<td></td>
<td>644</td>
<td>623</td>
<td>0.75</td>
<td>1.66</td>
<td>3 (°)</td>
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<tr>
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<td>777</td>
<td>777</td>
<td>0.75</td>
<td>1.72</td>
<td>3 (°)</td>
</tr>
<tr>
<td>Spitzer</td>
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<td></td>
<td>392</td>
<td>374</td>
<td>0.6</td>
<td>1.88</td>
<td>3 (°)</td>
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<tr>
<td>Spitzer</td>
<td>8.0 μm (IRAC-4)</td>
<td></td>
<td>411</td>
<td>392</td>
<td>0.6</td>
<td>1.98</td>
<td>3 (°)</td>
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<tr>
<td>Spitzer</td>
<td>24 μm (MIPS-1)</td>
<td></td>
<td>491</td>
<td>477</td>
<td>2.4–2.6</td>
<td>6</td>
<td>5 (°)</td>
</tr>
<tr>
<td>Spitzer</td>
<td>70 μm (MIPS-2)</td>
<td></td>
<td>198</td>
<td>177</td>
<td>4</td>
<td>18</td>
<td>10 (°)</td>
</tr>
<tr>
<td>Spitzer</td>
<td>160 μm (MIPS-3)</td>
<td></td>
<td>184</td>
<td>171</td>
<td>8</td>
<td>38</td>
<td>12 (°)</td>
</tr>
<tr>
<td>PACS</td>
<td>70 μm –</td>
<td></td>
<td>255</td>
<td>244</td>
<td>2</td>
<td>9 (5.8–12.2)</td>
<td>7 (°)</td>
</tr>
<tr>
<td>PACS</td>
<td>100 μm –</td>
<td></td>
<td>716</td>
<td>701</td>
<td>3</td>
<td>10 (6.9–12.7)</td>
<td>7 (°)</td>
</tr>
<tr>
<td>PACS</td>
<td>160 μm –</td>
<td></td>
<td>771</td>
<td>753</td>
<td>4</td>
<td>13 (12.1–15.7)</td>
<td>7 (°)</td>
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<tr>
<td>SPIRE</td>
<td>250 μm (PSW)</td>
<td></td>
<td>844</td>
<td>844</td>
<td>6</td>
<td>18</td>
<td>5.5 (°)</td>
</tr>
<tr>
<td>SPIRE</td>
<td>350 μm (PMW)</td>
<td></td>
<td>844</td>
<td>844</td>
<td>8</td>
<td>25</td>
<td>5.5 (°)</td>
</tr>
<tr>
<td>SPIRE</td>
<td>500 μm (PLW)</td>
<td></td>
<td>844</td>
<td>844</td>
<td>12</td>
<td>36</td>
<td>5.5 (°)</td>
</tr>
<tr>
<td>Planck</td>
<td>350 μm (857 GHz)</td>
<td></td>
<td>875</td>
<td>394</td>
<td>102.1</td>
<td>278</td>
<td>6.4 (°)</td>
</tr>
<tr>
<td>Planck</td>
<td>550 μm (545 GHz)</td>
<td></td>
<td>875</td>
<td>279</td>
<td>102.1</td>
<td>290</td>
<td>6.1 (°)</td>
</tr>
<tr>
<td>Planck</td>
<td>850 μm (353 GHz)</td>
<td></td>
<td>875</td>
<td>197</td>
<td>102.1</td>
<td>296</td>
<td>0.78 (°)</td>
</tr>
<tr>
<td>Planck</td>
<td>1.38 mm (217 GHz)</td>
<td></td>
<td>875</td>
<td>97</td>
<td>102.1</td>
<td>301</td>
<td>0.16 (°)</td>
</tr>
<tr>
<td>Planck</td>
<td>2.10 mm (143 GHz)</td>
<td></td>
<td>875</td>
<td>29</td>
<td>102.1</td>
<td>438</td>
<td>0.07 (°)</td>
</tr>
<tr>
<td>Planck</td>
<td>3.00 mm (100 GHz)</td>
<td></td>
<td>875</td>
<td>19</td>
<td>205.7</td>
<td>581</td>
<td>0.09 (°)</td>
</tr>
<tr>
<td>Planck</td>
<td>4.26 mm (70 GHz)</td>
<td></td>
<td>875</td>
<td>11</td>
<td>205.7</td>
<td>799</td>
<td>0.20 (°)</td>
</tr>
<tr>
<td>Planck</td>
<td>6.81 mm (44 GHz)</td>
<td></td>
<td>875</td>
<td>18</td>
<td>205.7</td>
<td>1630</td>
<td>0.26 (°)</td>
</tr>
<tr>
<td>Planck</td>
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<td></td>
<td>875</td>
<td>35</td>
<td>205.7</td>
<td>1940</td>
<td>0.35 (°)</td>
</tr>
<tr>
<td>IRAS</td>
<td>12 μm –</td>
<td></td>
<td>–</td>
<td>598</td>
<td>–</td>
<td>270</td>
<td>20 (°)</td>
</tr>
<tr>
<td>IRAS</td>
<td>25 μm –</td>
<td></td>
<td>–</td>
<td>578</td>
<td>–</td>
<td>276</td>
<td>20 (°)</td>
</tr>
<tr>
<td>IRAS</td>
<td>60 μm –</td>
<td></td>
<td>–</td>
<td>675</td>
<td>–</td>
<td>282</td>
<td>20 (°)</td>
</tr>
<tr>
<td>IRAS</td>
<td>100 μm –</td>
<td></td>
<td>–</td>
<td>682</td>
<td>–</td>
<td>300</td>
<td>20 (°)</td>
</tr>
</tbody>
</table>

Notes. For FUV-Ks bands, we refer to each band by its listed “Band name”; otherwise we refer to bands by wavelength. The “Imagery present” and “Photometry present” columns give the number of galaxies in each band for which we present imagery and photometry (there are some targets for which we present imagery, but where photometry was not possible; see Sect. 3.6). For supplementary photometry, we give the number of galaxies in each band for which photometry was available from IRAS SCANIPI and Planck CCS2 (see Sect. 5). References for calibration uncertainties are provided below. For bands with significantly varying resolutions, a typical value is given, followed by the full range in parentheses.

References. (a) Morrissey et al. (2007); (b) SDSS DR12 Data Release Supplement: https://www.sdss3.org/dr12/scope.php; (c) Cohen et al. (2003); (d) WISE All-Sky Release Explanatory Supplement: http://wise2.ipac.caltech.edu/docs/release/allsky/expsup/sec4_4h.html; (e) IRAC Instrument Handbook https://irsa.ipac.caltech.edu/data/SPITZER/docs/ircalinstrumenhandbook/17/7_Toc410728385; (f) MIPS Instrument Handbook: https://irsa.ipac.caltech.edu/data/SPITZER/docs/mips/mipsinstrumenhandbook/42/7_Toc280032317; (g) PACS Instrument & Calibration Wiki: http://herschel.esac.esa.int/twiki/bin/view/Public/PacsCalibrationWeb; (h) SPIRE Instrument & Calibration Wiki: http://herschel.esac.esa.int/twiki/bin/view/Public/SpireCalibrationWeb; (i) Planck Collaboration VII (2016); (j) Planck Collaboration V (2016); (k) Sanders et al. (2003), Miville-Deschênes et al. (2005).

the final maps have a diameter of 0.35°, centred upon the target galaxies; this is smaller than for the SPIRE maps, due to the additional computational complexity of reducing PACS data. In total, 771 (88%) of the DustPedia galaxies have usable PACS coverage. PACS maps are calibrated assuming a ν−1 reference spectrum.
2.2. Ancillary imagery

Our ancillary imaging data consists of observations from 8 facilities; the GALaxy eXplorer (GALEX; Morrissey et al. 2007), the Sloan Digital Sky Survey (SDSS; York et al. 2000; Eisenstein et al. 2011), the Digitized Sky Survey (DSS), the 2 Micron All-Sky Survey (2MASS; Skrutskie et al. 2006), the Wide-field Infrared Survey Explorer (WISE; Wright et al. 2010), the Spitzer Space Telescope, (Werner et al. 2004), and Planck (Planck Collaboration I 2011). The basic parameters of each band are listed in Table 1. The following sections will detail the specifics of how the data for each facility was processed; however the same general procedure was employed throughout.

The various facilities from which we take our ancillary data were chosen on the basis of observation quality, data availability, and wavelength coverage. All of the facilities we employ provide coverage for the vast majority of the DustPedia galaxies (>75%). Whilst data with improved resolution and/or sensitivity is available in certain wavelength ranges, coverage is only ever available for a small minority of our targets. The only potentially “borderline” case for inclusion was the UKIRT Infrared Deep Sky Survey (UKIDSS; Lawrence et al. 2007). Whilst UKIDSS data is higher-resolution and deeper than 2MASS, it still only provides coverage for a minority (40%) of the DustPedia galaxies, and questions remain regarding possible photometric inconsistencies when compared to the Visible and Infrared Survey Telescope for Astronomy (VISTA), 2MASS, and SDSS surveys (see Driver et al. 2016). As such, we opted not to include UKIDSS data at present; alternative facilities operating in the UKIDSS wavelength range (such as VISTA) only provide coverage for <15% of our target galaxies. Nonetheless, we do not rule out incorporating UKIDSS, or other datasets, in to the DustPedia database at some point in the future.

In general, the ancillary imagery takes the form of 0.5° × 0.5° cutouts (oriented East-North), centred on their respective target galaxies. However, for galaxies sufficiently extended that they would span more than 20% of the diameter of such a cutout (i.e., those with D25 > 6”), larger cutouts of 1° × 1° were used (although less than 10% of the DustPedia galaxies are this extended); this ensures that all cutouts can contain enough sky to allow for proper background measurement, aperture noise estimation, etc., even for extremely extended sources.

For a particular target galaxy, observed in a given band, by a given facility, all available imagery covering the region of interest was identified. The data was retrieved from the official online archive of the facility in question, and then mosaiced to produce the cutout map.

The tool employed to carry out the mosaicing was Montage⁹. Montage was used to re-grid the images to a common pixel projection, match their background levels, and co-add them; when the observations being co-added were of different depths, the contribution of each image to the final co-addition was weighted appropriately (by considering exposure time information, or using error maps, etc.).

Mosaicing via Montage has been found by previous authors to give rise to no consequential degradation of image quality (Blanton et al. 2011). However, when reprojecting images to a different coordinate grid rotation (via either Montage or other tools), moiré patterns can appear, particularly if the noise budget is dominated by Poisson statistics. As well as being a conspicuous artefact, moiré patterns also have an effect upon map noise characteristics at smaller scales. An illustration of a reprojection-induced moiré pattern is shown in Fig. 2. If necessary, it is possible to prevent the appearance of moiré patterns by slightly (~10%) enlarging the pixel size of the reprojected maps. This was necessary for our GALEX and SDSS data (the specifics are discussed in their respective sections below); all photometry and other analysis was performed on maps that had undergone the pixel enlargement, hence preventing moiré patterns from affecting our results.

Fig. 2. Illustration of the moiré patterns which can arise when reprojecting an image to a different orientation, using the example of an SDSS r-band image, shown before and after lower reprojection. The RMS pixel noise at the “peaks” of the pattern is ~50% greater than in the “troughs”.

Note that the background-matching carried out by Montage is not a background subtraction. Rather, Montage adjusts the level of each map, so that it matches, as closely as possible, those maps it overlaps with. This background matching does not scale the maps to any absolute physical zero-point – it merely minimises the difference between overlapping maps. This keeps the scaling of the map units as close as possible to that of the input data (whilst still providing good mosaicing), permitting users to convert the maps back to their native data units (for instance, in order to consider pixel Poisson statistics for resolved analyses) as per Table A.1. As such, local background subtraction (e.g., with a background annulus) must still be carried out when performing photometry on the resulting maps. Moreover, users will find local background subtraction unavoidable when working with much of the data, given the necessity of accounting for sky brightness, Galactic cirrus, etc.

2.2.1. GALEX

UV observations made by the GALaxy eXplorer (GALEX; Morrissey et al. 2007) were acquired from the GR6/7

⁹ http://montage.ipac.caltech.edu/
data release (Bianchi 2014), hosted at the Mikulski Archive for Space Telescopes (MAST\(^\text{10}\)).

GALEX coverage is available for the vast majority of the DustPeda \textit{Herschel} galaxies; 797 (91\%) in FUV, and 832 (95\%) in NUV. Coverage is more extensive in the NUV due to the fact that GALEX’s FUV camera failed in June 2009, whilst the NUV camera continued operating until the satellite was decommissioned.

The GALEX \textsc{view} utility provided by MAST was used to identify all GALEX tiles in the vicinity of the DustPeda \textit{Herschel} galaxies. These tiles were retrieved and co-added in line with the process detailed in Sect. 2.2, but with some additional considerations specific to handling GALEX data, described here.

The final GALEX cutouts have pixel sizes of 3.2'' in both FUV and NUV, in contrast to the standard GALEX pixel sizes of 1.5''; there are two reasons for this. In shallow GALEX tiles, particularly in the FUV, an extremely large fraction of the pixels have a value of zero – pixels where no photons were detected during the exposure. Indeed, for the extremely shallow GALEX “all-sky” survey, over 99\% of the pixels in a FUV tile can be expected to be zero-value, and most of the pixels that do contain flux have a pixel value that represents only a single photon. This presents an issue when co-adding GALEX tiles. In order to perform the co-addition, it is necessary to re-sample all of the input images to a common pixel grid. But when re-sampling shallow GALEX tiles, this has the effect of spreading the flux in single-photon pixels over several reprojected pixels. Not only does this manifest as a very obvious aliasing artefact when inspected visually, but spreading the flux from a single photon over an extended area of sky is also physically dubious. Furthermore, in reprojected GALEX tiles with smaller numbers of zero-value pixels, noticeable moiré patterns often appear, introducing poorly-behaved noise on some scales. We avoided both of these artefacts by re-projecting the tiles to larger pixel sizes when performing the co-addition. Whilst this means that our GALEX maps are no longer Nyquist sampled, we will primarily be using them in concert with much lower-resolution data (WISE, Spitzer, \textit{Herschel}, etc.), rendering this of less concern.

Care needs to be taken to correctly handle zero-value pixels in GALEX tiles. A zero-value pixel can either represent that a pixel lies outside the coverage region for the observation, or it can represent a pixel that lies inside the coverage region but where no photons were detected during the exposure. We used the GALEX response maps corresponding to each tile to determine the exact region observed; only the zero-value pixels within the coverage area were included in mosaics.

Coverage in the outer portions of GALEX tiles tend to be of very low quality, due to aperture vignetting (Morrissey et al. 2007). To prevent this leading to degradation of the final mosaics, we follow WiggleZ (Drinkwater et al. 2010) and GAMA (Driver et al. 2016) in masking all pixels outside the central 35'' of each tile, excluding the outer -9\% of the coverage area.

