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Keywords Force measurements; impact; bird strike; numerical simulation

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1 **Development and validation of a set-up to measure the transferred multi-**
2 **axial impact momentum of a bird strike on a booster vane**

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7 **Abstract**

8 Reaction force has always been one of the main characterization parameters for impact events.

9 Today, a set of force transducers are a common and valuable tool to measure reaction forces. But the
10 force signals are often influenced by vibrations of the supporting structures. Many other attempts
11 were already taken in the past to use other methods to measure force, such as ballistic pendulums,
12 Hopkinson bars, etc., all having their advantages and disadvantages. In this work, a multi-axial force
13 measurement tool is developed to serve in a test campaign of bird strike experiments on booster
14 vanes. The idea is to give some well-chosen mass three rotational degrees of freedom and acquire
15 the transferred rotational momentum from an optical measurement, which is a direct measure for
16 the impact force. The tool is validated experimentally and numerically using a simplified steel vane.

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18 1) Introduction

19 Certification by analysis is a hot topic these days. But, a lot of research is still required to be able to
20 prove that numerical methods are fully capable for simulating bird strike. The work described in this
21 paper is part of the European FP7 project *E-Break or Engine BREAKthrough Components and*
22 *Subsystems*, where the key jet engine subsystem technologies are further developed to incorporate
23 in ultra-high overall pressure ratio (OPR) and high bypass ratio (BPR) engines. In this project, a small

24 task is devoted to the development of a numerical model that is able to validate the design rules of
25 the booster vane in terms of bird strike robustness and investigate the possibilities of Variable Stator
26 Vane (VSV) systems.

27 The validation of the numerical models for bird strike requires quantitative measurements. Strain
28 gauges on one hand can tell something about local deformations at discrete points on a structure.
29 The measurement of residual energy after impact [ref eigen werk] and reaction forces on the other
30 hand are valuable parameters that give an idea of the global performance. Optical measurements
31 can also provide full field displacement and strain fields. The optical view however is often disturbed
32 in bird strike experiments, and a stereo set-up of high-speed cameras dedicated to the measurement
33 would be required. The measurement of reaction forces therefore remains one of the primary
34 parameters for characterizing the impact event. Several techniques were already successfully used,
35 from ballistic pendulums to Hopkinson bars and load cells and methods in between.

36 The oldest technique is the ballistic pendulum. The original idea dates back from the reference work
37 of Robins in 1742 [1], where it was used to measure the momentum of a bullet. An application of the
38 pendulum in bird strike research can be found in the work of Bertke et al. [2], where a 5 wire
39 pendulum was used to measure the total transferred momentum of a bird strike on titanium blades.
40 They calculated the transferred momentum from the chord length and the oscillation period after
41 impact.

42 Hopkinson introduced a first version of the Hopkinson bar in 1914 [3], which was basically an
43 advanced version of the ballistic pendulum. Hopkinson proposed a co-axial system of two bars,
44 where the second bar is suspended and able to trap a part of the momentum depending on its
45 length. The strain waves in the first bar however can, in the ideal case, be directly related to the
46 impact force, as was tried in the reference works on bird strike from Barber et al. [4] and Wilbeck [5],
47 in which forces were measured of bird strike on flat and inclined surfaces using a Hopkinson bar set-
48 up. But they had to integrate the momentum signals. Because of the high frequencies that were

49 dispersed due to the large diameters of the bars, exact force time signals could not be obtained. A
50 more recent attempt was taken in the work of Seidt [6].

51 In the work of Allcock [7], the targets were attached to a set of calibrated beams. The deflection of
52 the target was measured, from which the impact force was derived. In more recent work [ref eigen
53 werk], the targets were able to move in the direction of impact, from which the impact force could
54 be derived directly.

55 To test bigger and complex full scale structures such as leading edge wings, flaps or windshields and
56 to acquire force time signals, a set of load cells or instrumented links are often used to measure the
57 reaction force at discrete points [8-16]. The problem with load cells is that the force signal is often
58 influenced by vibrations of the supporting structure. Numerical simulations are capable of
59 incorporating a part of the boundary conditions, but the interpretation of the signals is nevertheless
60 not straightforward.

61 Prior to the booster vane experiments that were to be executed in the course of this project, the
62 purpose was to develop a tool able to measure reaction forces in multiple directions. This is achieved
63 by allowing movement of the set-up, which has the advantage to decouple the experiment from the
64 environment to some extent and to guarantee safety. More specific, three rotational degrees of
65 freedom are given to a well-chosen mass to acquire the transferred rotational momentum, which is
66 an idea that originates from the work of Premont et al. [17] and Steinhagen et al. [18]. Premont and
67 Steinhagen mounted vanes onto a rigid object that is able to pivot around one point, and measured
68 the transferred momentum around three axes using accelerometers.

69 Contrary to the work of Premont and Steinhagen, the moment arm from the point of rotation to the
70 impact location is much shorter, reducing the influence of Eigen frequencies of the force
71 measurement tool. Also, the momentum in this work is derived optically from the images of one high
72 speed camera (HSC). The HSC is already necessary to acquire the horizontal offset of the bird. A so-
73 called cone structure was designed onto which multiple vane configurations can be mounted. To

74 verify the set-up and to have an intermediate step between the initial calibration experiments [Ref
75 eigen werk] and the booster vane experiments, a rigid steel vane was used as target object. This
76 paper will introduce the cone as the tool to allow the rotational movement. The different steps
77 necessary to derive the momentum of the cone will be explained in detail. Finally, some results will
78 be shown as well as a comparison with a numerical model, using SPH modelling for the bird.

79 In the next section, the test set-up to launch the birds and the steel vane will be introduced. In
80 section 3, the main principle of the force measurement with the cone will be explained. Section 4
81 contains the actual derivation of the rotational momentum. The next sections contain the results of
82 some experiments, a comparison with simulations and finally a conclusion.

