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title

FRAP on the Confocal Laser Scanning Microscope: a Generalized Model without Restriction on the Size of the Photobleached Disk.

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Abstract

Fluorescence recovery after photobleaching (FRAP) carried out on a confocal laser scanning microscope (CLSM) performs well for photobleached disks that are large compared to the resolution of the bleaching beam. For smaller disks approaching this resolution, current FRAP models providing a closed-form solution, do not allow to extract the diffusion coefficient accurately. The new generalized disk model presented here addresses this shortcoming by bringing into account the bleaching resolution and the total confocal imaging resolution. A closed-form solution is obtained under the assumption of linear photobleaching. Furthermore, simultaneous analysis of FRAP data collected at various disk sizes allows for the intrinsic determination of the instrumental resolution parameters, thereby obviating the need for an extrinsic calibration. A new method to estimate the variance of FRAP data was introduced to allow for proper weighting in this global analysis approach by nonlinear least squares. Experiments were performed on two independent CLSMs on homogeneous samples providing validation over a large range of diffusion coefficients.

Keywords (max. 8):

Fluorescence Recovery After Photobleaching

FRAP / FPR

CLSM / LSCM

Diffusion

Artifacts

Global analysis

Weights

1 Introduction

Fluorescence recovery after photobleaching (FRAP, also named fluorescence photobleaching recovery FPR) is a well-known microfluorimetric technique for measuring the diffusion of fluorescently labeled molecules on a micrometer scale.¹ Fluorescent molecules in a defined area are quickly and irreversibly bleached by irradiation with light of high intensity. These photobleached molecules will subsequently be replaced by diffusion of intact fluorescent molecules from the surroundings. The resulting gradual recovery of the fluorescence over time in the defined area, observed by using excitation light of low intensity, holds information on this diffusion process. Subsequent analysis of the fluorescence recovery data using a suitable FRAP model yields the translational diffusion coefficient and the mobile fraction.

In contrast to its fundamental principle, the instrumental implementation of FRAP changed over time. The regular fluorescence widefield microscope with a stationary laser beam for bleaching is often replaced by a confocal laser scanning microscope (CLSM). Commercial availability, user-friendliness, and optimized optical performance have lead to the widespread use of these CLSMs in life science laboratories. The combination of raster scanning and fast modulation of the laser beam intensity enables a CLSM to bleach and monitor arbitrary regions and makes it an easily accessible FRAP tool.²

Together with this technical evolution new FRAP models are created for quantitative diffusion analysis. The models for non-scanning microscopes and two-dimensional (2D) diffusion^{1,3,4} are replaced by CLSM-dedicated models. Most of these models, however, are either limited to 2D diffusion,⁵ or apply a numerical approach with associated complexity.⁶⁻⁸ One of the first FRAP models for analyzing CLSM FRAP data by using a simple closed-form equation, was based on the photobleaching of a uniform disk that is much larger than the effective resolution of the bleaching beam. We will refer to this model as the uniform disk

model (UDM).² As demonstrated by Braeckmans *et al.*, the assumed bleach profile with sharp boundaries is no longer obtained when using smaller bleach regions (region of interest, ROI) due to the finite resolution of the bleaching beam. In this context, small or intermediate ROIs are defined as having a radius smaller than five times the resolution of the bleaching beam.² Diffusion coefficients obtained with these small ROIs will significantly underestimate the actual diffusion values when analyzed by the UDM.

Some applications require small or intermediate sized ROIs or do not allow for large ROIs at all. A clear example can be found in biological cells, which are limited in size by nature. Also investigation of anomalous diffusion of proteins requires the use of a range of ROI sizes.⁹⁻¹¹ Here we present a new, closed-form generalized FRAP disk model (GDM) which does not impose restrictions on the size of the circular bleached area. This is achieved through modification of the UDM by bringing into account both the bleaching resolution as well as the confocal imaging resolution. A procedure with simultaneous analysis of recovery curves derived from bleach regions of various sizes, i.e. a global analysis of the resulting multidimensional data surface, is introduced for increased accuracy and a calibration free approach. A new method to estimate the variance of FRAP data was introduced to allow for proper weighting in this global analysis approach by nonlinear least squares. By a detailed experimental validation, we demonstrate that a wide range of diffusion coefficients can be accurately retrieved independent of the size of the bleached disk.