When performing the co-addition of the GALEX images, each tile’s contribution was appropriately weighted according to the exposure-time information provided with the tile. The large numbers of pixels in shallow GALEX tiles that contain flux from few or zero photons also causes problems when attempting to match the background levels of tiles being co-added. The distribution of pixel values in shallow GALEX tiles is highly discrete and asymmetrical, especially at the distribution peak; this makes it difficult to ascertain the actual background level – and hence difficult to match to the background levels of other tiles when co-adding. To address this, a copy of each tile was convolved with a 30'' (20 pixel) diameter top hat kernel; this had the effect of smoothing out discrete pixel value levels. The pixel values of each smoothed map were iteratively sigma-clipped to a three-sigma threshold in order to remove particularly bright pixels; a flat plane was fit in a least-squares manner to the remaining pixels. The level of this plane was taken to represent the background level of the tile in question, and subtracted prior to co-addition.

The native GALEX tile pixel units of counts s\(^{-1}\) pix\(^{-1}\) were translated to Jy pix\(^{-1}\) using conversion factors of 1.076 \times 10\(^{-4}\) Jy counts\(^{-1}\) s in the FUV and 3.373 \times 10\(^{-4}\) Jy counts\(^{-1}\) s in the NUV; these correspond to the standard GALEX AB magnitude zero points of 18.82 mag and 20.08 mag in FUV and NUV respectively (Morrissey et al. 2007).

### 2.2.2. SDSS

The Sloan Digital Sky Survey (SDSS; York et al. 2000; Eisenstein et al. 2011) provides UV, optical, and NIR imaging of 35\% of the sky in the \textit{ugriz} bands. SDSS data was acquired from SDSS DR12 (Alam et al. 2015), hosted at the SDSS Science Archive Server\(^\text{11}\), yielding coverage for 656 (75\%) of the DustPeda galaxies.

The \textsc{archiveList} tool, part of Montage, was used to identify all SDSS fields in the vicinity of the DustPeda galaxies. For each band, most places in the SDSS footprint have been observed more than once; however, for each point, a particular field – the one deemed of the best quality – is designated to be primary. Primary frames are those used by SDSS for photometry, mosaic generation, etc. Only fields designated as being primary by the SDSS were selected to be retrieved and co-added. By only using primary fields we ensure that only science-quality SDSS data is used.

Once the primary SDSS fields had been retrieved, they were co-added in line with the process described in Sect. 2.2, with the exception that the final cutouts have pixel sizes of 0.45''. This was done to minimise the appearance of the conspicuous moiré patterns that otherwise occurred when reprojecting to a North-East orientation. Given the median seeing full-width half-maximum (FWHM) of 1.43'' for SDSS imaging fields\(^\text{12}\), a pixel size of 0.45'' is within the limit required for Nyquist sampling. Indeed, given that Nyquist sampling is achieved in diffraction-limited observations provided that pixel sizes are smaller than FWHM/2.44, the 0.45'' pixels are suitable for frames with seeing-limited resolution >1.09'', which is the case for >97\% of SDSS frames\(^\text{13}\).

Maps retrieved from the SDSS DR12 Science Archive Server have native units of “nanomagnitudes” (an SDSS convenience unit of linear flux\(^\text{14}\)), scaled such that a flux in nanomagnitudes can be converted to an SDSS magnitude using a zero point magnitude of 23.5 mag. SDSS magnitudes were designed to be AB magnitudes, but in practice have been found to exhibit offsets\(^\text{15}\) of \(-0.04\) mag offset in \(g\)-band (such that \(u_{AB} = u_{\text{SDSS}} - 0.04\)), and

\(^{10}\) http://galex.stsci.edu/GR6/

\(^{11}\) http://dr12.sdss3.org/

\(^{12}\) http://www.sdss.org/dr12/imaging/other_info/


\(^{14}\) http://www.sdss.org/dr12/algorithm/magnitudes/

\(^{15}\) http://www.sdss.org/dr12/algorithm/fluxcal/

#SDSSStoAB
0.02 mag in z-band (such that $z_{AB} = z_{SDSS} + 0.02$). We corrected the pixel units for these offsets, scaling them to the AB magnitude scale, and thereby allowing conversion to Jy pix$^{-1}$.

### 2.2.3. DSS

The Digitized Sky Survey (DSS) provides UV, optical, and NIR coverage of the entire sky, in the form of scans of photographic survey plates from the Samuel Oschin telescope at Palomar Observatory, and the UK Schmidt Telescope at the Anglo-Australian Observatory.

As the DSS is produced from photographic plates, the pixel values are not in units of linear flux; rather, the units are “photographic densities”. The photographic density scale is not constant across the DSS, and varies from plate to plate. Whilst there are some general prescriptions for photometric calibration of the DSS, these only apply to individual contributing sub-surveys, and typically have calibrations no better than 0.5 mag (Doggett et al. 1996). Whilst a photometric calibration can be determined for each individual plate by comparison to standard reference stars, great care needs to be taken when doing so, particularly with regards to saturation of reference stars, and whether or not any portion of each target galaxy is also saturated. As such, we opt not to attempt a calibration of the DSS imagery, instead leaving it in its native units of photographic densities. However, even without photometric calibration, the DSS data is valuable, as it provides consistent optical coverage of our entire sample – including the 25% of DustPedia not covered by the SDSS. For example, the DSS data can still be used to work out the optical axial ratio, position angle, and extent (not in an absolute isophotal sense, but using signal-to-noise analysis) of a galaxy; all useful information when determining apertures for our multiwavelength photometry (see Sect. 3.4).

DSS B-band and R-band imagery for the DustPedia Herschel galaxies was retrieved from the NASA SkyView service\(^\text{16}\), and used to produce standard cutouts as described in Sect. 2.2. However, given the complications detailed above, we made no attempt to convert the pixel values to units of Jy pix$^{-1}$\(^\text{17}\); instead, they were left as photographic densities. Whilst the DSS provides UV and NIR imagery for some of the sky, we did not make use of these data, as GALEX, SDSS and 2MASS provide vastly superior coverage. The DSS is divided into two phases, DSS1 and DSS2, each made up of different sets of contributing sub-surveys. As the different contributing sub-surveys operate at a range of resolutions, no attempt was made to co-add the DSS1 and DSS2 imagery; rather, cutouts for each were produced separately. As many, source lists have both DSS1 and DSS2 cutouts available. The DSS cutouts have pixel sizes of either 1″ or 1.7″, depending upon which contributing sub-survey the relevant photographic plate scans came from.

### 2.2.4. 2MASS

The 2 Micron All-Sky Survey (2MASS; Skrutskie et al. 2006) provides NIR imaging of the entire sky in J-band, H-band, and $K_s$-band. 2MASS tiles in the vicinity of the DustPedia Herschel galaxies were identified and acquired from the NASA/IPAC Infrared Science Archive (IRSA)\(^\text{17}\). The retrieved 2MASS tiles were then co-added in line with the process described in Sect. 2.2. The final cutouts retain the standard 2MASS All-Sky Data Release pixel size of 1″; the pixel units were rendered in Jy pix$^{-1}$ according to the zero points provided in the headers of the input archival maps.

Note that severe sky brightness can give rise to significant aperture noise in the 2MASS bands – especially in $H$-band, primarily due to OH emission in the atmosphere (Jarrett et al. 2003). The sky brightness is typically seen to vary across angular scales of a few arcminutes, making it particularly troublesome for photometry of nearby galaxies. Because 2MASS observed in JHKs concurrently, an almost identical sky brightness structure will be seen in all three bands, varying only in magnitude. This means that the resulting aperture noise will be strongly correlated between bands. If a given $H$-band flux is erroneously high, the $J$- and $K_s$-band fluxes are likely to be too high also – although usually not to the same degree as in $H$-band, given the much worse sky brightness $H$-band tends to suffer.

### 2.2.5. WISE

The Wide-field Infrared Survey Explorer (WISE; Wright et al. 2010) provides NIR and MIR coverage of the entire sky, observing at 3.4 µm, 4.6 µm, 12 µm, and 22 µm. WISE imagery was obtained from the AllWISE data release Image Atlas, which combines WISE cryogenic and NEOWISE (Near Earth Object WISE; Mainzer et al. 2011) survey phases. AllWISE Atlas images in the vicinity of the DustPedia Herschel galaxies were identified and retrieved from IRSA, and co-added in line with the process described in Sect. 2.2 (however, given the large area covered by each AllWISE Atlas image, co-addition was very rarely necessary).

The final WISE cutouts retain the standard AllWISE Image Atlas pixel size of 1.375″. The map pixel units were rendered in Jy pix$^{-1}$ by using the zero-point magnitudes provided in the headers of the AllWISE Atlas images to convert pixel units to Vega magnitudes, which were then converted to AB magnitudes, and thence to Jy, according to the conversions given in Table 1, Sect. IV.3.a of the Explanatory Supplement to the AllWISE Data Release Products\(^\text{18}\). Extended-source corrections need to be applied when performing aperture photometry upon WISE maps, as per Table 5, Sect. IV.4.c.vii of the WISE All-Sky Data Release Explanatory Supplement\(^\text{19}\), to account for the fact that the calibration of WISE data is based upon profile fitting of point sources.

Noticeable image artefacts are occasionally encountered in the final WISE cutouts. However, as these are present in the input AllWISE Atlas images, and as there is typically only a single AllWISE Atlas image covering a given location on the sky, there is little that can currently be done to address them. The artefacts commonly encountered in WISE data are discussed in Sects. IV.4.g and VI.2.c of the WISE All-Sky Data Release Explanatory Supplement.

It should also be noted that Wright et al. (2010) and Brown et al. (2014) have found evidence that the in-orbit spectral response functions of the WISE W3 (12 µm) and W4 (22 µm) filters deviate from the laboratory-measured pre-launch functions. For monochromatic fluxes, Brown et al. (2014) report that W4 filter’s effective wavelength should be revised from 22.1 µm to 22.8 µm. However, the change for any given source depends on its spectral index in wavelength range of the W3

\(^{16}\) [http://skyview.gsfc.nasa.gov/current/cgi/query.pl](http://skyview.gsfc.nasa.gov/current/cgi/query.pl)

\(^{17}\) [http://irsa.ipac.caltech.edu/frontpage/](http://irsa.ipac.caltech.edu/frontpage/)

\(^{18}\) [http://wise2.ipac.caltech.edu/docs/release/allwise/expsum/](http://wise2.ipac.caltech.edu/docs/release/allwise/expsum/)

\(^{19}\) [http://wise2.ipac.caltech.edu/docs/release/allsky/expsum/](http://wise2.ipac.caltech.edu/docs/release/allsky/expsum/)
and W4 filters. Ultimately, this means W3 and W4 bands require more severe colour-correction than would otherwise have been necessary. As for other instruments, we perform no WISE colour-correction ourselves, as the full SED modelling necessary to determine the appropriate corrections is beyond the scope of this work. However, users of our WISE data products are reminded to be aware of the particular considerations necessary for the WISE 12 and 22 µm bands, and to ensure that they are correctly accounting for the filters’ spectral response functions. WISE maps are calibrated assuming a ν−2 reference spectrum (Wright et al. 2010).

2.2.6. Spitzer

The Spitzer Space Telescope (Werner et al. 2004) provides imaging across 7 photometric bands in the NIR to FIR, observed using two cameras. The InfraRed Array Camera (IRAC; Fazio et al. 2004) observes at 3.6 µm, 4.5 µm, 5.8 µm, and 8.0 µm; whilst the Multiband Imager for Spitzer (MIPS; Rieke et al. 2004) observed at 24 µm, 70 µm, and 160 µm. Since depletion of the helium cryogen in 2009, only the IRAC 3.6 µm and 4.5 µm bands are able to operate. For each DustPedia Herschel galaxy, we acquired Spitzer data from the best archive source available. For both IRAC and MIPS, an order of preference was used to decide which data to use, described below. Once data had been retrieved from the preferred source, it was otherwise handled as per Sect. 2.2. The pixel units of all maps were converted from MJy sr−1 to Jy pix−1.

In total, 826 of the DustPedia galaxies (94% of the total) have Spitzer coverage in at least one photometric band (with data in 4 Spitzer bands being the median for those sources that have coverage); 808 (92%) have IRAC coverage, whilst 493 (56%) have MIPS coverage.

Spitzer-IRAC. For IRAC, the preferred data source was the Spitzer Survey of Stellar Structure in Galaxies (S4G20; Sheth et al. 2010; Muñoz-Mateos et al. 2013; Querejeta et al. 2015), hosted by IRSA. S4G was a Spitzer post-cryogenic Exploration Science Legacy Program which provides high-quality dedicated imaging of nearby galaxies at 3.6 µm and 4.5 µm.

When S4G data was not available, the next preferred IRAC data source was the Spitzer Enhanced Imaging Products (SEIP), produced and hosted by IRSA. SEIP images combine Spitzer observations made of a given part of the sky into “Super Mosaics”, reduced using the latest version of the Spitzer pipeline. Whilst SEIP data provides consistent data products for a large fraction of all Spitzer observations, conspicuous imaging artefacts are sometimes encountered; these are most often due to the co-addition and mosaicing process being confounded by edge-of-field effects and/or by artefacts around bright sources.

When SEIP data was not available, IRAC data was retrieved from the Spitzer Heritage Archive (SHA)21. When using SHA data, the procedure laid out in Sect. 2.2 was followed to produce the final cutouts; this process quite closely mimics the way in which the SEIP “Super Mosaics” were produced22. However, SEIP was preferred over SHA, as when IRSA produced the SEIP maps, they were able to tailor their mosaicing process to partially mitigate the effects of Spitzer-specific pathologies (e.g., limiting the influence of artefacts during co-addition).

When performing aperture photometry on IRAC maps, it is necessary to perform calibration corrections, as given in Table 4.8, Sect. 4.11.2 of the IRAC Instrument Handbook23, to account for the fact that IRAC maps are calibrated using aperture photometry of point sources, using apertures of 12” radius. The calibration corrections ensure that the correct fluxes are measures when performing photometry of non-point sources, and/or when using apertures with radii other than 12”24. Additionally, users who wish to perform surface-brightness analyses with the maps will have to consider the surface-brightness corrections described in Sect. 4.11.3 of the IRAC Instrument Handbook. IRAC maps are calibrated assuming a flat νSν reference spectrum, as described in Sect. 4.4 of the IRAC Instrument Handbook.

Spitzer-MIPS. For MIPS, the preferred data source was the MIPS Local Galaxies Program (Bendo et al. 2012), hosted by IRSA, which compiled and re-reduced all Spitzer archive observations of local galaxies that were selected to be observed by 4 Herschel key programs – Herschel Virgo Cluster Survey (HeVICS; Davies et al. 2010, 2012), the Dwarf Galaxy Survey (DGs; Madden et al. 2013), the Very Nearby Galaxy Survey (VNGS), and the Herschel Reference Survey (HRS; Boselli et al. 2010) – in order to provide standardised ancillary data for those surveys. This makes the MIPS Local Galaxies Program the ideal source of MIPS data for the sources it covers (25% of the DustPedia Herschel galaxies).