83 2) Test set-up

84 a. Ugent bird strike test set-up

85 The experiments were performed on the Ghent University bird strike set-up (Figure 1). The set-up is
86 capable of shooting birds up to 42 kJ. Birds can be launched with a weight of maximum 4 lb
87 (according to the regulations [19]) at speeds up to 250 m/s. At the beginning of each experiment, a
88 projectile called a sabot is filled with foam in accordance to the desired shape, after which gelatine is
89 moulded into the acquired foam shape. The sabot is mounted in front of a pressure vessel and
90 released at the required pressure. After the release trigger, the sabot launches through a barrel and
91 strips off from the bird in the stripper chamber using a cone shaped stripper, after which the stripped
92 bird flies into the test chamber and impacts on the required target. Before each experiment, the test
93 chamber is evacuated up to 0.2 bar absolute pressure to be able to perform precise velocity
94 measurements.



Figure 1: Ghent University bird strike test set-up

95

96

b. Target: steel vane

97

The set-up to measure the transferred rotational momentum will be tested and validated using a

98

simplified steel vane (Figure 2). Tests on the actual vanes cannot be disclosed. The vane consists of a

99

V-shaped steel bar welded to a plate. The holes in the plate will be used to mount the vane to the

100

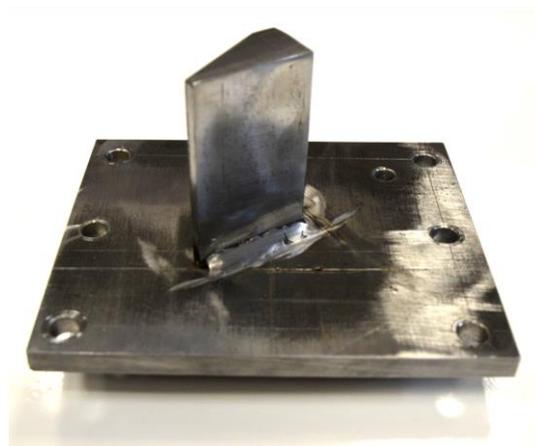
cone structure. Making a (construction steel) vane with similar behaviour (in terms of elastic energy,

101

deformation and thickness) as the Titanium vanes is almost impossible due to the difference in yield

102

strength and stiffness.



103

104

Figure 2: Simplified steel vane

105 The steel vane will serve as an intermediate step between the initial calibration experiments on rigid
106 targets [Ref eigen werk] and the experiments on the booster vanes because it will combine a
107 significant amount of change of momentum with the splitting of the bird and the surface over which
108 mass is deflected is significantly shorter, as will also be the case in the booster vane experiments.

109 3) Principle rotational momentum measurement

110 In order to acquire reaction forces in multiple directions, multiple degrees of freedom (DOF) are
111 necessary. It proved to be practically difficult to allow translational movement in multiple directions
112 for impact measurements, which is why the idea of Premont et al. [17] and Steinhagen et al. [18] to
113 use rotational DOFs was further investigated. Practically, the rotational DOFs are realized by a cone
114 structure (Figure 3 shows the first version). The ball joint at the tip of the cone is the centre of
115 rotation. The vane fixtures should be mounted on top of the flange attached to the cone. The
116 rectangular shaped protrusions on each side of the flange are used for the optical measurements.

117 At impact, the cone structure with the vane fixture starts rotating around the ball joint. From this
118 movement the momentum can be calculated. The weight and rotational inertia was designed in such
119 a way that the displacement of the flange during impact is in the order of millimetres (1-5 mm),
120 which minimizes the influence of the set-up on the experiment. This way the experiment is also
121 decoupled from the environment as much as possible (the structure to hold the ball joint, which is
122 assumed rigid sees no moment) and is therefore less dependent on supporting structures. The
123 distance from the point of impact to the rotation point is also much larger than the radius of the ball
124 joint, which decreases the possible influence of friction forces on the momentum measurement.



125

126

Figure 3: the cone structure

127 Newton's second law for rotational systems is the following:

128

$$\tau = I\alpha$$

129 With τ , the torque in Nm, α , the rotational acceleration in rad/s² and I , the inertia tensor in kg.m²,

130 which can be obtained from CAD software (containing the parts with the actual dimensions):

131

$$I = \begin{bmatrix} I_{xx} & I_{xy} & I_{xz} \\ I_{xy} & I_{yy} & I_{yz} \\ I_{xz} & I_{yz} & I_{zz} \end{bmatrix}$$

132 The angular momentum L can be calculated by integrating the torque:

133

$$L = \int \tau dt$$

134 in kg.m²/s or also Nms.

135 The idea is to have a tool that does not influence the experiment. This means that, if the vane would

136 be mounted on a rigid surface, the same impact forces should be measured. Applying this to the

137 equations above, the angular momentum from the torque in an impact event with rigid boundary

138 conditions obtained in equation 3 should give the same results as the angular momentum obtained

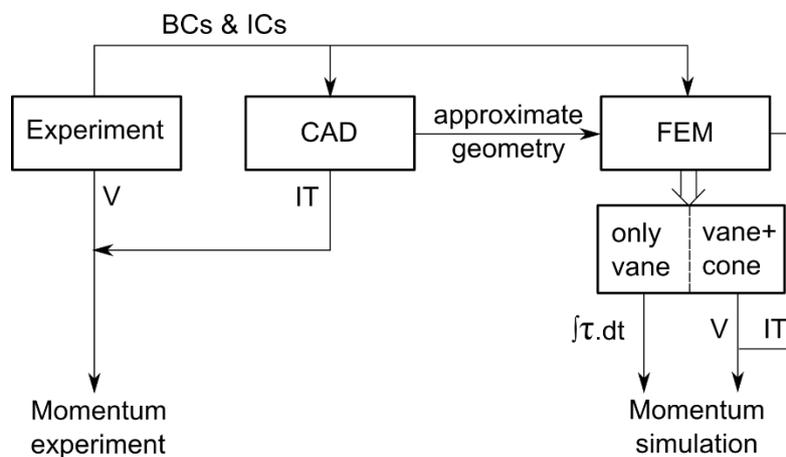
139 from the rotational accelerations in equation 1:

140

$$L = \int I\alpha dt = I \int \alpha dt = I\omega$$

141 This is only valid when I is not dependent on time, or also, that the entire structure is rigid and fixed.
 142 This is why the displacement during impact should be limited. For deforming objects such as a
 143 booster vane, I is not exactly constant. But, simulations showed that the average transferred
 144 momentum is equal in both situations with and without cone, as well as the amplitude of the
 145 oscillations superimposed on this average value. Small delays in the momentum signals showed to be
 146 the biggest influence of the cone. The influence of the cone will be further discussed in the results
 147 section.