When MIPS Local Galaxies Program data was not available, the next preferred source was data provided by the Spitzer Legacy/Exploration Science Programs, hosted by IRSA. These are the final data products used and supplied by each program’s team, and are consistently free of artefacts. We made use of MIPS data provided by the C2D (from molecular Cores To planet-forming Disks; Evans et al. 2003, 2009), SEP (Spitzer MIPS 24 and 70 µm imaging near the south Ecliptic Pole; Scott et al. 2010), and SWIRE (SIRTF Wide-area Infrared Extragalactic Survey; Lonsdale et al. 200324) programs.

When neither MIPS Local Galaxies Program nor Spitzer Legacy/Exploration Science Program MIPS data was available, we made use of SEIP MIPS data. Note that SEIP only includes 24 µm MIPS data.

If no other MIPS data source is available, we make use of SHA data in the same manner as for IRAC.

The reference spectrum for MIPS maps has the shape of a 104K blackbody, as described in Sect. 4.3.2 of the MIPS Instrument Handbook25.

2.2.7. Planck

The Planck satellite (Planck Collaboration I 2011) mapped the entire sky in 9 bands, in the 350 µm to 1 cm wavelength range, making Planck well-suited to detecting dust, free-free, and synchrotron emission from nearby galaxies. However, Planck has extremely low resolution compared to the other datasets we employ, with a beam FWHM ranging from 5–32” in size; a factor

20 https://irsa.ipac.caltech.edu/data/SPITZER/docs/S4G/SHA/
21 http://sha.ipac.caltech.edu/applications/Spitzer/SHA/
22 https://irsa.ipac.caltech.edu/data/SPITZER/Enhanced/SEIP/docs/seip_explanatory_supplement_v3.pdf
23 https://irsa.ipac.caltech.edu/data/SPITZER/docs/irac/iracinstrumenthandbook/
25 http://irsa.ipac.caltech.edu/data/SPITZER/docs/mips/mipsinstrumenthandbook
of 8–54 worse than the 36′′ resolution of the SPIRE 500 μm band (which is the lowest-resolution imagery we present aside from Planeck). This stark difference in resolution makes it impractical to perform aperture-matched CAAPR photometry with Planeck maps in the same manner as for the other bands; instead, we will use catalogue photometry provided by the Planck Collaboration (see Sect. 5).

However, for completeness, we nonetheless produced Planeck imaging cutouts for the DustPedia galaxies. Planeck data was accessed from the NASA SkyView service, which uses the all-sky maps26 from the Planeck Public Data Release (Planck Collaboration I 2014).

Due to the exceptionally large size of the Planeck beam, our Planeck imaging cutouts have diameters 4 times larger than our cutouts for other instruments; hence galaxies with optical angular sizes <6′ are in 2′ × 2′ cutouts, whilst those with optical angular sizes >6′ are in 4′ × 4′ cutouts.

All Planeck maps are in our standard units of Jy pix⁻¹. The 350 μm (857 GHz) and 550 μm (545 GHz) cutouts were converted from their native units of MJy sr⁻¹, whilst the other 7 bands were converted from their native units of K cm⁻² rad⁻¹. For the 350–3.00 mm Planeck bands, the data is calibrated assuming a constant vˢ⁻¹, reference spectrum (Planck Collaboration VII 2016); for the 4.26–10.60 mm Planeck bands, the reference spectrum is the Cosmic Microwave Background (CMB), a 2.73 K blackbody (Planck Collaboration II 2016).

3. The CAAPR photometry pipeline

Here we describe the Comprehensive & Adaptable Aperture Photometry Routine (CAAPR), which we use to produce consistent multiwavelength aperture-matched photometry of the DustPedia galaxies from our GALEX, SDSS, 2MASS, WISE, Spitzer, PACS, and SPIRE imagery.

CAAPR is a development of the photometry pipeline employed by Clark et al. (2015), De Vis et al. (2017a,b), and Keenan et al. (in prep.). Its development and functionality are driven by three main motivations.

Firstly, CAAPR was created to be able to generate aperture-matched photometry that is cross-comparable, even in the face of the broad gamut of data that are required for multiwavelength astronomy – producing fluxes and uncertainties in a consistent manner, despite the great variation in the characteristics of observations ranging from the UV to the IR to the submm.

Secondly, CAAPR allows automated standardised photometry for local Universe surveys, as has long been the norm for intermediate-to-high redshift astronomy. Photometry of nearby galaxies generally entails a large degree of “by hand” tweaking, such as altering aperture dimensions to maximise flux and signal-to-noise ratio (S/N), masking of contaminating foreground and background sources, and manually selecting an appropriate background region. This manual approach was taken by, for instance, the Herschel Reference Survey (HRS, see Sect. 6.3.1), the Third Reference Catalogue (de Vaucouleurs et al. 1991), the SCUBA Local Universe Galaxy Survey (SLUGS; Dunne et al. 2000), the 2MASS Large Galaxy Atlas (Jarrett et al. 2003), the Spitzer Infrared Nearby Galaxy Survey (SINGS; Dale et al. 2005), the GALEX Ultraviolet Atlas of Nearby Galaxies (Gil de Paz et al. 2007), the Key Insights on Nearby Galaxies Far-Infrared Survey with Herschel (KINGFISH; Dale et al. 2012), the WISE High-Resolution Galaxy ATLAS (Jarrett et al. 2012), ad nauseam. However, this manual approach is a luxury we cannot afford with DustPedia, as we deal with a larger number of targets, and a much more extensive range of bands, than is typically the case for nearby galaxy surveys. As such, CAAPR requires minimal human interaction once started (except for uncommon occasions where visual inspection leads us to exclude a certain band from contributing to a source’s aperture fitting, or disable star removal for a particular map, etc.). This automated approach allows us to handle the scale of the DustPedia dataset, and be confident that all of our measurements have been conducted in a consistent and cross-comparable way.

Thirdly, CAAPR is designed to be able to handle the wide range of shapes, sizes, brightness distributions, and other morphological traits found in nearby galaxies (and to do so across the very different resolutions and noise environments found in multiwavelength datasets). Naturally, aperture photometry requires that the shape and size of a target source is well constrained, in order to construct an appropriate aperture – however this is notoriously difficult for nearby galaxies, which are often “shredded” (i.e., erroneously identified as multiple separate sources). The widely-used SEXTRACTOR software (Bertin & Arnouts 1996) is prone to shredding extended sources (Hammer et al. 2010a,b; Wright et al. 2016), and the shredding of nearby galaxies in the SDSS is well-documented (Blanton et al. 2001; Hall et al. 2012; Budavári et al. 2009). Indeed, when describing the LAMBDA photometry pipeline (which is designed for higher-redshift galaxies), Wright et al. (2016) give nearby galaxy NGC 5690 as an example of where their multi-pass SEXTRACTOR-derived aperture generation fails – even though NGC 5690 is relatively “well-behaved” from the standpoint of nearby galaxy aperture fitting, being bright, edge-on, and only somewhat flocculent (by coincidence, NGC 5690 is one of the galaxies in the sample of Clark et al. 2015, for which CAAPR was originally developed, and where CAAPR’s aperture generation performs successfully).

CAAPR takes the form of a Python 2.7 package, and is available on GitHub28. CAAPR is operated via a single high-level function, with two primary inputs – a table describing the properties of each source, and a table describing the properties of each band that is to be processed – along with a number of optional arguments (e.g., stating whether CAAPR is to operate using parallel processing, or specifying that only the aperture-fitting phase of the pipeline should be run). The input tables allow the user to alter how the pipeline runs for individual sources (e.g., the user might want to exclude the r-band map of a given source from contributing to the shape of the final photometric aperture, due to contamination from a satellite trail), and set up the pipeline to run appropriately in each band (e.g., the user is likely to want to enable foreground star removal at 3.4 μm, but disable it at 500 μm).

3.1. Preprocessing

CAAPR can operate with separate input maps provided for each target in a given band, or with a single input map for all targets in a given band. In either case, the user has the option of

26 Although the Planeck all-sky maps are in the HEALPIX projection (Górski et al. 2005), the maps provided by SkyView are in the standard gnomonic TAN projection.
27 https://irsa.ipac.caltech.edu/data/Planck/ercsc_v1.3/explanatory_supplement_v1.3.pdf
28 https://github.com/Stargrazer82301/CAAPR
extracting a cutout map centred, on the target source, which will then be used throughout the rest of the pipeline (this is often necessary if the input map is particularly large, and hence memory-intensive); this cutout will be reprojected to a gnomonic TAN projection. Note that for purposes of consistency, the maps being processed by CAAPR will be referred to as cutouts for the rest of this section.

### 3.2. Foreground star removal

Contamination from foreground stars in UV-MIR bands is minimised using the foreground star removal functionality included in the Python Toolkit for SKIRT (PTS\(^{29}\); Camps et al. 2015; Verstocken et al., in prep.).

The foreground star removal procedure works on a band-by-band basis and runs semi-automatically, requiring a small amount of adjustment of the configuration parameters. For a particular image, the appropriate entries from the 2MASS All-Sky Catalog of Point Sources (Cutri et al. 2003) are retrieved as a starting point. The catalogue is queried through the VizieR interface of Astroquery\(^{30}\) Python library. Around each catalogued position, a small patch of the image is subtracted by the estimated background in that area and subsequently a local peak is searched, using tools from the Photutils\(^{31}\) Python package. If a peak with a reasonable S/N is not found within a radius of a few pixels, the position is ignored. If a matching peak is found, it has to show no deviations from being a true point source (other than the peak exhibiting saturation in the case of the brightest stars), otherwise it will also be ignored. This is an effective way of preventing most compact sources in the target galaxies, primarily HII regions, from being removed – as such objects tend to lie within local peaks in a galaxy’s brightness distribution. For example, even in the uncommon cases where HII regions display no extended nebulosity at all, they generally lie along the brightness “ridge” of a galaxy’s spiral arms; this makes their brightness profile deviate from being a true point source.

The FWHM of a positively identified point source determines the size of the area which will be masked and replaced with an interpolation. For the brightest point sources, image segmentation is used to detect saturation bleed, ghosts, and diffraction artefacts. The segmentation is also replaced with an interpolation. This interpolation is performed by approximating the source neighbourhood by a 2-dimensional polynomial function, whilst including Gaussian noise based on the average deviation of the pixel values from this polynomial. Figure 3 shows a WISE 3.4 \(\mu\)m map of IC 0342 before and after undergoing the foreground star removal process.

The effectiveness of the foreground star removal is explored in-depth in Sect. 6.1. In uncommon instances where visual inspection revealed that the foreground star removal caused problems (such as confusing a bright HII region in the target galaxy for a star, and hence removing it), then CAAPR was run with the star removal disabled; this was required for approximately 1–2\% of cases.

### 3.3. Polynomial sky flattening

After star-subtraction, CAAPR removes large-scale sky structure in each cutout, if necessary, by means of fitting and subtracting

\(^{29}\) http://www.skirt.ugent.be/pts/
\(^{30}\) https://github.com/astropy/astroquery
\(^{31}\) https://github.com/astropy/photutils

![Fig. 3. Demonstration of our foreground star removal process, using the example of WISE 3.4 \(\mu\)m observations of IC 0342, over a 45’ × 30’ area of sky, shown before (upper) and after (lower) undergoing star removal. Both images are displayed using the same logarithmic brightness scale, to allow both bright and faint features to be seen. Photometric master aperture ellipse, as determined by CAAPR, is marked in green. IC 0342 lies in the plane of the Milky Way, well within the zone of avoidance (b = 10°), leading to a high degree of foreground stellar contamination. Nonetheless, our foreground star removal process successfully subtracts the vast majority of all emission from foreground stars in the vicinity of IC 0342.](image-url)
it is unnecessary, and at worst it may actually degrade the sky quality. Therefore, CAAPR evaluates the cutout’s pixel distribution before and after the sky polynomial is subtracted. For each distribution, only pixel values below the peak are considered in this evaluation, as these values should represent pixels corresponding to “empty” sky. In cases where the sky polynomial subtraction improves the sky quality – flattening the sky – the range of pixel values below the distribution peak should be narrowed, as there should be a narrower range of pixel values corresponding to empty sky. To quantify this width, CAAPR finds the root-mean-squared (rms) deviation between the peak pixel value and all of the values beneath it, for both the polynomial-subtracted and polynomial-unsubtracted distributions. If the rms width for the subtracted distribution is at least 10% smaller than that of the unsubtracted distribution, it is assumed that the polynomial subtraction has lead to the sky being appreciably flattened, and the polynomial-subtracted cutout is utilised for all subsequent stages of the pipeline; otherwise, the original cutout is used, unaltered. The polynomial sky subtraction process is demonstrated in Fig. 4, using the example of NGC 6946. Note that the polynomial sky flattening is only conducted as part of our CAAPR photometry; the FITS images provided in our data products (see Sect. 7.3) have not undergone this process.

3.4. Aperture fitting

CAAPR’s aperture fitting process generates an elliptical aperture in every band for a given target; these apertures are then combined (as described in Sect. 3.5) to provide a final aperture for the target. The aperture fitting process consists of two distinct stages; first determining the position angle and axial ratio (i.e., the shape) of the aperture ellipse, then second determining the semi-major axis (i.e., the size) of the aperture ellipse. The stages of the aperture-fitting process are illustrated in Fig. 5.

To determine the position angle and axial ratio of the aperture ellipse, CAAPR starts by estimating the noise in the cutout, by finding the iteratively sigma-clipped (with a 3σ threshold) standard deviation of the cutout’s pixel values. CAAPR then identifies all contiguous groups of five or more pixels with a \(S/N > 3\) (Panel 2 of Fig. 5); if there is a group of contiguous \(S/N\) pixels associated with the target source, it is isolated (Panel 3 of Fig. 5).

If no group of \(S/N > 3\) pixels is found at the coordinates of the target source, it is assumed that the source is non-detected and CAAPR defaults to a circular aperture (i.e., an axial ratio of 1, and a position angle of 0 deg).

If CAAPR succeeds in finding and isolating a group of contiguous \(S/N > 3\) pixels associated with the target source, then CAAPR subjects that group of pixels to binary opening (i.e., a succession of binary erosion then binary dilation), with an erosion element of size equal to the FWHM of the band’s beam (Panel 4 of Fig. 5). This removes small protrusions from the periphery of the group of pixels, as such features are often due to background sources, foreground stars, map artefacts, etc. For example, in Panel 3 of Fig. 5 the removed feature at the north-easterly edge of the isolated group of pixels is associated with a background galaxy.

Next, CAAPR finds the convex hull (the tightest polygon that will contain a set of points) of the remaining pixels (Panel 5 of Fig. 5), then fits an ellipse to the vertices of the convex hull (Panel 6 of Fig. 5); the position angle and axial ratio of this ellipse will be used for the final aperture ellipse for the band in question. The semi-major axis of this ellipse will be employed by CAAPR as an initial reference value for determining the size of the aperture ellipse, as described below.

Before determining the semi-major axis of the aperture ellipse, CAAPR convolves a copy of the cutout with a Gaussian filter, of FWHM twice the beam width of the band in question. This serves to enhance any lower-surface-brightness features associated with the target source. This is particularly important for sources that are very well resolved, where individual regions of emission within the source can be separated by apparently “empty” sky (e.g., individual star-forming regions in FUV for well-resolved face-on spirals); the Gaussian smoothing helps to mitigate this effect.

The semi-major axis of the aperture ellipse is determined by placing concentric elliptical annuli on the smoothed cutout, centred at the coordinates of the target source, with the position angle and axial ratio determined previously. Initially, 10 such elliptical annuli are placed, with semi-major axes ranging from 0.5–3.0 times the initial reference value determined above, and with widths such that the boundaries of each subsequent annulus are in contact. The average per-pixel S/R in each elliptical annulus is calculated, to identify the pair of annuli that bracket the semi-major axis at which \(S/N = 2\) (Panel 7 of Fig. 5).

Within this range of semi-major axes, a differential evolution optimisation (Storn & Price 1997), using elliptical annuli of one beam-width, is then used to find the precise semi-major axis at which \(S/N = 2\) (Panel 8 of Fig. 5). This two-step approach is much faster than a full high-resolution radial profiling. The semi-major axis of the \(S/N = 2\) isophote is then multiplied by
an expansion factor of \(1.25^{32}\) to give the final semi-major axis of the aperture ellipse for the band in question (Panel 9 of Fig. 5); applying this expansion factor accounts for the fact that some fraction of a source’s flux will always fall outside any practical S/N cutoff (Overcast 2010; Blanton et al. 2001). If no annulus had \(S/N > 2\), then we revert to a default semi-major axis of twice the beam FWHM in the current band (note that, by design, all galaxies in the DustPedia sample are detected in at least one band).

3.5. Aperture combining

The aperture-fitting process described in Sect. 3.4 is repeated for every band for which data is available for a given target. We include DSS data in our aperture-fitting, as it provides consistent optical coverage for all of our galaxies (particularly useful for determining the stellar dimensions of galaxies without GALEX and/or SDSS data); although we exclude DSS data from the actual photometry as the maps are in photometrically uncalibrated units of photographic densities. To allow a valid comparison despite a potentially wide range of resolutions, the beam-width in each band is subtracted in quadrature from the major and minor axes of that band’s aperture ellipse (the axial ratio is recomputed accordingly), effectively deconvolving them. These corrected ellipses are compared, to find a “super-ellipse” – the smallest ellipse which will contain all of the deconvolved ellipses. This is the master aperture for this source (Panel 10 of Fig. 5), to be used for the aperture photometry and uncertainty estimation described in Sects. 3.6 and 3.7.

The user can intervene to exclude specific bands from the aperture fitting and combining process, either for an individual target or all targets, by adjusting the corresponding entries in the input tables passed to CAAPR.

3.6. Aperture photometry

The photometric measurements made by CAAPR take the form of standard aperture photometry, with local background subtraction conducted using an elliptical background annulus around the elliptical master aperture, both centred on the coordinates of the target source. The user can opt to use apertures generated by CAAPR as per Sects. 3.4 and 3.5; or the user can provide their own apertures.

An example of the final master aperture for a given source, applied to all bands, is shown in Fig. 6 for NGC 4559.

In general, for earlier-type galaxies (types E–Sa), the final master aperture size was often driven by the WISE 3.4 and 4.6 \(\mu m\) bands, thanks to their sensitivity to low-surface-brightness NIR emission (earlier-type galaxies being dominated by evolved stellar populations which are luminous in the NIR). Conversely, the master apertures of later-type galaxies (types Sb–Sm) were typically dictated by the apertures fit in the GALEX and UV-optical SDSS bands.
Fig. 6. Photometry thumbnail image grid for NGC 4559. CAAPR produces an image like this for every source. Solid lines show master aperture, whilst dashed lines demark background annulus. Several of the fluxes for NGC 4559 have flags (flags are defined in Sect. 4). WISE 12 µm flux has “A” flag due to negative flux artefact in master aperture, whilst WISE 22 µm flux has “a” flag, due to a less pronounced version of the same feature. The Spitzer 3.6, 4.5, 5.8, and 8.0 µm fluxes have “N” flags (although these are borderline cases), as <40% of pixels in their sky annuli have coverage, so it is assumed that any background estimate will be unreliable (similarly, the lack of coverage in their master apertures would, independently, warrant “n” flags). Spitzer 24 µm and PACS 70, 100, and 160 µm fluxes have “n” flags, as <80% of pixels in their sky annuli have coverage, potentially reducing reliability of background estimate.

(in particular, we repeated the photometry for all sources in the GALEX FUV and WISE 12 µm bands; these data have sufficiently high resolution for convolution to 36ʺ to represent a significant change, without being computationally excessive). The photometry obtained in these repeats agreed to within 1%, with the typical difference in flux being 0.4σ.

The flux in a source aperture is simply the sum of the values of all the pixels it contains. Bands with larger pixel sizes require consideration of partial pixels; CAAPR determines what fraction of each pixel’s flux falls within the source aperture by dividing each pixel into a number of sub-pixels (a standard technique; see Bourne et al. 2013; Wright et al. 2016, and references therein). For the DustPedia photometry we divide each input pixel into 10 × 10 grids of 100 sub-pixels for the Spitzer 70 and 160 µm bands, all three PACS bands, and all three SPIRE bands.

The background annulus for a given target and band has the same axial ratio and position angle as the master aperture it surrounds. The inner and outer semi-major axes of the annulus are at 1.250 and 1.601 times the beam-adjusted semi-major axis of the source aperture (Panel 10 of Fig. 5); these proportions mean that the background annulus has the same area as the master aperture (although they can be adjusted by the user). The average background level in the annulus is determined by taking the iteratively sigma-clipped (with a 3σ threshold) mean of the pixel values contained within the annulus, with partial pixels considered in the same manner as for the master aperture. This average background is then subtracted from the flux measured in the master aperture, appropriately weighted for the aperture’s area.

For maps so small that the background annulus lies wholly outside their coverage (for example, in the instance of some Spitzer and Herschel observations which observed only the nucleus of the target galaxies), then no background can be measured, and so a null flux is recorded for the source.

It is possible for a “negative” flux to be recorded, if the locally-determined background is brighter than the target source. This is to be expected; consider sources where there is essentially no flux present, such as SPIRE observations of dust-free early-type galaxies. In these cases, the flux level measured in the background annuli should, on average, be the same as the flux level measured in the elliptical master apertures. However, they will never be perfectly identical for any individual instance, due to aperture noise – so 50% of the time there will be a brighter flux level in the background annulus than in the elliptical master aperture, which will hence result in a negative flux measurement. Negative fluxes can also occur when measuring faint sources in fields with complex backgrounds (e.g., Galactic cirrus in MIR–submm bands). The vast majority of negative flux measurements are compatible with a flux of 0 when the uncertainty is considered.

For bands with beam FWHM > 10ʺ, an aperture correction is applied to account for the fraction of the source flux spread outside the master aperture by the PSF. Most instrument handbooks only provide such corrections for point sources, as corrections for extended sources (such as the DustPedia galaxies) require a model for the underlying unconvolved flux distribution. CAAPR assumes that each target galaxy, as observed in a
CAAPR needs to not only be able to measure photometric instrumental noise in the MIR, to confusion noise in the submm. Uncertainties in a wide range of noise regimes; from photon-

Multiwavelength photometry requires determining photometric uncertainties in all of these situations, but to do so in a legitimately cross-comparable manner. As such, CAAPR takes a robust “brute force” approach to estimating uncertainties, by measuring the variation in the fluxes recorded in sky apertures placed in the vicinity of the target master aperture. The aperture noise measured in this way will encompass the uncertainty contributions arising from instrumental variation in pixel values, and from confusion, and from the risk of unrelated sources contaminating the aperture, and from large-scale noise such as sky brightness and Galactic cirrus.

The random sky apertures all have the same area as the master aperture, and are background subtracted using an encircling background annulus in the same manner as for master aperture. CAAPR places each random sky aperture quasi-randomly on the cutout, with its coordinates being drawn from a Gaussian distribution centred on the target source, with a standard deviation five times the semi-major axis of the source aperture. This distribution means that the random sky apertures are most likely to be positioned in the immediate vicinity of the master aperture, where the measured aperture noise will be most pertinent.

However, each randomly generated sky aperture has to pass three tests before it is accepted for use. Firstly, no random sky aperture is permitted to overlap the master aperture; otherwise the inclusion of the target source in the aperture noise estimation would artificially drive up the reported uncertainty. Secondly, no random sky aperture is permitted to have more than 50% of its pixels be sampled by any other random sky aperture; if a large number of random sky apertures happened to be positioned in the same location, this would artificially reduce the reported aperture noise. Thirdly, no random sky aperture may extend beyond the boundaries of the cutout. Figure 7 shows the distribution of accepted and rejected random sky apertures for the example case of NGC 5147.

CAAPR keeps generating random sky apertures, and accepting those that pass the three tests listed above, until certain criteria are met. If 100 random sky apertures are successfully generated, then the generation process stops and those apertures will be used to determine the aperture noise. However, if the situation arises where 250 random apertures have been generated in a row without a single one of them being accepted, it is assumed that there are no more valid positions left in the cutout for any more apertures to be successfully placed; if this occurs, one of two things can happen. If at least 50 random sky apertures have already been accepted, these apertures will be used to determine the aperture noise. But if fewer than 50 random sky apertures have been accepted, this is deemed an insufficient number, and aperture noise will instead be determined using the mini-aperture extrapolation method described in Sect. 3.8.

This method of determining aperture noise has the additional benefit of being highly stable. Running repeated aperture noise estimates on the same cutout yields aperture noise estimates that typically vary by <10%. Compare this to the high degree of variation – as much as a factor of several – that can be suffered when placing sky apertures across a square cutout totally at random (Clark 2015).

Once a sufficient number of random sky apertures has been successfully generated, CAAPR calculates the iteratively sigma-clipped (with a 3σ threshold) standard deviation of the distribution of flux values measured from those apertures. This value is taken to be the aperture noise for the aperture size and cutout in question.

The final uncertainty for a given flux is calculated by adding in quadrature the aperture noise and photometric calibration uncertainties.
CAAPR first finds the largest size of random mini-aperture that can successfully be placed in sufficiently large numbers, and then measures the aperture noise for that size mini-aperture, doing so in the exact same way as laid out in Sect. 3.7 – with the only difference being that a smaller radius is being used. CAAPR then places a set of even smaller mini-apertures across the map, and likewise determines the aperture noise for mini-apertures of that size. CAAPR repeats this process, measuring the aperture noise for successively smaller sets of mini-apertures. The radii of each successive set of mini-apertures are evenly-spaced in log space, following a power law of index 1.2 (an empirically-selected value, which provides sampling dense enough to be effective, whilst coarse enough to still be reasonably fast).

CAAPR keeps on calculating the aperture noise associated with smaller and smaller sizes of mini-aperture, stopping when one of two conditions is reached; until 10 increments of mini-aperture size have been successfully processed, or until the mini-aperture radius has fallen to 1 pixel (making smaller apertures impossible).

CAAPR then produces a log-log plot of the aperture noise determined for each set of random mini-apertures, against the mini-aperture area used for each set. An example of such a plot, for the SPIRE 250 μm photometry of M 101 (NGC 5457) is shown in Fig. 8. In these area-noise plots, there is consistently a strong correlation between the area of the random mini-apertures, and their associated aperture noise. CAAPR uses these area-noise plots to extrapolate what the aperture noise associated with the area of the full-size random sky apertures would have been, had there been sufficient map area available; this extrapolation is performed by fitting a power law to the area-noise points, in a chi-squared minimising manner. For fitting, the error bar on each point is derived from the Poisson uncertainty given the number of random mini-apertures employed, which is then multiplied by a weighting that takes the form $1 + \log_{10}(\Delta_A)$, where $\Delta_A$ is the factor difference between the areas of the random mini-apertures and the master aperture. This empirical weighting accounts for the fact that random mini-apertures with areas more similar to that of the master aperture should have greater influence over the extrapolated aperture noise than random mini-apertures with areas a lot smaller than the master aperture.

We demonstrate the validity of this approach, and quantify its effectiveness, in Sect. 6.2. This method allows CAAPR to produce reliable estimates of aperture noise for almost any photometric measurement; if there is sufficient map area to determine a flux, there is sufficient map area to estimate the aperture noise, in all but the most extreme cases. In the very rare instances where there is not enough map area for even the mini-aperture approach to work, we record a negative uncertainty in our photometry; this indicates that that uncertainty represents the calibration uncertainty alone.

### 3.8. Aperture noise extrapolation

The approach of determining aperture noise via the placement of randomly-positioned sky apertures has the virtue of being effective for essentially any noise regime (e.g., instrumental noise, confusion noise, contaminating foreground/background sources, Galactic cirrus, etc.). However, it has the distinct disadvantage that it requires maps with a sufficient amount of coverage area around the target source to permit the placement of enough random sky apertures.

Moreover, consider the case of an extremely extended nearby galaxy, such as M 101 (NGC 5457), for which we employ a photometric aperture with a major axis of 35′. For a map of M 101 to be large enough for the placement of 50 random sky apertures (our threshold minimum acceptable number) around the target, with perfect tessellation, the map would have to cover an area of over 13 sq deg. Whilst large-area surveys like SDSS might observe an area of that size, most observations by telescopes like *Spitzer* and *Herschel* would not map even a fraction of that amount of sky when observing a galaxy like M 101 (for instance, Fig. 6 shows how small the *Spitzer* and *Herschel* maps of NGC 4559 are in relation to the size of the actual galaxy). And even if maps were consistently that large, it would not be appropriate to use such a vast area of sky for determining aperture noise; an aperture noise estimate determined using random sky apertures placed several degrees away from the target will not be representative of the aperture noise at the location of the target itself. To remedy this, CAAPR features a novel method of extrapolating the aperture noise, for maps of any size.

In instances where CAAPR is unable to place a sufficiently large number of full-size random sky apertures on a cutout (as described in Sect. 3.7), CAAPR will instead switch to placing random “mini-apertures”, of a range of sizes, across the map.

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**Fig. 8.** Example of how CAAPR extrapolates aperture noise, for SPIRE 250 μm photometry of M 101 (NGC 5457). Blue circles indicate aperture noise determined when using mini-apertures of different sizes. Note that the error bars do not represent uncertainties, but rather indicate weightings assigned to each point, with greater weight given to mini-apertures closer in size to the master aperture. Blue dashed line indicates best-fit power law to trend exhibited by these points. Red hexagon indicates extrapolated aperture noise for M 101, given area of its master aperture.