148 The fact that a well-designed inertial tensor does not influence the momentum transfer can ease the
 149 comparison with the simulations. It allows to make abstraction from the actual geometry in the
 150 simulation (bolts, accelerometers, stiffeners, little plates with optical patterns, etc.) to speed up the
 151 process of meshing and reduce the model size. While boundary and initial conditions are acquired
 152 from the actual experiment, the inertial tensor can be approximated in the numerical analysis, as
 153 schematically shown in Figure 4.



154

155 Figure 4: Comparison between experiment and simulation (IT: inertial tensor, V: rotation speed, τ :
 156 torque, BC: boundary condition, IC: initial condition)

157 The rotational momentum is a measure for the impact force and will be used to compare simulations
 158 and experiments because it is fairly independent of the set-up. If the actual impact forces would have
 159 to be calculated, the impact radius would need to be known which (a) is an estimation but worse, (b)

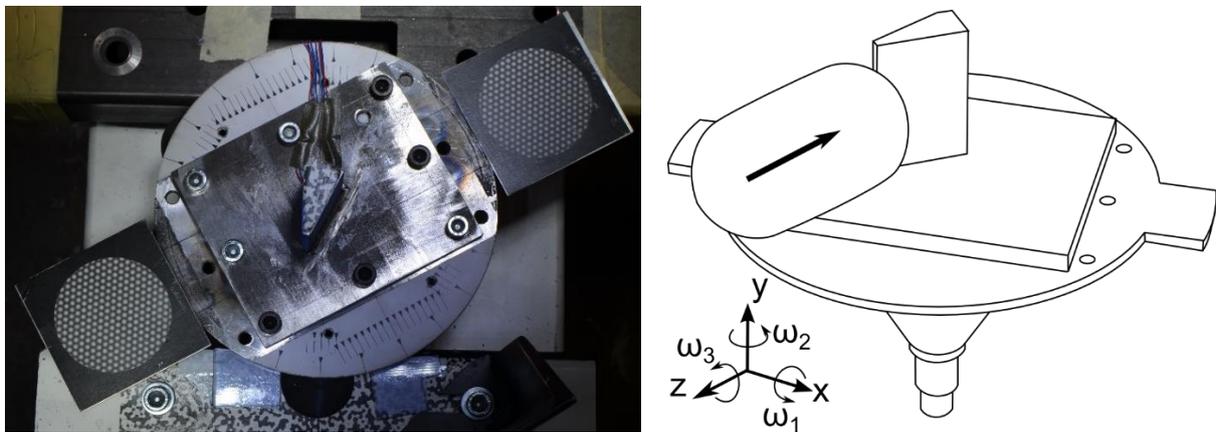
160 is not constant through time and will therefore introduce another unknown and error in the process.
161 Therefore, the momentum will remain the parameter for comparison throughout the remainder of
162 this paper.

163 The next section will explain how the rotational speeds are derived.

164 4) Deriving the rotational speeds

165 a. Overview

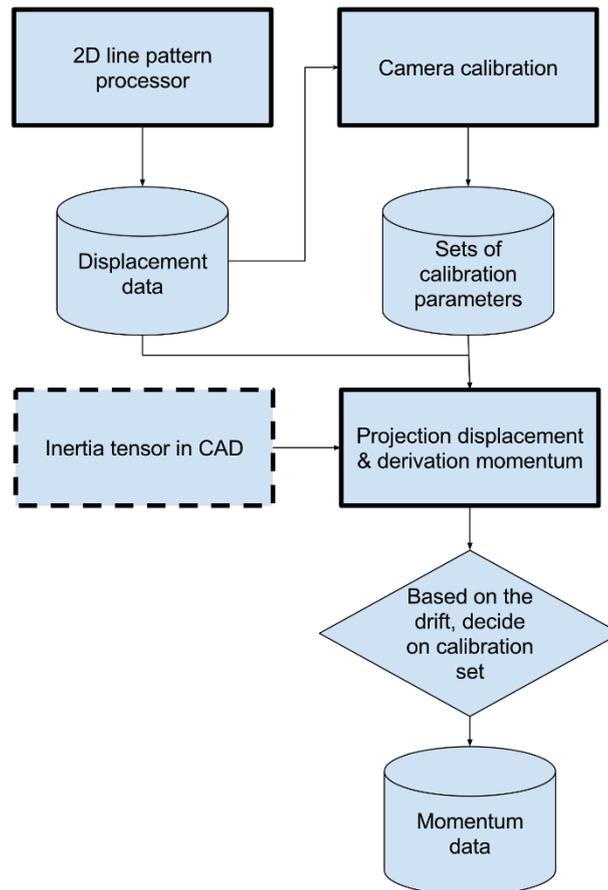
166 As mentioned in the previous sections, the kinematics of the cone are derived from an optical
167 measurement. Figure 5 shows the final concept of the cone including the steel vane with two strain
168 gauges, indications on the flange to be able to set the initial position of the cone and two optical
169 patterns at each side of the cone. On the right, the convention for the coordinate system is shown.
170 The z-axis is aligned with the impact direction, the y-axis is aligned according to the axis of the cone
171 or also the vertical upward direction and the x-axis completes the orthogonal right-handed
172 coordinate system.



173
174 Figure 5: Cone set-up with steel vane and optical patterns (left) and convention of the coordinate
175 system (right)

176 The displacements from the optical patterns are derived using a new Fourier based algorithm [Ref
177 pattern paper?]. The main advantage above a correlation based technique is that it is very robust.
178 During bird strike, a lot of debris, foam particles and pieces of bird fly above the patterns, which can

179 introduce a lot of noise in the results and therefore requires a robust algorithm. The patterns move
180 on a spherical path, which is a boundary condition to the problem and makes it possible to derive
181 kinematics with only one camera. A schematic overview of how the kinematics are derived is shown
182 in Figure 6.