4. Photometry flagging

The CAAPR photometry of the DustPedia galaxies provides an enormous number of fluxes, produced in a consistent manner between galaxies and across bands. The size of our dataset makes it impractical to manually adjust the photometry in the case of individual sources where issues may have been encountered (and regardless, such tweaking would risk introducing inconsistencies). Instead, we opted to manually inspect all of the photometry points for the DustPedia galaxies, providing an example of which is shown in Fig. 6 for...
NGC 4559. At least two members of the DustPedia team inspected the thumbnail grid for each galaxy, along with the corresponding fluxes, and recorded flags for instances where issues were encountered.

Three different types of flag are used to indicate when there may be issues with a particular flux in our aperture-matched photometry (along with an additional flag used only for supplementary photometry, discussed below). Each flag is split into “minor” and “major” categories. Minor flags are indicated by lowercase letters, whilst major flags are indicated by uppercase letters. The types of flag are defined as follows:

\( a/A \) – Artefact flag. This indicates that some sort of map artefact is present in the region of the target source, as determined by visual inspection. Examples of this include poor mosaicing in the archival data, satellite trails, or evidence of saturation in bright pixels.

\( c/C \) – Contamination flag. This indicates contamination of the photometry due to a nearby source, as determined by visual inspection. Examples of this include foreground stars that were not fully removed, or other galaxies close to the target source.

\( n/N \) – Null-coverage flag. This flag is automatically assigned where necessary. This indicates that the observation did not provide full coverage of the area of the source, limiting CAAPR’s ability to suitably measure the source flux and/or background. If more than 5% of the pixels in the master aperture and/or more than 20% of the pixels in the background annulus are found to be NaN – “Not a Number” – then a minor flag is raised. If more than 20% of the pixels in the master aperture and/or more than 60% of the pixels in the background annulus are found to be NaN, a major flag is raised. Instruments for which small maps are common (such as Spitzer and PACS) are particularly likely to have fluxes with this flag.

\( e/E \) – Excluded extended emission flag. This flag is only used with our supplementary photometry (as described in Sect. 5), and is not assigned to any of our aperture-matched CAAPR fluxes. It indicates that the target is sufficiently extended that there is a risk some of its emission will have been excluded from the measurement in question, which would lead to the flux being an under-estimate.

A major flag indicates that the associated flux probably is not suitable for general use – although it may still be useful for specific considerations. For example, the most common reason for a “C” flag is because the target galaxy is interacting with another galaxy inside the master aperture; although the recorded flux will not be of use as a measure of the emission from the target galaxy alone, it will nonetheless be useful to anyone interested in the flux of the overall system.

A minor flag indicates that the associated flux is probably still useful for most purposes – although it is advisable for users to inspect the corresponding data to confirm that the flux is still valid for their specific applications.

Some sources have a “global” flag applied. In this case, an issue affects so many bands that all bands receive the relevant flag. The issue may not be apparent in every band – for example, if the target galaxy is a spiral that overlaps with an elliptical companion, the elliptical might not be visible in SPIRE bands. Nonetheless, if an issue affects enough bands, all bands get flagged by default. Only “c/C” flags can be applied globally.

The flagging is not intended to cover the types of minor issues which are encompassed by the quoted uncertainties – minor background objects within the master aperture, foreground Galactic cirrus, etc. Rather, the flagging is designed to highlight aberrations that uncertainties would not reflect.

Amongst the small number of our sources for which a negative flux is measured (see Sect. 3.6), the flux can occasionally be “significantly” negative, and hence unphysical. This can arise due to a particularly inopportune distribution of Galactic cirrus surrounding the target source, for example. In the rare instances where the flux has a \( S/N < -2 \), we automatically assigned an “A” flag. This was necessary for 0.4% of our fluxes.

In total, 4081 (22.4%) of our aperture-matched CAAPR fluxes have flags associated with them: 2262 (12.4%) have major flags, whilst 1819 (10.0%) only have minor flags. Whilst this might seem high, a large fraction of these are due to “N” or “n” flags being automatically assigned to Spitzer fluxes, because of the very small map sizes of many Spitzer observations; excluding Spitzer photometry, only 2234 (14.6%) of our fluxes possess flags, of which 1202 (7.8%) have major, and 1032 (6.8%) have only minor.

Some examples of issues for which flags are raised are illustrated in Fig. 6, for NGC 4559.

5. Supplementary photometry

5.1. Planck CCS2

5.1.1. Planck CCS2 photometry

Planck (Planck Collaboration I 2011) photometry was obtained from the Second Planck Catalogue of Compact Sources (Planck CCS2; Planck Collaboration XXV 2016). The poor resolution of Planck made it impractical to include in our aperture-matched CAAPR photometry. Entries in the Planck CCS2 (from the full range of Galactic latitudes) were matched to DustPedia sources within a matching radius equal to the FWHM of each Planck band.

The Planck CCS2 provides photometry in all 9 Planck bands, with a focus on source reliability. The Planck CCS2 provides 4 different flux measurements for each source. Of these different fluxes, Planck Collaboration XXV (2016) find that their aperture photometry (the APERFLUX field in the published Planck CCS2 tables) is the measure that compares best with Herschel aperture photometry of HRS galaxies. Therefore the Planck CCS2 aperture photometry values are the fluxes we provide here. The Planck CCS2 fluxes are not colour corrected, and assume the standard Planck reference spectra (see Sect. 2.2.7).

The calibration uncertainties of each Planck band (as provided in Planck Collaboration V 2016; Planck Collaboration VII 2016, including the systematic uncertainties quoted for the 350 and 550 \( \mu m \) bands) are listed in Table 1, and were added in quadrature to the photometric uncertainties given for each source, to produce the final flux uncertainties.

In total, the Planck CCS2 provides an additional 1079 fluxes for the DustPedia Herschel sample. The number of detections ranges from 394 (45% of the DustPedia galaxies) at 350 \( \mu m \), to just 11 at 4.26 \( \mu m \). Perhaps the most useful Planck band for our purposes is 850 \( \mu m \), as it is the most sensitive Planck band to cover wavelengths longer than those observed by Herschel; 197 (22%) of the DustPedia galaxies are detected at 850 \( \mu m \).

5.1.2. Planck CCS2 flagging

The Planck CCS2 aperture photometry uses circular apertures with radii equal to 1 FWHM in each band. Emission more extended than 1 FWHM from the target source’s centre will therefore be lost (not counting flux that falls outside the aperture...
due to the instrumental PSF, which the CCS2 corrects for). To identify which of our sources are vulnerable to this effect, we first measured the angular scale of the FIR/submm emission of the DustPedia galaxies, by running the aperture-fitting phase of CAAPR (see Sect. 3.4) on the SPIRE 250 µm map of each target (or the PACS 160 µm map for targets without SPIRE coverage); no expansion factor was applied, so the semi-major axes of the resulting ellipses represent the actual extent of the detectable FIR/submm emission associated with each target. In each Planck band, sources where the semi-major axis of this emission ellipse was larger than the CCS2 aperture radius (being the effective FWHM listed in Table 2 of Planck Collaboration XXV 2016) were identified, and the corresponding fluxes flagged with an “e” flag – indicating a source’s emission may be extended beyond the region being measured, leading to missed flux. Planck’s resolution at 4.26, 6.81, and 10.60 mm is poor enough that no sources were identified as being at risk of losing flux in this way. At 1.380 and 2.100 mm, >45% of fluxes required flagging. In all other bands, <27% of fluxes are affected. Note that the presence of an “e” flag does not necessarily mean any flux has been missed, as the emission scales in the Planck bands will not be identical to the FIR/submm flux extent. A difference in emission scales is particularly likely at the longer Planck wavelengths, where emission processes are different.

If a “C” or “c” flag was associated with the SPIRE 350 or 500 µm photometry for a given source, as per Sect. 4, then that flag was propagated to all Planck submm fluxes (i.e., 350, 500, or 850 µm). Additionally, we visually inspected the 350 µm SPIRE maps for each source with Planck submm photometry, to look for any potential contaminating source within the CCS2 aperture, and applied “c” and “C” flags as appropriate. We do not apply contamination flags derived from submm wavelengths to the longer-wavelength Planck bands, due to the very different emission processes that dominate there. Moreover, the much lower source density at the longer Planck wavelengths – even when accounting for the larger Planck beam and CCS2 aperture sizes – means that there is a much smaller contamination risk.

5.2. IRAS SCANPI

5.2.1. IRAS SCANPI photometry

Photometry from the InfraRed Astronomical Satellite (IRAS; Neugebauer et al. 1984) was acquired using the Scan Processing and Integration Tool (SCANPI)35), has extremely poor resolution, and the highly asymmetrical PSF; made it impractical to perform aperture-matched photometry using IRAS maps. Nonetheless IRAS photometry is valuable for DustPedia, particularly in the 60 µm band; fluxes in this part of the spectrum are important for constraining dust emission and properties (see Jones et al. 2016, 2017, and references therein).

In each of the 4 IRAS bands, a calibration uncertainty of 20% (Miville-Deschênes et al. 2005; Sauvage 2011) was added in quadrature to each photometric uncertainty. IRAS fluxes are calibrated assuming a ν−1 reference spectrum.36 In total, IRAS provides an additional 2533 fluxes for the DustPedia Herschel sample.

SCANPI operates by fitting the IRAS timelines. Indeed, it was not originally intended that IRAS observations would be used to be produced imagery; whilst IRAS maps do of course exist, the PSF asymmetry necessitates dramatic down-sampling of the data, and the pointing of the resulting maps suffers from some variability37). For these reasons, and in order to avoid presenting maps that are at odds with our photometry, we opt not to include IRAS imagery in the DustPedia database.

5.2.2. IRAS SCANPI flagging

Because SCANPI works by fitting to the IRAS timelines, the resulting fluxes cannot necessarily be directly compared to the aperture-matched photometry. In the case of sources with large angular sizes, there is a risk that SCANPI will miss faint extended emission, especially if the emission distribution of the target source is not smooth; in this scenario, SCANPI might fit only the bright central region of a source, essentially shredding it. A good example of this is M81 (NGC 3031), where the SCANPI 60 µm flux represents only the flux associated with the bright central bulge.

We consider any source more extended than twice the resolution of IRAS in a given band to be at risk of being shredded by SCANPI in this way. Given the highly asymmetrical resolving power of IRAS, we assume a worst-case scenario by taking the best possible resolving power in each band37. IRAS has a maximum resolution of 0.75” at 12 and 25 µm, 1.5” at 60 µm, and 3” at 100 µm. Flux densities identified as being vulnerable to shredding were flagged with an “e” flag, indicating that some emission may be extended beyond the region being measured.

Given that the DustPedia sample consists only of sources larger than 1”, all 12 and 25 µm SCANPI fluxes receive “e” flags. However, since we have WISE aperture-matched photometry covering both the 12 and 25 µm bands, at greater resolution and sensitivity than provided by IRAS, we envisage no scenario where the SCANPI photometry in these bands would be preferred; however, we include it for completeness. Note, also, that SCANPI photometry does not benefit from the foreground star removal used for our CAAPR aperture-matched photometry.

To determine which 60 and 100 µm fluxes are at risk of shredding, we consider the FIR/submm extent of each DustPedia target, determined in the same manner as in Sect. 5.1.2. Any galaxy where the major axis of the FIR/submm emission ellipse is larger than two IRAS resolving elements is deemed to be vulnerable to being shredded. Of the 675 DustPedia targets with SCANPI 60 µm fluxes, 254 (36%) exceed this threshold and hence get flagged; however, 81 of them have reliable (i.e., no major flags) Spitzer 70 µm and/or PACS 70 µm fluxes. As a result, 502 (75%) of our sources have good photometry at 60–70 µm. Of the 682 DustPedia targets with SCANPI 100 µm fluxes, 97 (14%) are extended enough in the FIR/submm to require flagging; of these, 68 have reliable PACS 100 µm fluxes. Therefore 713 (81%) of our galaxies have good 100 µm photometry.

Contamination flags from our aperture-matched photometry were propagated to the corresponding SCANPI fluxes. For SCANPI 12 µm photometry, if a given source had a “C” or “c” flag associated with its WISE 12 µm flux, then that flag was applied to the SCANPI flux also. Similarly, contamination flags associated with our WISE 22 µm photometry were applied to the corresponding SCANPI 25 µm fluxes. For SCANPI 60 µm fluxes, we preferred to take contamination flags from PACS 70 µm; if PACS 70 µm photometry was not available, we instead used flags from, in order of preference: PACS 100 µm.

35 http://irsa.ipac.caltech.edu/applications/Scanpi/
36 IRAS Explanatory Supplement: https://lambda.gsfc.nasa.gov/product/iras/docs/exp.supp/
37 Band resolving powers taken from the Survey Array table in the Instrument Summary: http://irsa.ipac.caltech.edu/IRASdocs/iras_mission.html
PACS 160 µm, or SPIRE 250 µm. For SCANPI 100 µm photometry, we used the same order of preference, except that PACS 100 µm flags were preferred over PACS 70 µm flags.

6. Photometry validation

6.1. Foreground star removal validation

We tested the foreground star removal method we described in Sect. 3.2 by using it to remove the stars from observations of 16 patches of sky at a range of wavelengths.

We produced cutout maps for 16 patches of sky in SDSS u, SDSS r, 2MASS J, and WISE 3.4 µm, following the procedures laid out in Sect. 2.2; this selection of data should provide a fair sampling of the range of wavelengths and noise regimes in which we perform star subtraction in our photometry. In all bands, each of the 16 maps had dimensions of 0.5° × 0.5°, the same size as our standard cutout images. These test cutouts were all located within 3° of coordinates α = 136.6° 6 δ = 0.45° (J2000); this area of space is known to be extremely under-dense at low-redshifts (Driver et al. 2011), and contains no large galaxies in the volume sampled by DustPedia. Each of the maps was also visually inspected to confirm the absence of any nearby galaxies. As such, the vast majority of the emission seen in each of these maps should come from foreground stars, with only a small fraction coming from distant background galaxies, etc.

To remove any large-scale emission from the test cutouts, we conducted a polynomial sky subtraction, as per Sect. 3.3. Having done this, we assumed that the background level associated with “empty” sky was represented by the peak of the pixel value distribution; for each test cutout, we found this level and subtracted it. Having performed this background subtraction, the sum total of the pixel values in each cutout should represent the combined flux of all the sources it contains – primarily stars.