183

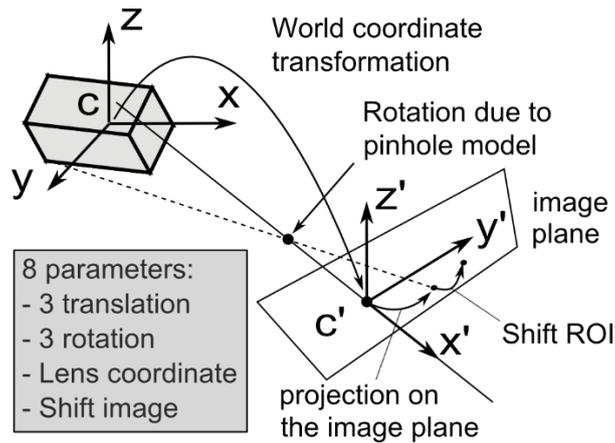
184 Figure 6: Schematic overview derivation cone kinematics

185 In a first step, the images are processed using the Fourier based technique [Ref pattern paper?]. The
186 calculated displacements however are only a projection of the actual displacements of the patterns
187 of the cone, so some transformation needs to be done. To be able to do this right, the extrinsic and
188 intrinsic camera parameters need to be defined through camera calibration. An optimization
189 procedure is used to calculate these parameters. The solution space proved to be very noisy with a
190 lot of local optima, multiple sets of calibration parameters were therefore generated after which an
191 additional parameter called the drift (which will be introduced further on) will decide on the final

192 calibration set. Once this is done, the displacements can be projected onto the sphere (the spherical
193 path on which the patterns move) and the rotational speeds and momentum can be calculated. In
194 the following sections the main steps are described in more detail.

195 b. Calibration camera

196 The camera is calibrated using 15 points, partly on the cone, partly on the patterns. Specifying these
197 coordinates in a world coordinate system and linking them to the corresponding image pixels makes
198 it possible to calibrate the cameras using an optimization technique [20]. In this work, a pinhole
199 camera model without distortion is assumed. To be able to link the world to the image coordinates,
200 four transformations are needed (Figure 7): a transformation from the world to the camera
201 coordinate system (3 translation and 3 rotation parameters), a projection on the CCD (1 parameter),
202 a 180° rotation of the CCD and a shift of the image (1 parameter). All the parameters except for the
203 last one are quite common. The last parameter is intrinsic to the high speed cameras (HSC's), as
204 there is a limit to the amount of data the HSC can save in a certain amount of time. At higher
205 framerates, the amount of pixel information or rather resolution is reduced. In this work, most
206 experiments were recorded with a framerate of 27.000 fps, which resulted in a reduced resolution of
207 448x288 (compared to the full resolution of 1024x1024). For the Photron SA-4 cameras, the position
208 of this smaller region (which will be further referred to as the region of interest or ROI) on the CCD is
209 vertically in the middle of the CCD, while the position of the ROI in the horizontal direction can be
210 chosen in steps of 32 pixels. This parameter was never noted in the experiments and is therefore an
211 additional optimization parameter.



212

213

Figure 7: Transformations from world to CCD coordinates

214

A lot of optimization techniques exist. Essentially, they can be subdivided in local and global

215

techniques. Local optimization starts from an initial guess or also set of parameters and tries to find

216

the solution starting from this input vector. A global optimization scheme on the other hand tries to

217

examine the entire solution space (in a clever way). In general, local optimization schemes are used

218

for the calibration of the camera. Multiple experiments however showed that the solution of the

219

local optimization is very dependent on the initial guess. Also with a global optimization scheme,

220

different solutions were found each run, indicating that the solution space is very noisy with a lot of

221

local optima. This phenomenon is partly the result of a poor set of almost coplanar calibration points,

222

but also due to the nature of the problem. For each experiment, therefore, multiple calibration sets

223

were generated. An additional parameter introduced in the next section was used to decide on the

224

final calibration set.

225

For the optimization, the sum of the squares of the differences between the world coordinates

226

projected on the image plane, and the corresponding actual image coordinates for each of the 15

227

points is minimized using a genetic algorithm in Matlab.

228

c. Calculation of the actual displacements

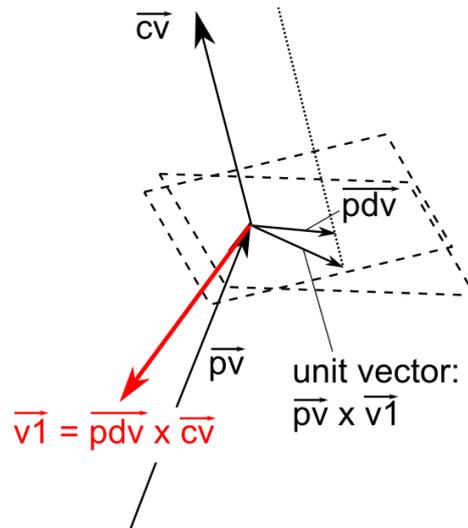
229

The displacements of the patterns are only a projection of the actual displacements. How the actual

230

displacements are determined is shown in the schematic in the figure below. In the figure, each

231 vector has an abbreviated name: cv stands for camera vector (starting from the centre of the
 232 pattern, pointing in the direction of the camera), pdv stands for projected displacement vector, this is
 233 the vector calculated by the pattern software, pv stands for position vector, this is the vector from
 234 the centre of rotation to the centre of a pattern. All the other vectors are derived from these three
 235 known vectors.



236

237 Figure 8: Calculation of the actual displacements

238 The camera vector can directly be derived from the calibration parameters. The projected
 239 displacement vector lies in the plane defined by cv, at an orientation which can also be determined
 240 from the calibration parameters. The initial position vector is defined by the initial position of the
 241 cone.

242 The actual displacement vector lies in the plane defined by pv (for small displacements, the
 243 displacement can be assumed tangential to the sphere's surface), but also in the plane defined by cv
 244 x pdv. Therefore, first a vector v1 is created, after which the wanted unit vector is obtained from the
 245 cross product of pv and v1.