In SDSS u, SDSS r, 2MASS J, and WISE 3.4 µm, the median amount of flux removed from the 16 test cutouts in each band was 88%, 90%, 93%, and 86% respectively. This demonstrates that our star removal process is able to consistently remove the vast majority of the stellar flux from maps, across a wide range of wavelengths. And even in instances where our star removal is less effective than usual, a large majority of the stellar flux is still successfully removed – for 90% of the test cutouts, flux removal levels better than 81%, 71%, 84%, and 68% were achieved in each band. Moreover, note that these values are only lower limits on the fraction of the stellar flux that is removed – even if all stellar emission was successfully removed, a low level of background flux due to distant galaxies would remain in the test cutouts. This is the reason why we omitted GALEX data from this test – the contribution of distant galaxies to the background flux level in GALEX bands is large enough to no longer be negligible compared to that of foreground stars, preventing this method of testing from returning valid results.

Additionally, we wished to evaluate whether our foreground star removal method was prone to removing flux associated with compact sources within the target galaxies – HII regions being the most obvious examples. Section 6.3 compares our photometry to that of a number of independent external sources. Therein, we find no evidence of our photometry being systematically fainter than that produced by previous authors who performed by-hand identification and masking of foreground stars. We note that, as described in Sect. 3.2, we repeated our photometry with star removal disabled for maps where visual inspection revealed that the process had done more harm than good by erroneously removing flux associated with the target galaxy.

6.2. Aperture noise extrapolation validation

As described in Sect. 3.8, in cases where a map has insufficient coverage area around the target source to allow determination of aperture noise by placing full-size sky apertures in its vicinity, CAAPR employs a novel alternate approach. In these circumstances, CAAPR uses “mini-apertures” of various sizes to find the relation between aperture area and aperture noise in the map, and hence extrapolate the aperture noise associated with the aperture area of the master aperture.

To evaluate the effectiveness of this approach, we use five bands where a large proportion of the DustPedia imagery comes from wide-area surveys: GALEX NUV, SDSS r-band, 2MASS J-band, WISE 12 µm, and SPIRE 250 µm. In each of these bands, we isolated the instances where the observations provide coverage of the entire cutout. For photometry of all of these maps, the aperture noise would have proceeded according to the standard uncertainty estimation technique described in Sect. 3.7, without having to resort to mini-aperture extrapolation.

We then repeated the photometry in each of these five bands – but did so on smaller cutout maps, with diameters only 3 times larger than the major axis of the master aperture. These smaller cutouts force the aperture noise estimation process to resort to the mini-aperture extrapolation approach. Having applied the mini-aperture noise extrapolation technique to the smaller cutouts, we compared the resulting uncertainty estimates to those from our actual photometry (which were determined with full-sized randomly generated sky apertures on the standard maps).

The results of this comparison are shown in Fig. 9. As can be seen, the uncertainty estimates produced with the mini-aperture extrapolation method agree well with the standard uncertainty estimates. The average ratios between the extrapolated and standard uncertainties are 0.94 ± 0.18, 0.99 ± 0.19, 0.94 ± 0.19, 0.87 ± 0.19, and 0.90 ± 0.16, for GALEX NUV, SDSS r-band, 2MASS J-band, WISE 12 µm, and SPIRE 250 µm respectively (the ± value indicating the RMS deviation from a ratio of 1.0 in each band). It seems that the mini-aperture extrapolation method shows a tendency to result in slightly smaller uncertainty estimates, however for all bands the scatter encompasses the 1:1 relation. As such, we deem the mini-aperture extrapolation method to be suitably accurate – and certainly far superior than the alternative, of not attempting to estimate aperture noise when performing photometry on small maps.

6.3. External validation

In order to provide an objective means of validating our photometry, particularly with regards to identifying any systematic methodological issues, we compare our measurements to those from several independent external sources. Given that the scope of the DustPedia photometry is much larger than typical nearby galaxy surveys (both in terms of number of galaxies and number of bands; see Sect. 3), and given that our methodology therefore has to be much more automated, we aim to be thorough in establishing that our photometry is reliable.

When comparing DustPedia photometry to the corresponding external photometry in a given band, we consider a number of figures of merit to quantify the consistency of the two sets of fluxes. The relative scale of any offset between the sets of fluxes is quantified by the median flux ratio \( \bar{R} \), being the median value of \( S_{DP}/S_{ext} \), where \( S_{DP} \) is the DustPedia flux, and \( S_{ext} \) is the external measurement. The scatter in the relationship between two set of fluxes is described by \( \Delta R \), the median absolute deviation from \( \bar{R} \). For two sets of fluxes in good agreement with
one another, we would expect $\bar{R}$ to be smaller than $\Delta \bar{R}$, and also smaller than the calibration uncertainty of the band in question.

The other main figure of merit we consider when comparing photometry is $\chi^{[-1:1]}$, which is the fraction of galaxies for which $\chi$ lies in the range $-1 < \chi < 1$; to clarify the significance of this, we must first explain precisely what is meant by $\chi$. When comparing two sets of fluxes, we calculate $\chi$ for each galaxy, defined as:

$$\chi = \frac{S_{\text{DP}} - S_{\text{ext}}}{\sqrt{\Delta S_{\text{DP}}^2 + \Delta S_{\text{ext}}^2}},$$

where $\Delta S_{\text{DP}}$ and $\Delta S_{\text{ext}}$ are the uncertainties on the DustPedia and external fluxes respectively. The quadrature sum term $\sqrt{\Delta S_{\text{DP}}^2 + \Delta S_{\text{ext}}^2}$ represents the mutual uncertainty between a pair of measurements. Assuming well-behaved Gaussian uncertainties, then 68% of values of $\chi$ should lie in the range $-1 < \chi < 1$ – i.e., the $\chi$ distribution should have a mean of 0 and a standard deviation of 1. The figure of merit $\chi^{[-1:1]}$ describes what fraction actually lie in the range $-1 < \chi < 1$ for a given set of fluxes. $\chi^{[-1:1]}$ is a useful value; merely knowing that the median flux ratio for a set of fluxes is $\bar{R} = 1.05$, for example, tells us very little if we do not also know the uncertainties on those fluxes. If the typical flux uncertainty is 20%, then $\bar{R} = 1.05$ indicates good agreement for that set of fluxes – whereas if the typical flux uncertainty is 1%, then that same median flux ratio of $\bar{R} = 1.05$ indicates rather poor agreement. By using $\chi^{[-1:1]}$, we are able to quantify how good the agreement is for a set of fluxes with consideration given to the uncertainties on the measurements.

Throughout this section, we only consider fluxes with significance $> 5\sigma$, and which have no major flags associated with them, nor “e” flags (in the case of Planck CCS2 and IRAS SCANPI).

### 6.3.1. Herschel reference survey

The Herschel Reference Survey (HRS; Boselli et al. 2010) is a volume- and magnitude-limited survey of 323 nearby galaxies conducted by Herschel, with published multiwavelength photometry in 13 UV-submm bands. The HRS is ideally suited to serving as a reference for comparison with DustPedia; we have 288 galaxies in common, and our photometry includes all 13 bands for which they have published measurements. The HRS SPIRE photometry is presented in Ciesla et al. (2012), PACS photometry in Cortese et al. (2014), WISE photometry in Ciesla et al. (2014), and SDSS and GALEX photometry in Cortese et al. (2012).

The published GALEX UV-MIR photometry of the HRS is not corrected for Galactic extinction. However, 46% of the galaxies of the HRS are located in the Virgo Cluster (as defined by objects included in the Virgo Cluster Catalogue of Binggeli et al. 1985), which lies in a region of sky noted for having conspicuously large amounts of Galactic cirrus, despite its high galactic latitude (Auld et al. 2013; Bianchi et al. 2017). As such, to allow for a fair comparison, we extinction-correct all HRS fluxes at wavelengths $<10$ μm in the same manner as for own fluxes (as per Sect. 3.6).

Figure 10 illustrates how our photometry compares to that of the HRS – both in terms of $\chi$, and the direct relationships between the reported fluxes. The various figures of merit we consider are provided for each band in Table 2.
Fig. 10. Comparison of DustPedia CAAPR and HRS photometry. Left-hand plot for each band directly plots CAAPR and HRS photometry against each other, with 1:1 relationship indicated by dashed black line. Right-hand plot for each band shows $\chi^2$ distribution for that set of fluxes (binning dictated by the Freedman-Diaconis rule; Freedman & Diaconis 1981). Noticeable deviations for GALEX FUV, GALEX NUV, and WISE 12 $\mu$m are addressed in Sects. 6.3.1 and 6.3.1.
anything suggests a possible over-estimation of the photometric uncertainties).

In contrast to the concurrence at WISE 22 µm, the WISE 12 µm fluxes measured by CAAPR are seriously offset from those quoted by the HRS, with $R_{12} = 0.776$ – far exceeding anything that could be ascribed to simple scatter or normal methodological differences. The details of our investigation into the source of this discrepancy are given in Appendix B. But to summarise, the final result of our investigation was that we found three reasons why the WISE 12 µm fluxes measured by the HRS will differ from those measured by CAAPR: a factor 1.585 difference due to their incorrect use of the Preliminary Data Release unit calibration when they were in fact using All-Sky Data Release maps; a factor 0.929 difference due to the fact they used the DN-to-Jy given by the WISE documentation, as opposed to the zero-point magnitudes from the map headers (an apparent contradiction in the WISE data); and a factor 0.909 difference due to their smaller apertures and lack of foreground star removal. Combined, these three effects should make the HRS fluxes brighter than our own by a factor of 1.338 – i.e., an expected median offset of $R = 0.747$, which is very similar to the actual median offset of $R_{12} = 0.776$. The remaining 2.9% difference is much smaller than the WISE 12 µm calibration uncertainty of 4.6%, and well within the $\Delta R_{12} = 0.035$ scatter – suggesting we have successfully isolated all of the significant causes of disagreement between our respective photometry.

**Spitzer.** Our Spitzer 8.0 µm photometry is in good agreement with that of the HRS; the median flux ratio of $R_{8.0} = 0.985$ is well within both the 3% calibration uncertainty of the band and the $\Delta R_{8.0} = 0.032$ scatter. Although the $\chi^{[-1,1]} = 0.587$ is somewhat smaller than optimal, DustPedia and the HRS only share 65 galaxies with high-quality Spitzer 8.0 µm CAAPR fluxes, so small-number statistics prevent us from inferring too much from the $\chi$ distribution.

**SDSS.** Our SDSS photometry shows only a small offset from the HRS values, with median flux ratios $R < 1.037$ in all three bands (the HRS only report SDSS photometry for gri-bands). Nonetheless, this is larger than both $\Delta R$ and the calibration uncertainties for each band, so it is probably a real effect. However a discrepancy this small is within the realm of what can be ascribed to unremarkable methodological differences. Cortese et al. (2012) do not describe their foreground star removal process in detail, which in itself could account for an offset of this size through minor methodological differences (e.g., differences in distinguishing between bright HII regions in the target galaxies from true foreground stars). The SDSS $\chi$ distributions in Fig. 10 are well-behaved, showing few outliers, although they are a little broad (with $\chi^{[-1,1]} < 0.49$ for all bands).

**GALEX.** Our GALEX photometry shows a systematic offset from that of the HRS, with $R_{FUV} = 1.105$ and $R_{NUV} = 1.053$; both of these offsets are significant, given the bands’ respective calibration uncertainties of 4.5% and 2.7%, and scatters of $\Delta R_{FUV} = 0.075$ and $\Delta R_{NUV} = 0.047$. The cause of this offset is not immediately obvious. As with the HRS SDSS photometry, no detailed description is given of their foreground star removal process, so this is a potential methodological source of some of the discrepancy. Also, the HRS fluxes are asymptotic measurements; as this method assumes a smooth curve-of-growth, it is possible that our apertures (designed to ensure all practically-recoverable flux is recorded) encompass additional flux excluded by their technique. Cortese et al. (2012) find that their fluxes are on average ∼9% fainter than those reported by Gil de Paz et al. (2007) for the 62 galaxies they have in common (though this may in part be due to the fact that Gil de Paz et al. 2007, use the older GR2/GR3 GALEX data release) – similar to the discrepancy we find in FUV, and even larger than we find in NUV. On the other hand, we note that the HRS GALEX fluxes are systematically brighter than those reported by Bai et al. (2015; who use the current GR6/GR7 GALEX data release, and have 90 galaxies in common with the HRS), by an average of 4.0% in FUV and 2.4% in NUV. It therefore seems that the disagreement between our GALEX photometry and that of the HRS is in line with the typical variation between authors.

### 6.3.2. Planck CCS2

The wavelength coverage of Planck and Herschel overlap by design, to allow them to conduct complimentary observations; they have similar spectral response functions at 500–550 µm, and nearly identical response functions at 350 µm (Bertincourt et al. 2016). As such, our supplementary Planck CCS2 photometry is well suited to being cross-validated with our SPIRE CAAPR photometry. Plots comparing the two sets of fluxes are shown in Fig. 11, whilst figures of merit for the cross-validation are given in the upper block of Table 3. For the CCS2 photometry, we excluded any fluxes with a “c” flag (in addition to the standard exclusion of major flags).

At 350 µm the agreement between the two sets of fluxes is very good; there is no meaningful offset, whilst the $\chi$ distribution is tight and Gaussian.

For 500–550 µm we first colour-corrected the SPIRE 500 µm fluxes to allow for a valid comparison; we followed Baes et al. (2014) in multiplying the SPIRE 500 µm fluxes by a correction factor of 0.83, which assumes a 20 K modified black body spectrum with an emissivity slope of $\beta = 2.0$. The two sets of photometry seem reasonably consistent. Whilst it appears that the CCS2 fluxes might be slightly brighter (with $R_{SPIRE} = 1.05$), this is not only within the calibration uncertainty of both bands, but also within the range of deviation that could be ascribed to our colour correction. For example, were we to assume a

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**Table 2. Figures of merit for comparison of DustPedia CAAPR and HRS photometry in the 13 bands shared between the samples.**

<table>
<thead>
<tr>
<th>Band</th>
<th>$\bar{R}$</th>
<th>$\Delta R$</th>
<th>$\chi^{[-1,1]}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>GALEX FUV</td>
<td>1.105</td>
<td>0.075</td>
<td>0.390</td>
</tr>
<tr>
<td>GALEX NUV</td>
<td>1.053</td>
<td>0.047</td>
<td>0.460</td>
</tr>
<tr>
<td>SDSS g</td>
<td>1.030</td>
<td>0.049</td>
<td>0.488</td>
</tr>
<tr>
<td>SDSS r</td>
<td>1.033</td>
<td>0.032</td>
<td>0.421</td>
</tr>
<tr>
<td>SDSS i</td>
<td>1.037</td>
<td>0.040</td>
<td>0.492</td>
</tr>
<tr>
<td>Spitzer 8.0 µm</td>
<td>0.985</td>
<td>0.032</td>
<td>0.587</td>
</tr>
<tr>
<td>WISE 12 µm</td>
<td>0.776</td>
<td>0.035</td>
<td>0.159</td>
</tr>
<tr>
<td>WISE 22 µm</td>
<td>1.006</td>
<td>0.043</td>
<td>0.917</td>
</tr>
<tr>
<td>PACS 100 µm</td>
<td>1.036</td>
<td>0.053</td>
<td>0.800</td>
</tr>
<tr>
<td>PACS 160 µm</td>
<td>1.030</td>
<td>0.047</td>
<td>0.891</td>
</tr>
<tr>
<td>SPIRE 250 µm</td>
<td>1.031</td>
<td>0.046</td>
<td>0.848</td>
</tr>
<tr>
<td>SPIRE 350 µm</td>
<td>1.044</td>
<td>0.059</td>
<td>0.844</td>
</tr>
<tr>
<td>SPIRE 500 µm</td>
<td>1.061</td>
<td>0.067</td>
<td>0.801</td>
</tr>
</tbody>
</table>

**Notes.** Given figures are $\bar{R}$ (median flux ratio), $\Delta R$ (median absolute deviation in flux ratio), and $\chi^{[-1,1]}$ (fraction of values for which $-1 < \chi < 1$); all of these terms are explained in detail at the start of Sect. 6.3.
A power law slope of $\alpha = -2.0$ at 60 $\mu$m and $\alpha = -1.0$ at 100 $\mu$m would approximate a fairly unreasonable galaxy SED at those wavelengths – for example, that of a two-component modified blackbody SED with a cold dust temperature of 20 K, a warm dust temperature of 45 K, a cold-to-warm mass ratio of 100:1, and an emissivity slope of $\beta = 2$. See Smith et al. (2012), Ciesla et al. (2014), and Clark et al. (2015) for examples of typical FIR-submm SEDs of nearby galaxies.