246 This procedure can be repeated for each time step to obtain the actual displacements of the centre
 247 of both patterns. The mean of those displacements (from both patterns) gives the displacement of
 248 the centre of the flange of the cone, from which the two main rotation components can be derived.

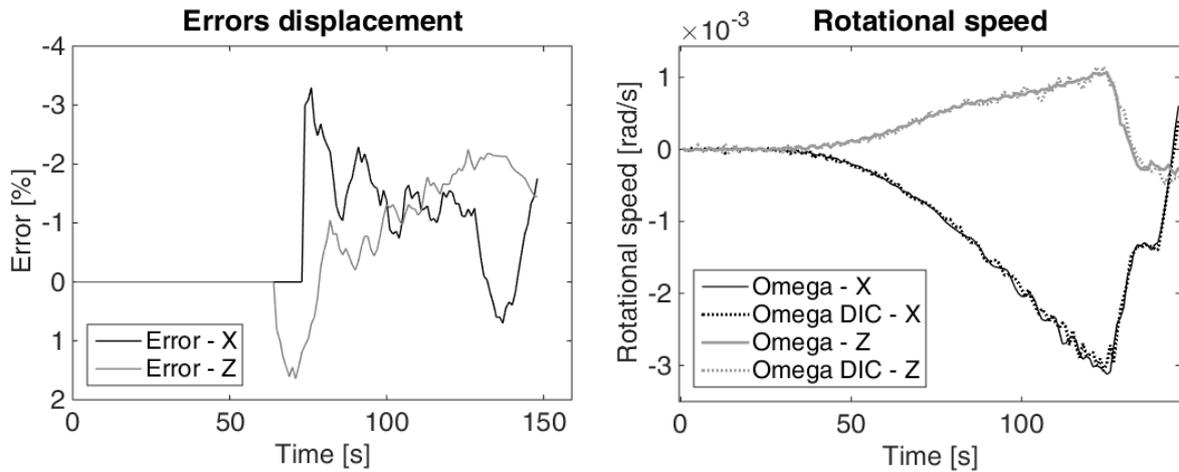
249 The calculation of the third rotation component along the axis of the cone makes use of the fact that
250 this rotation is very small. The difference between the displacement of both patterns in the direction
251 perpendicular to the axis of the cone and the vector between the patterns divided by the distance
252 between the patterns gives a good estimate of the rotation about the cone's axis. From the
253 displacements and rotations, the rotation speeds can be calculated. Together with the inertia tensor
254 obtained from the CAD drawing, the momentum can be determined.

255 The calculation is separately done for each pattern. The rigid motion of the cone implies that the
256 vector between both patterns should remain constant, but due to the fact that the whole structure is
257 not perfectly rigid and due to errors in the optimization of the calibration parameters, this is not the
258 case. Multiple tests however showed that the drift on the norm of this vector is a good measure for
259 the quality of the calibration. From the different calibration sets that are generated with the global
260 optimization scheme, the set with the least drift is therefore chosen.

261 d. Error quantification

262 The methodology to derive the momentum, including the calculation of the displacements of the
263 patterns, the calibration of the cameras and the projection on the sphere was validated using a set of
264 quasi-static and impact tests. For these experiments, speckle patterns were attached to the cone and
265 vane and tracked with a stereo DIC set-up. DIC or digital image correlation is a speckle pattern
266 tracking method with sub-pixel accuracy, which results in high resolution full-field displacement and
267 strain maps. The tracking is done by correlating small parts of the subsequent images called subsets,
268 each containing an almost unique set of speckels [21]. The orientations derived from these DIC
269 measurements were a reference to calculate the errors, because DIC is a well-known technique that
270 can be used in a stereo set-up to directly acquire 3D displacements and includes a distortion model.
271 From these experiments, it could be deduced that the errors were in general less than 5%. In Figure
272 9, the results of the two main axes of a static experiment are shown. In this quasi-static test, the cone
273 was rotated manually. The time scale is a fictitious one corresponding with one image per second. To

274 the left, the error on the displacements can be seen (cut off for low displacements). To the right, the
275 rotational speeds are shown. The figures show that good correlation is achieved.



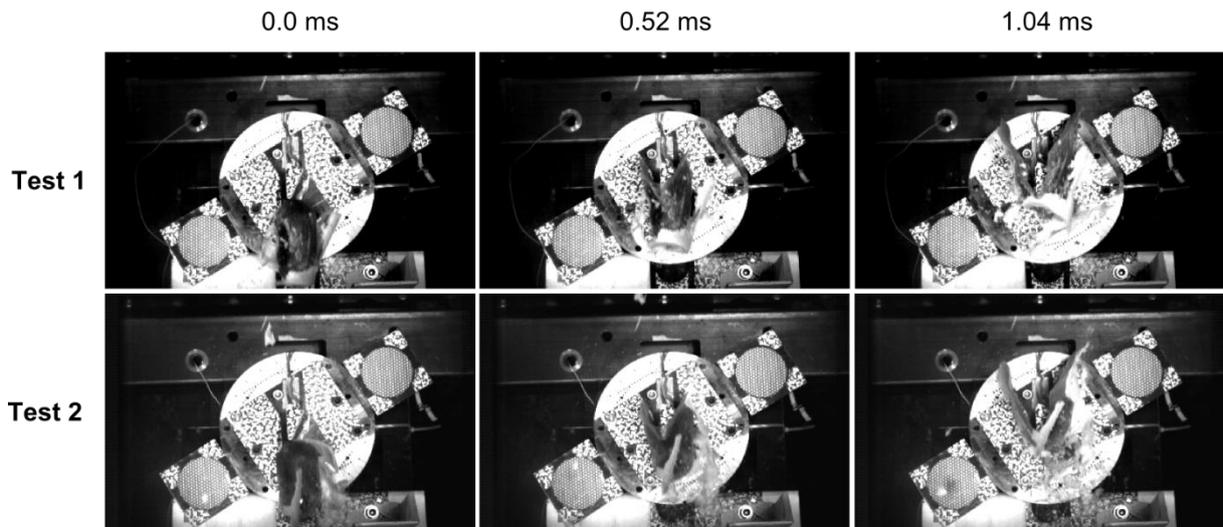
276 Figure 9: Errors and rotational speeds of a quasi-static test

277 5) Results

278 a. Impact test

279 In this section, the results of two bird strikes with the same impact conditions will be shown. Two
280 tests were executed with a gelatin bird of 300 gram with a gelatin mixing ratio of 1:6 at an impact
281 speed of approximately 110 m/s. The initial position of the cone was twice 30 degrees turned to the
282 left.