Notes. The bands being compared are denoted by $S_1$ and $S_2$ (such that flux ratios are $S_1/S_2$); for consistency, $S_1$ represents the shortest wavelength band of each pair. Other column definitions the same as for Table 2.

6.3.3. IRAS SCANPI

It is important to ensure that our supplementary IRAS SCANPI photometry is consistent with our aperture-matched CAAPR photometry, especially at 60 $\mu$m (an important part of the spectrum for SED modelling, as previously discussed) where IRAS provides the only available photometry for $>40\%$ of the DustPedia galaxies. Plots comparing PACS 70 and 100 $\mu$m fluxes to SCANPI 60 and 100 $\mu$m fluxes are shown in Fig. 12, whilst for the cross-validation figures of merit are given in the lower block of Table 3.

To allow for a valid comparison between instruments, colour corrections for IRAS were taken from Sect. V.L.C, Table V.LC.6 of the IRAS Explanatory Supplement$^{36}$, whilst colour corrections for PACS were taken from Müller et al. (2013). Emission from galaxies in the IRAS–PACS wavelength regime tends to be due to emission from a mixture of sources (cold dust, warm dust, and small-grain emission). The full SED modelling required to describe the relative contributions of these components is beyond the scope of this work; for the purposes of colour-correction, here we simply assume a reference spectrum described by an appropriate choice of power law, of the form $S_\nu \propto \nu^\alpha$.

To compare SCANPI 60 $\mu$m and PACS 70 $\mu$m photometry, we colour-corrected both sets of fluxes assuming a power law index of $\alpha = -2.0$ as the reference spectrum$^{36}$; this

Notes. The upper block of figures are for the direct comparison between the DustPedia CAAPR and 2MASS-LGA photometry; the lower block of figures (in italic) compare the 2MASS-LGA photometry to CAAPR photometry repeated using the 2MASS-LGA apertures. Column definitions the same as for Table 2.

6.3.4. 2MASS large Galaxy Atlas

The only instrument for which HRS, Planck CCS2 and IRAS SCANPI photometry provide no independent external validation for our aperture-matched CAAPR photometry is 2MASS. To remedy this, we refer to the 2MASS Large Galaxy Atlas (2MASS-LGA: Jarrett et al. 2003). The 2MASS-LGA has 129 galaxies in common with DustPedia, spread across the sky (including in the plane of the Milky Way). We specifically compare to the 2MASS-LGA large-aperture Kron photometry, as it is the most similar of their photometric measures to our own. Plots comparing the CAAPR and 2MASS-LGA photometry for our common sources are shown in Fig. 13, and the validation figures of merit are given in the upper block of Table 4.

See also Casey (2012) for discussion of $\alpha = -2.0$ being a typical power law in the $~60 \mu$m range.
In all three \( JHK_s \)-bands, the 2MASS-LGA photometry compares poorly with our own. Our fluxes are systematically brighter by an average of 19\%, and the \( \chi \) distributions are broad, non-Gaussian, and asymmetrical in all bands. However, the photometric apertures used by the 2MASS-LGA are significantly smaller than our own; on average, our apertures encompass 8 times more area. Visual inspection of the 2MASS-LGA apertures reveals that they often do not encompass all of the stellar emission of the target galaxies – especially when compared to the deeper WISE 3.4 \( \mu m \) data, which shows an appreciable fraction of some galaxies’ NIR flux can extend beyond the 2MASS-LGA apertures. This strongly suggests that the reason we record systematically brighter \( JHK_s \)-band photometry was that we recovered flux excluded by the much smaller apertures of the 2MASS-LGA.

To test this hypothesis, we repeated our photometry, but instead used the apertures employed by the 2MASS-LGA; the resulting comparison plots are shown in Appendix C, with validation figures of merit in the lower block of Table 4. As can be seen, this photometry is in near-perfect agreement with that of the 2MASS-LGA. The offsets have been eliminated (to the <1\% level), and the \( \chi \) distributions are Gaussian and narrow (with \( \chi^{[-1,1]} > 0.7 \)) in all bands. We can therefore confidently state that difference in aperture sizes is the cause of the disagreement between our fluxes. Given that our larger apertures contain flux missed by the 2MASS-LGA (as we as being matched to the rest of our photometry), we deem our own photometry to be preferable. The high degree of scatter seen in Fig. 13 is an unavoidable consequence of our larger apertures, driving up the aperture noise associated with our measurements (an effect to which 2MASS is particularly vulnerable to due to sky brightness; see Sect. 2.2.4).

For a number of our galaxies, visual inspection of the multiwavelength SEDs indicated that the \( H \)-band flux can appear conspicuously offset above or below the \( J \)-band and \( K_s \)-band fluxes\footnote{NB, here we are once again referring to our standard aperture-matched CAAPR photometry, not the repeated 2MASS photometry employed in the previous paragraph.}. Our investigation of this is detailed in Appendix D; to summarise, we find that our \( H \)-band photometry does not exhibit any systematic offset from that of the 2MASS-LGA, but there is evidence of non-Gaussianity in our \( H \)-band photometric uncertainties.

### Table 5. Figures of merit for internal comparison of DustPedia photometry in similar bands.

<table>
<thead>
<tr>
<th>( S_1 )</th>
<th>( S_2 )</th>
<th>( \bar{R} )</th>
<th>( \Delta R )</th>
<th>( \chi^{[-1,1]} )</th>
</tr>
</thead>
<tbody>
<tr>
<td>WISE 3.4 ( \mu m )</td>
<td>Spitzer 3.6 ( \mu m )</td>
<td>1.027</td>
<td>0.044</td>
<td>0.636</td>
</tr>
<tr>
<td>Spitzer 4.5 ( \mu m )</td>
<td>WISE 4.6 ( \mu m )</td>
<td>0.901</td>
<td>0.041</td>
<td>0.308</td>
</tr>
<tr>
<td>WISE 22 ( \mu m )</td>
<td>Spitzer 24 ( \mu m )</td>
<td>1.040</td>
<td>0.064</td>
<td>0.686</td>
</tr>
<tr>
<td>PACS 70 ( \mu m )</td>
<td>Spitzer 70 ( \mu m )</td>
<td>1.083</td>
<td>0.106</td>
<td>0.563</td>
</tr>
<tr>
<td>PACS 160 ( \mu m )</td>
<td>Spitzer 160 ( \mu m )</td>
<td>0.976</td>
<td>0.091</td>
<td>0.838</td>
</tr>
</tbody>
</table>

**Notes.** Column definitions the same as for Tables 2 and 3. For consistency, \( S_1 \) represents the shortest wavelength band of each pair.

### 6.4. Internal validation

Thanks to wavelength overlap between the various instruments used for our photometry, there is good scope for internal validation, comparing CAAPR photometry in bands that are similar to one another. Figure 14 shows plots comparing the overlapping

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\( JHK_s \)-bands

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\( JHK_s \)-bands

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\( JHK_s \)-bands

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\( JHK_s \)-bands

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\( JHK_s \)-bands

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\( JHK_s \)-bands

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\( JHK_s \)-bands
bands; the cross-validation figures of merit are listed in Table 5. We employ the same figures of merit we did for external validation; for instance, when comparing two bands, $S_1$ and $S_2$, the median flux ratio between the two is denoted by $R_{S_1|S_2}$.

6.4.1. WISE-Spitzer

The WISE 3.4 μm and Spitzer 3.6 μm bands have extremely similar relative system responses, as do WISE 4.6 μm and Spitzer 4.5 μm (see Fig. 1 of Jarrett et al. 2011). To allow comparison between both pairs of bands, we colour correct all fluxes assuming a Rayleigh-Jeans reference spectrum; colour corrections for WISE were taken from Table 1 of Wright et al. (2012), whilst corrections for Spitzer were taken from Table 4.3, Sect. 4.4 of the IRAC Instrument Handbook. These corrections are small, being <1.2% in all cases.

We find our WISE 3.4 μm and Spitzer 3.6 μm fluxes to be in excellent agreement, with the median offset smaller than either bands’ calibration uncertainty, and a good $\chi^2_{3,4,5}$ value of 0.636. This is in keeping with the findings of Jarrett et al. (2012) and Papovich et al. (2016), who report that the two bands exhibit no noticeable offset.

However, we do find a definite offset between our WISE 4.6 μm and Spitzer 4.5 μm photometry, with $R_{4.6|4.5} = 0.901$. This is twice the size of the 4.5% mutual calibration uncertainty between the bands (i.e., $\sqrt{3\%^2 + 3.4\%^2}$). Previous authors have reported varying degrees of disagreement between these bands, with Jarrett et al. (2012) finding offsets in the range $0.89 < R_{4.6|4.5} < 1.02$ for a small sample of extended nearby galaxies, whilst Papovich et al. (2016) find $R_{4.6|4.5} = 0.98$ for a large sample of high redshift galaxies. Jarrett et al. (2012) point out that the expected flux offset depends on the nature of the target galaxy, from 0.94 for ellipticals, to 0.93–0.99 for standard spirals, to 1.09 for starbursts; given that it is beyond the scope of this work to perform the full SED-fitting necessary to provide source-specific reference spectra, such differences could account for a large proportion of the offset we find. As the matter stands, users of this photometry should be aware of the potential disagreement here; we are hopeful that the full SED-fitting to be performed in future DustPedia papers will allow for source-specific colour corrections that negate the issue. And notwithstanding the offset, the $\chi$ distribution for these bands is excellent.

Given that the part of the spectrum sampled by the WISE 22 μm and Spitzer 24 μm bands has contributions from stellar emission, hot dust emission, and PAH-like emission, in proportions that vary greatly between sources, we are unable to apply any colour-corrections in the absence of full SED-fitting. As per the corrections given in Sect. 4.3.5 of the MIPS Instrument Handbook, adjusting Spitzer 24 μm fluxes to the reference spectrum employed by WISE requires a correction of only 0.1% (thanks to the close similarity at 24 μm between the 10 K blackbody reference spectrum used by WISE and the perfect Rayleigh-Jeans reference spectrum used by Spitzer-MIPS). The two sets of fluxes are entirely compatible; the median flux ratio of $R_{22|24} = 1.04$ is smaller than the WISE 22 μm calibration uncertainty, and the minimal scatter yields an almost ideal $\chi^2_{22|24} = 0.686$. This agreement between the two sets of fluxes is good enough that it essentially rules out the risk of egregious differences arising when source-by-source colour corrections are applied during full SED-fitting.

6.4.2. PACS-Spitzer

PACS and Spitzer have two almost-identical bands in common, at 70 and 160 μm. To colour-correct the 70 μm fluxes for direct comparison, we follow the same procedure as in Sect. 6.3.3, assuming an $S_e \propto \nu^\beta$ reference spectrum. At 160 μm we assume a 20 K blackbody reference spectrum. In both cases we adjust the PACS fluxes as per Müller et al. (2013; including the translation...
factor to account for the difference in central wavelengths between the PACS and Spitzer bands), and adjust the Spitzer fluxes as per Sect. 4.3.5 of the MIPS Instrument Handbook.

At both 70 and 160 μm, the PACS and Spitzer fluxes are in good agreement, with average offsets smaller than the Spitzer calibration uncertainties in both cases (and also smaller than the PACS calibration uncertainty at 160 μm). And both χ distributions appear acceptable, given the small number of sources involved.

Whilst it is satisfying to find such good agreement between PACS and Spitzer at 160 μm, it is worth noting that there are only 5 sources without major flags where Spitzer 160 μm coverage is available when PACS 160 μm coverage is not (and only 3 of those are >3σ). As such, there are almost no instances at 160 μm where users will be unable to use superior PACS fluxes instead of those from Spitzer.

7. Data products

7.1. Photometry

The DustPedia photometry tables can be accessed from the DustPedia database42, and from the VizieR catalogue service43. Separate tables are provided for our aperture-matched CAAPR photometry, the supplementary IRAS SCANPI photometry, and the supplementary Planck CCS2 photometry. Each table has a row for every DustPedia galaxy. Each band has 3 columns; one giving the flux, one giving the uncertainty on the flux (with calibration uncertainty included, as per Sect. 3.7), and one stating any flags associated with that flux. If the flux and error are blank for a given source in a particular band, this indicates that there is no data for the target in that band, or that there was insufficient data to make a measurement (i.e., that the map in question was so small that no pixels were located in the background annulus). Whenever a measurement could be made, the resulting value is given, regardless of whether or not the measurement represents a detection. For maps where a flux could be measured, but where too little sky area was available to determine the aperture noise (even for CAAPR’s mini-aperture extrapolation), then we quote a negative uncertainty; this indicates the uncertainty only incorporates the contribution of the instrument’s calibration uncertainty.

Extended-source corrections have been applied to the WISE and Spitzer-IRAC photometry, as per Sects. 2.2.5 and 2.2.6.

In total, we present 21,857 photometric measurements; 18,254 fluxes from our aperture-matched CAAPR photometry, along with 1079 Planck CCS2 fluxes, and 2533 IRAS SCANPI fluxes. A typical DustPedia source has photometry in 25 bands; Fig. 15 shows the example UV-mm SED of NGC 3683, a DustPedia galaxy with rich photometric coverage, having fluxes in 32 bands.