283 In Figure 10, three subsequent high speed images are shown of both tests (a top view, where the bird
284 comes from the bottom of the image). The total ROI of 448x288 pixels is shown. Little space is
285 foreseen to make sure that patterns are always in the ROI throughout the entire movement of the
286 cone. In the third image, the bird has travelled through its own length, while the movement of the
287 cone can barely be seen.



288

289 Figure 10: High speed images of two bird strikes with similar impact conditions on the steel vane

290 mounted on the cone set-up.

291 For both tests, multiple calibration sets are generated, from which the one with the least drift is

292 chosen (being 0.25 and 0.27 mm peak to peak respectively). The displacements from the patterns are

293 projected on the spherical path and the rotational speeds are calculated. The entire structure is

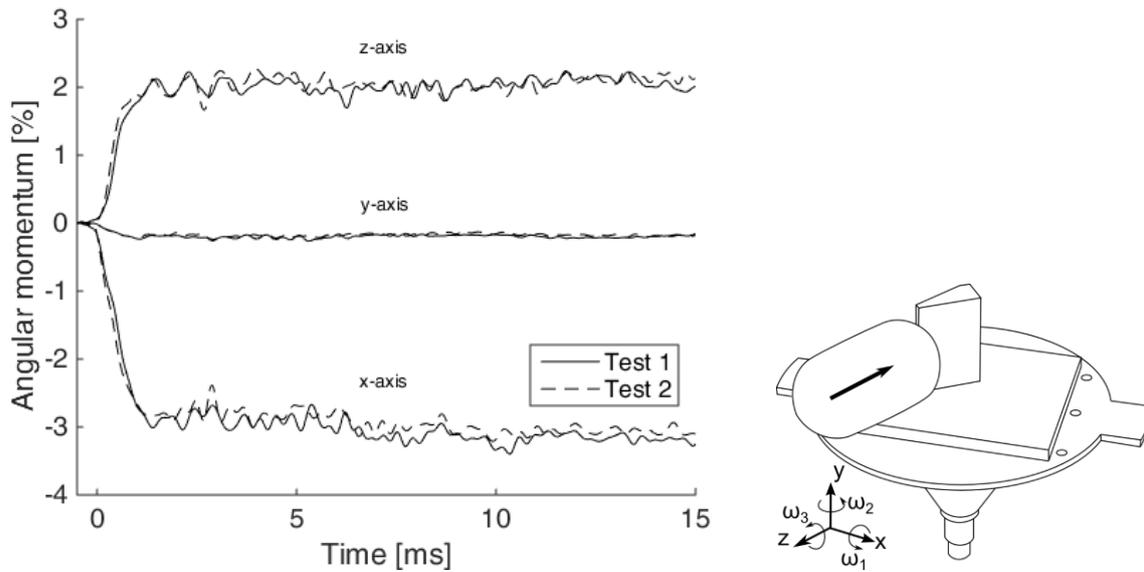
294 drawn in Solid Works to extract the inertial tensor. From the rotational speeds and the inertial

295 tensor, the momentum is calculated. To get a better relative comparison and to take the different

296 speed and mass of both birds into account, the calculated momentum of the cone is divided by the

297 initial momentum of the bird. To get a better overview of the results, a low pass filter of 3 kHz was

298 applied. The result of this process is shown in Figure 11.



299

300

Figure11: Normalized momentum data test 1 and 2

301 The two tests give very similar results, for all the three axes. Most of the mass is deflected to the
 302 right (positive x-axis), which makes the cone move to the left and back (the negative x- and z-axis).
 303 This corresponds with a positive and negative rotational momentum around the z-axis and x-axis
 304 respectively. The magnitude of the momentum around the z-axis is smaller than around the x-axis
 305 because some mass is also deflected to the left (negative x-axis). The rotational momentum around
 306 the y-axis is a lot smaller because the radius of the impact force with respect to the axis of the cone is
 307 a lot smaller than for the other axes. The sign is negative because most of the deflection happens at
 308 the front of the steel vane, which is positioned slightly to the left of the axis of the cone from the
 309 point of view of the bird.