Our photometry demonstrates that the DustPedia sample spans a wide range of luminosities; amongst the DustPedia sources with >5σ detections, SPIRE 250 μm luminosity ranges from 6.1–10.6 log10 L⊙, whilst WISE 3.4 μm luminosity (a standard proxy for stellar mass) ranges from 5.0–10.4 log10 L⊙. We span a similarly wide range of colour, with NUV-r colour ranging from 0.8–6.3 mag, and FUV-Ks colour (a good proxy for morphology, see Gil de Paz et al. 2007) ranging from 0.6–8.8 mag.

7.2. Sample catalogue

Also provided is the sample catalogue, providing the key properties of each DustPedia galaxy, such as position, morphology, inclination, optical D25, etc. This catalogue provides a number of distance measures; because of the difficulty in establishing reliable distances to nearby galaxies, we have established an order of preference.

Our preferred distance estimates are the redshift-independent distances provided by the HyperLEDA database, as these have been homogenised to account for methodological differences between references. For galaxies without redshift-independent HyperLEDA distances, our next preferred values are the redshift-independent distances provided by the

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42 http://dustpedia.astro.noa.gr/Photometry
43 Available at the CDS (see title footnote).
NASA/IPAC Extragalactic Database (NED); whilst these are not homogenised like the HyperLeda values, we take the quoted average value for each galaxy to minimise bias. If neither HyperLeda nor NED redshift-independent distances are available for a source, we use the flow-corrected redshift-derived values provided by NED; these distances were calculated assuming a Hubble constant of \( H_0 = 73.24 \text{ km s}^{-1} \text{ Mpc}^{-1} \) (Riess et al. 2016), and have been corrected for bulk deviation from Hubble flow arising from the influence of the Virgo Cluster, the Great Attractor, and the Shapley Supercluster (Mould et al. 2000). For each galaxy in the DustPedia database we list each distance measure available, along with the preferred value as per the hierarchy described here.

7.3. FITS images

The standardised DustPedia imagery (as described in Sect. 2) can be accessed from the DustPedia database, in the form of FITS images. The database may be queried according to recessional velocity, angular size, and Hubble type. Data for individual galaxies may also be searched for by name (note that, at present, queries by name must use the name by which each galaxy is listed in the sample catalogue). The maps may also be retrieved programmatically via URL. For example, the SPIRE 250 \( \mu \text{m} \) map of NGC 0891 can be directly downloaded with the following URL: http://dustpedia.astro.noa.gr/Data/GetImage?imageName=NGC0891_SPIRE_250.fits&instrument=SPIRE. Other maps can be retrieved in the same manner, by changing the galaxy, instrument, and band names given in the URL.

In total, the imagery presented in this work represents 21 724 FITS images (not counting error maps), including the 3297 non-photometric DSS maps. A typical DustPedia source has maps providing coverage of 25 bands.

8. Summary

We have presented the imagery and photometry of the DustPedia sample, covering 875 extended nearby galaxies in 42 UV-radio bands – every nearby extended galaxy that was observed by the Herschel Space Observatory. The centrepiece of the dataset we present is our consistent multiwavelength aperture-matched photometry, encompassing 27 bands (from GALEX, SDSS, 2MASS, WISE, Spitzer, and Herschel). In combination with supplementary IRAS and Planck photometry, this represents 21 857 fluxes; on average, each DustPedia galaxy possesses photometry in 25 bands. Additionally, we have produced imagery spanning 38 bands, including custom reductions of Herschel observations, along with standardised preparations of archival UV-radio data.

To perform our aperture-matched photometry, we developed the Comprehensive & Adaptable Aperture Photometry Routine (CAAPR), which is designed to determine fluxes and uncertainties that can be directly compared across a wide range of wavelengths and data types. CAAPR features a novel technique of determining aperture noise, even in maps with very little background coverage around the target source; this allows us to determine uncertainties consistently across the wide range of observations we employ.

We perform extensive validation of our photometry, comparing overlapping bands internally, and comparing to independent external sources. We find that our photometry is consistently reliable, with no excessive discrepancies.

The resulting database is being made publicly available, for the benefit of the astronomical community.

Acknowledgements. The authors kindly thank the anonymous referee. This work was improved by the referee’s attention to detail and constructive comments. The DustPedia project is funded as a European Revi by the European Union, the Comprehensive & Adaptable Aperture Photometry Routine (CAAPR), which is designed to determine fluxes and uncertainties that can be directly compared across a wide range of wavelengths and data types. CAAPR features a novel technique of determining aperture noise, even in maps with very little background coverage around the target source; this allows us to determine uncertainties consistently across the wide range of observations we employ.

We perform extensive validation of our photometry, comparing overlapping bands internally, and comparing to independent external sources. We find that our photometry is consistently reliable, with no excessive discrepancies.

The resulting database is being made publicly available, for the benefit of the astronomical community.
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Storn, R., & Price, K. 1997, J. Global Optim., 11, 341

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Smith, M. W. L. 2012, Ph.D. Thesis, Cardiff University, School of Physics and Astronomy, Cardiff
Storn, R., & Price, K. 1997, J. Global Optim., 11, 341

Sérsic, J. L. 1963, Boletin de la Asociacion Argentina de Astronomia La Plata Argentina, 6, 41
Smith, M. W. L. 2012, Ph.D. Thesis, Cardiff University, School of Physics and Astronomy, Cardiff
Storn, R., & Price, K. 1997, J. Global Optim., 11, 341
Appendix A: FITS image data unit conversions

Table A.1 provides information regarding the pixel units of the original maps for our ancillary data, as taken from the official archive for each facility, and the relevant quantities we employed to convert these pixel units to our consistent Jy pix$^{-1}$ units. See Sect. 2.2 for the full details of each ancillary data facility.

The $m_{ZP}$ column of Table A.1 provides the zero-point magnitude used to convert the native map units $S$ to VCG magnitudes $m$, as per $m = m_{ZP} - 2.5 \log_{10}(S)$. In the case of 2MASS, each tile has its own independently-calibrated zero-point magnitude, provided in the FITS header; as such no $m_{ZP}$ values can be listed for the 2MASS bands.

The $m_{ZP}$ values listed are the values provided by each facility, and as such give magnitudes in whatever native magnitude system each facility employs. Both 2MASS and WISE use Vega magnitudes, whilst the SDSS uses SDSS magnitudes (which are similar to AB magnitudes, but with small offsets in $u$ and $z$). The $\Delta m_{AB}$ column of Table A.1 gives the offset between these native magnitudes for each band, and the corresponding AB magnitudes, such that $m_{AB} = m + \Delta m_{AB}$; GALEX employs AB magnitudes by default, so there is no $\Delta m_{AB}$ offset.

Table A.1. Description of map unit conversions for various UV-MIR ancillary data for which pixel units of input archival maps were not already expressed in terms of Jy.

<table>
<thead>
<tr>
<th>Facility</th>
<th>Band</th>
<th>Native units</th>
<th>$m_{ZP}$</th>
<th>$\Delta m_{AB}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>GALEX</td>
<td>FUV</td>
<td>counts s$^{-1}$</td>
<td>18.82</td>
<td>-</td>
</tr>
<tr>
<td>GALEX</td>
<td>NUV</td>
<td>counts s$^{-1}$</td>
<td>20.08</td>
<td>-</td>
</tr>
<tr>
<td>SDSS</td>
<td>$u$</td>
<td>nanomaggies</td>
<td>22.50</td>
<td>-0.04</td>
</tr>
<tr>
<td>SDSS</td>
<td>$g$</td>
<td>nanomaggies</td>
<td>22.50</td>
<td>+0.00</td>
</tr>
<tr>
<td>SDSS</td>
<td>$r$</td>
<td>nanomaggies</td>
<td>22.50</td>
<td>+0.00</td>
</tr>
<tr>
<td>SDSS</td>
<td>$i$</td>
<td>nanomaggies</td>
<td>22.50</td>
<td>+0.00</td>
</tr>
<tr>
<td>SDSS</td>
<td>$z$</td>
<td>nanomaggies</td>
<td>22.50</td>
<td>+0.02</td>
</tr>
<tr>
<td>2MASS</td>
<td>$J$</td>
<td>DN</td>
<td>-</td>
<td>+1.39</td>
</tr>
<tr>
<td>2MASS</td>
<td>$H$</td>
<td>DN</td>
<td>-</td>
<td>+1.85</td>
</tr>
<tr>
<td>2MASS</td>
<td>$K_S$</td>
<td>DN</td>
<td>-</td>
<td>+2.669</td>
</tr>
<tr>
<td>WISE</td>
<td>3.4 $\mu$m</td>
<td>DN</td>
<td>20.5</td>
<td>+3.339</td>
</tr>
<tr>
<td>WISE</td>
<td>4.6 $\mu$m</td>
<td>DN</td>
<td>19.5</td>
<td>+3.339</td>
</tr>
<tr>
<td>WISE</td>
<td>12 $\mu$m</td>
<td>DN</td>
<td>18.0</td>
<td>+5.174</td>
</tr>
<tr>
<td>WISE</td>
<td>22 $\mu$m</td>
<td>DN</td>
<td>30.0</td>
<td>+6.620</td>
</tr>
</tbody>
</table>

Appendix B: External validation – investigating severe disagreement with HRS WISE 12 $\mu$m photometry

As described in Sect. 6.3.1, our photometry significantly differs from that of the HRS at 12 $\mu$m, with our fluxes being fainter by an average factor of 0.776. Here we discuss in detail how we determined the sources of this disagreement.

After investigation, it transpired that the HRS converted their WISE measurements from “DN” (the WISE term for instrumental pixel units) to Jy using the conversion provided in the WISE Preliminary Data Release Explanatory Supplement\textsuperscript{45} (Ciesla et al. 2014), despite the fact the maps they used came from the WISE All-Sky Data Release (L. Ciesla, priv. comm.). The 12 $\mu$m maps of the All-Sky Data Release have a very different zero-point from that of the Preliminary Data Release, causing the HRS to overestimate their 12 $\mu$m fluxes by a factor of 1.585 (the 12 $\mu$m zero-point of the AllWISE Data Release, used by us, is the same as that of the All-Sky Data Release).

Our investigation also revealed an apparent inconsistency in the WISE All-Sky and AllWISE data releases. As described in Sect. 2.2.5, we converted our WISE imagery from DN to Vega magnitudes using the zero-point magnitudes provided in the FITS headers, then to AB magnitudes using the offsets listed in the Table 1, Sect. IV.3.a AllWISE Explanatory Supplement, and thereby to Jy (we use this conversion method, as our preference is for using the photometric calibration information provided with the data itself). However, if we instead use the DN-to-Jy conversion factor listed in Table 1, Sect. IV.3.a of the AllWISE Explanatory Supplement (mirroring the approach of the HRS), the resulting final map units are not the same – map units produced by using the DN-to-Jy conversion from the documentation differ by a factor of 0.929 from the map units produced by using the zero-point magnitudes in the FITS headers.

Another potential source of discrepancy between the CAAPR and HRS 12 $\mu$m fluxes is the fact that the HRS do not perform any foreground star removal, and they use much smaller photometric apertures than our own (theirs being designed to contain only the visible MIR emission). To test the impact of these differences, we repeated our photometry but using the same apertures as Ciesla et al. (2014), with foreground star removal disabled – the resulting fluxes differed from our actual fluxes by a factor of 0.909 (for >5σ measurements).

The combined effect of these three differences is that we should expect our WISE 12 $\mu$m photometry to be fainter than the WISE 12 $\mu$m photometry of the HRS by a factor of 0.747 – very close to the actual factor of 0.776 that we encounter.

Appendix C: External validation – repeating the 2MASS LGA photometry

In Sect. 6.3.4, we repeat our 2MASS photometry, in order to compare it with the 2MASS-LGA. We use CAAPR to perform this repeat photometry, but instead of our own apertures we employ the apertures used by the 2MASS-LGA. This is to establish whether the fact that our apertures are consistently larger than those of the 2MASS-LGA is the reason why our photometry is systematically brighter. The plots of the results of this comparison are shown in Fig. C.1. The findings from this comparison are discussed in Sect. 6.3.4.

Appendix D: External validation – assessing the reliability of 2MASS $H$-band photometry

When visually inspecting the multiwavelength SEDs of our galaxies, we found that our $H$-band photometry sometimes appears conspicuously offset above or below the $J$-band and $K_S$-band, manifesting as a “bump” or “dip” in the SED at $H$-band; see Fig. D.1 for an example.

To establish whether this was indicative of some sort of systematic issue, we calculated the $[J/H]$ and $[H/K_S]$ colours of our galaxies using our aperture-matched CAAPR photometry. If the $H$-band fluxes were suffering from a systematic artefact, we would expect the colours calculated with our CAAPR photometry to be systematically offset from the equivalent colours calculated with the 2MASS LGA photometry. However, we find that the CAAPR colours show no systematic difference from the 2MASS LGA colours; an offset can be ruled out to the <1.4% level. What we do find is that the CAAPR photometry exhibits a

45 http://wise2.ipac.caltech.edu/docs/release/prelim/expsup/wise_prelrel_toC. J. R. Clark et al.: DustPedia photometry and imagery
Fig. C.1. Comparison of CAAPR photometry performed using the apertures employed by the 2MASS-LGA, to photometry the 2MASS-LGA themselves report. Plots as per Fig. 10.

much larger scatter in colour than the 2MASS LGA photometry. Whilst we should expect CAAPR colours to exhibit more scatter (due to the larger uncertainties of our aperture-matched photometry), the degree to which the scatter is larger for the CAAPR colours is greater than we would predict assuming the photometric uncertainties to be Gaussian – specifically, 65% greater for $[J/H]$ and 26% greater for $[H/K_S]$.

For comparison, we also calculated the $[J/K_S]$ colours for both the CAAPR and 2MASS-LGA photometry. Once again, we found greater scatter with the CAAPR photometry than with that of the 2MASS-LGA – however, the increase was within the bounds of what would be expected given the larger CAAPR uncertainties (assuming they behave in a Gaussian manner). We can therefore infer that the conspicuous $H$-band offsets found for some sources are due to non-Gaussianity in our $H$-band uncertainties. This results in there being more outliers amongst our $H$-band photometry than we would otherwise expect (such as that shown in Fig. D.1). Users who are concerned by the this non-Gaussianity have the option of not using the $H$-band photometry – the $J$-band $K_S$-band fluxes do not seem to suffer from a similar problem, given that $[J/K_S]$ colours behave as expected.

The non-Gaussianity of the $H$-band uncertainties could arise from the complex nature of the sky brightness suffered by 2MASS $H$-band observations. The arcminute scales of the $H$-band sky brightness make it particularly problematic for nearby-galaxy photometry.

Fig. D.1. UV-mm SED of NGC 3631, showing an especially pronounced example of a “bump” in our 2MASS $H$-band photometry. 2MASS photometry is plotted in dark blue, with $H$-band hence being the central of the three 2MASS fluxes plotted (at 1.65 $\mu$m).