310

b. Correlation with a numerical model

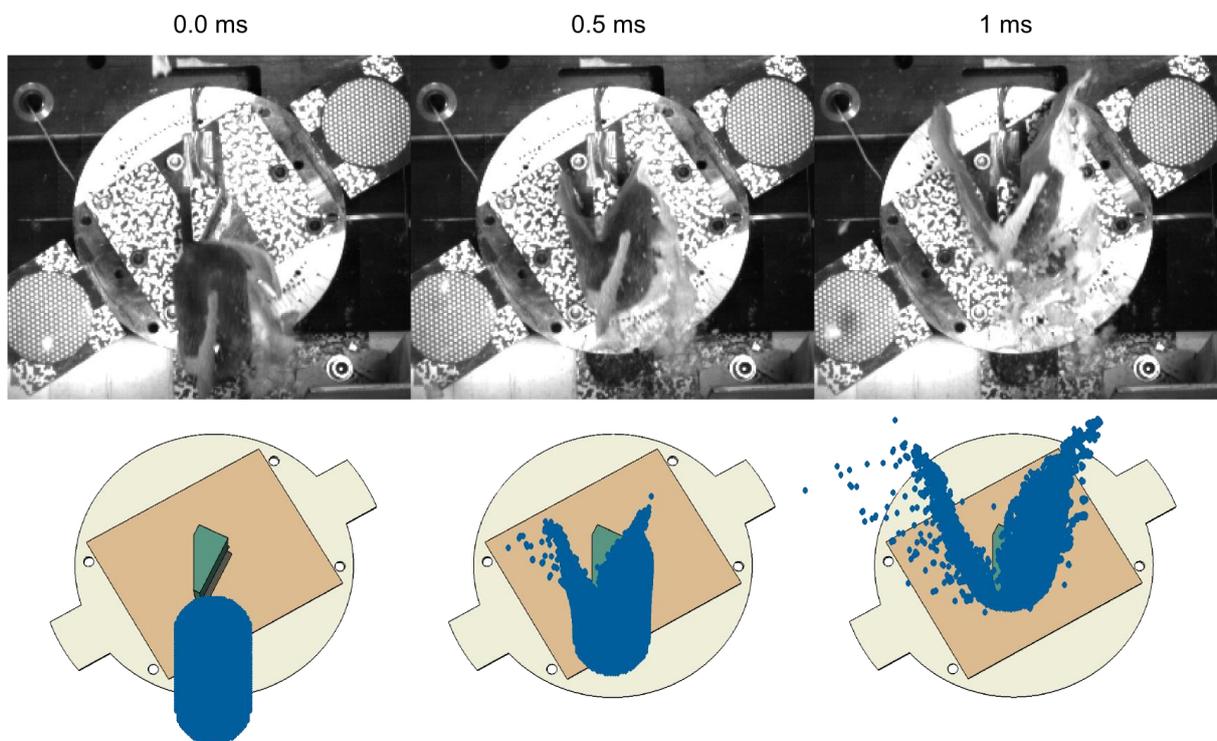
311 An explicit simulation was ran with the same impact conditions as test 2, where the bird was
 312 modelled with smoothed particle hydrodynamics or SPH. SPH is increasingly used in bird strike
 313 simulations as it already proved to be quite capable of simulating high deforming matter with
 314 defragmentation [22]. Also, tracking of field values is not a problem for SPH. A complete and clear
 315 explanation of SPH and its governing equations can be found in literature [23]. A structured mesh is

316 generated based on the shape of the mould, with a slight offset to exactly match the mass with the
317 experiments. The impact position is derived from the high speed images as best as possible (both in
318 the vertical as the horizontal direction).

319 For the numerical model, the cone was modelled as a deformable object, able to rotate around the
320 tip of the cone. The vane was modelled as accurately as possible, including the welds and
321 membranes for the strain gauges. Both the cone and the vane were modelled with reduced
322 integration hexahedral elements.

323 A linear Mie-Grüneisen EOS is used for the bird material model, which relates the pressure to the
324 density. Parameters for porcine gelatine were found in literature ($c_0 = 1570$ m/s and $s = 1.77$) [24],
325 which is very similar to water. The bird was tilted 8° in the vertical plane to have a better match with
326 the impact conditions of the experiment.

327 Figure 12 shows a qualitative comparison from the top view of the cone. The difference in time
328 should be 0.05 ms maximum.



329

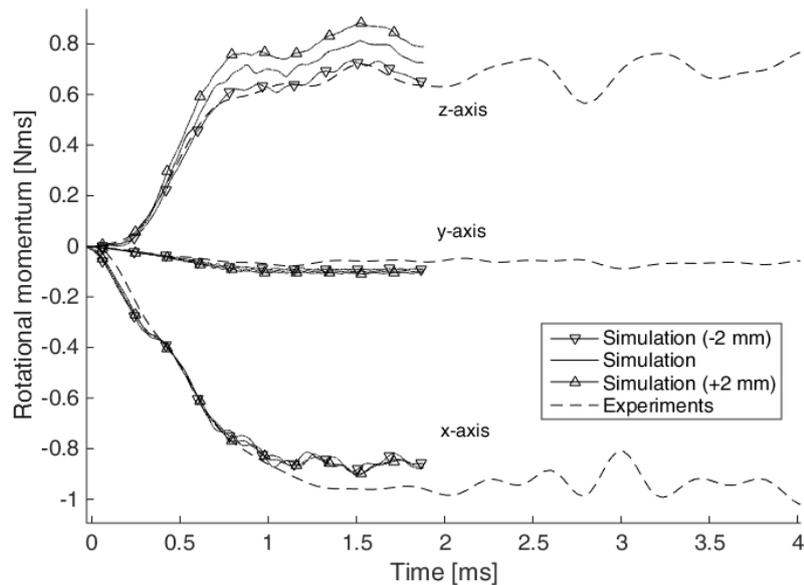
330

Figure 12: Qualitative comparison between experiment and simulation

331

332 The global behaviour correlates well with the experiment. Both the direction in which the mass is
333 deviated as the speed of the deviated mass after impact.

334 Figure 13 shows a comparison between the momentum obtained from the simulation and from the
335 experiments, including a 2mm offset on the estimation of the impact location both in the negative as
336 the positive x-direction in the simulation. For the momentum of the simulation, the speeds are
337 extracted at two locations on the cone. Together with the inertia tensor from the model, the
338 momentum can be obtained.



339

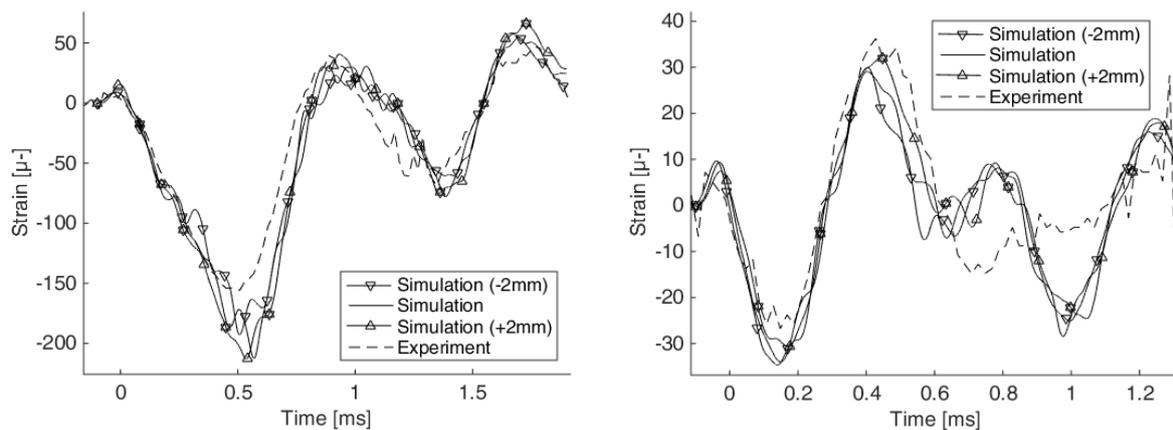
340 Figure 13: Comparison of the rotational momentum in the simulation and the experiment

341 The impact location of the bird is crucial. A 2mm offset error for example, which can be easily made,
342 corresponds with 4.5% of the bird mass that is deviated in the other direction. This kind of error is
343 therefore most represented in the momentum around the z-axis. The order of magnitude of
344 momentum transfer is approximated well, but some differences can be observed. In the simulation,
345 the momentum transfer to the x-axis is slightly lower. The graphs also shows that possibly, an error

346 of approximately 2mm was made in the estimation of the x-offset. The momentum transfer around
347 the x-axis also takes slightly longer in the experiment.

348 A comparison between the strain in the simulation and experiment at the back and the left side is
349 shown in Figure 14. The influence of the offset is less pronounced in the strain signals. The left strain
350 gauge measures the deformation along the weak axis of the steel vane and therefore sees more
351 strain. For the left strain gauge, the strain amplitude is 15-20% higher in the simulation. And after
352 0.6 ms, the correlation with the strain gauge at the back gets worse. Apart from that, the response of
353 the vane is capture quite well.

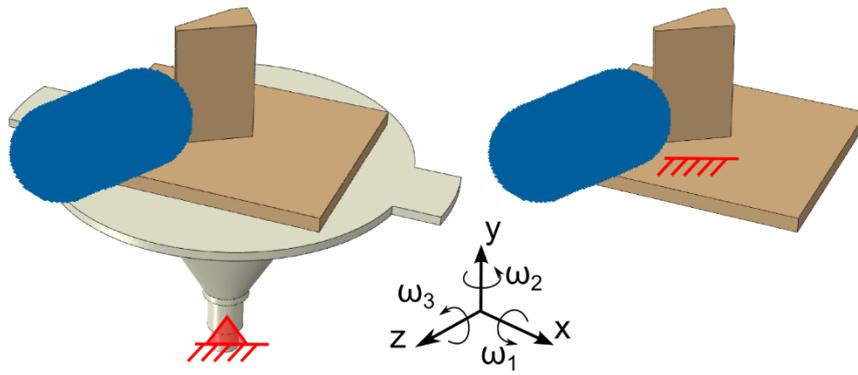
354



355 Figure 14: Comparison of the strain at the left side (left) and the back (right) of the vane in the
356 simulation and the experiment

357 c. Influence of the cone

358 To show that the influence of the cone is negligible (Figure 15), a reference simulation with and
359 without cone is performed as well. For the simulation without cone, the ties that connected the vane
360 to the cone are connected instead to a fixed reference point. The forces and moments acting on this
361 reference node are recorded throughout the simulation.



362

363

Figure 15: Simulation with and without cone

364

In a post processing step, the forces and moments are transformed to the actual location of the ball

365

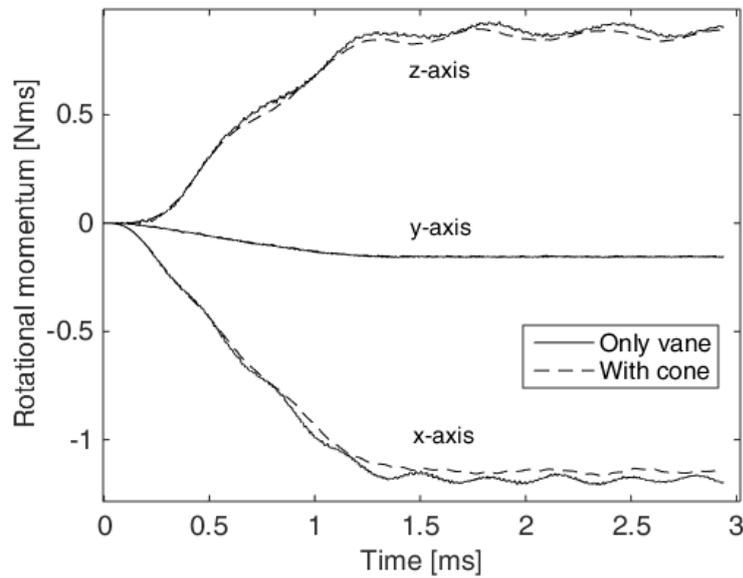
joint (with respect to the vane), integrated over time (as in equation 3) and plotted together with the

366

momentum obtained from kinematics of the simulation with the cone shown in the previous section.

367

The results are depicted in Figure 16.



368

369

Figure 16: Comparison rotational momentum with and without cone (simulation)

370

It can be observed that the transferred momentum is comparable. The momentum transfer is slightly

371

lower with cone, but still sufficiently small compared to other errors sources.

372

6) Conclusion

373 In this paper, the development and results of a multi-axial momentum measurement tool are
374 discussed. The multi-axial force measurement is realized using a well-chosen mass in the shape of a
375 cone which is able to rotate freely around the tip of the cone. The ability to rotate freely makes it
376 possible to determine the transferred rotational momentum. This requires that the kinematics of the
377 cone, obtained from an optical measurement in three main steps: calculation of the projected
378 displacements from the optical patterns, calibration of the camera using a global optimization
379 technique and transformation of the projected displacements on the spherical path on which the
380 optical patterns move.

381 Quasi-static and impact tests show that the developed tool can be used to get a reliable estimate of
382 the transferred rotational momentum. Two bird strike tests with the same impact conditions are
383 compared and give very comparable momentum transfer results. A numerical simulation of one of
384 these tests (using SPH for the bird) shows a good correlation in terms of momentum. Finally,
385 simulations indicate that the tool has a negligible influence on the experiment for the considered stiff
386 steel vane. The tool proved to be very useful and successful results have been obtained for multiple
387 sets of real Titanium vanes.

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392 More information can be found at <http://www.e-break.eu/>.

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