2 Theoretical framework

2.1 Effect of the finite width of the scanning laser beam on the bleached region

The derivation below of the GDM is valid for 2D diffusion in an infinite plane XY and single photon photobleaching and imaging. For this situation, it is shown by Braeckmans *et al.*² that

the concentration C_b of fluorophores after irreversible photobleaching of a 2D geometry $B(r)$ with rotational symmetry around the Z -axis by a scanning beam can be described by:

$$C_b(r) = C_0 e^{-\frac{\sigma q}{v\Delta y} K(r)} \quad (1)$$

where C_0 is the homogeneous initial fluorophore concentration, v the line scanning speed, Δy the distance between consecutive scanning lines, σ the cross-section for single photon absorption, and q the quantum efficiency for single photon photobleaching. The bleaching geometry $K(r)$ that results from scanning the bleaching geometry $B(r)$ with the effective bleaching intensity distribution $I_b(r)$ of the scanning beam, can be calculated from the convolution product of $B(r)$ and I_b according to [Figure 1]:

$$K(r) = B(r) \otimes I_b(r) = \int_{r'=0}^{+\infty} \int_{\theta'=0}^{2\pi} r' B(r') I_b \left(\sqrt{r^2 + r'^2 - 2rr' \cos \theta'}, t \right) d\theta' dr' \quad (2)$$

The effective bleaching intensity distribution $I_b(r)$ should not be confused with the intensity distribution of the illuminating laser beam. It was shown before that the effective photobleaching distribution is a complex function of the laser intensity distribution, the fluorophore photochemistry, the photon flux and the physicochemical local environment.¹²⁻¹⁴ Nor does $I_b(r)$ equal the imaging point-spread function (PSF) of a CLSM. The latter is the product of the illumination PSF at the excitation side and the detection PSF at the emission side,¹⁵ while during the bleaching process only the excitation side matters.

The bleaching intensity distribution is approximated by a 2D Gaussian characterized by the effective bleaching resolution r_b :

$$I_b(r) = I_{b0} e^{-2\frac{r^2}{r_b^2}} \quad (3)$$

The convolution in Eq. (2) implies a modulation of $B(r)$ with the effective bleaching intensity distribution whose effect on the recovery process will increase as the dimensions of $B(r)$ approach r_b .

In order to obtain a closed-form solution further on for the recovery process, only a small amount of photobleaching is assumed, i.e. $\frac{\sigma q}{v\Delta y} K(r) \ll 1$, such that Eq. (1) can be linearized.

Under this approximation, Eq. (1) can be expressed in terms of photobleached molecules C_b^* at time $t = 0$ as:

$$C_b^*(r) = C_0 - C_b(r) = C_0 \frac{\sigma q}{v\Delta y} K(r) \quad (4)$$

Combining Eq. (4) together with Eq. (2) and Eq. (3) yields:

$$\begin{aligned} C_b^*(r) &= \frac{\sigma q}{v\Delta y} I_{b0} C_0 e^{-\frac{2r^2}{r_b^2}} \int_{r'=0}^{+\infty} r' B(r') e^{-\frac{2r'^2}{r_b^2}} dr' \int_{\theta'=0}^{2\pi} e^{\frac{4rr'\cos\theta'}{r_b^2}} d\theta' \\ &= 2\pi \frac{\sigma q}{v\Delta y} I_{b0} C_0 e^{-\frac{2r^2}{r_b^2}} \int_{r'=0}^{+\infty} r' B(r') I_0\left(\frac{4rr'}{r_b^2}\right) e^{-\frac{2r'^2}{r_b^2}} dr' \end{aligned} \quad (5)$$

where I_0 is the modified Bessel function of the 0th order. In order to obtain the recovery of fluorescence after bleaching, Fick's second law of diffusion has to be solved for the initial condition in Eq. (5). A solution can be easily found by noting that this expression resembles closely the general solution of the diffusion equation in cylindrical coordinates for a radially symmetric initial distribution $f^*(r)$ of bleached molecules and for a diffusion constant D :

$$C_{bf}^*(r, t) = \frac{e^{-\frac{r^2}{4Dt}}}{2Dt} \int_{r'=0}^{+\infty} r' f^*(r') e^{-\frac{r'^2}{4Dt}} I_0\left(\frac{rr'}{2Dt}\right) dr' \quad (6)$$

where $C_{bf}^*(r, t)$ denotes the concentration of the bleached molecules arising from the initial distribution $f^*(r)$.

If now an initial concentration $f^*(r)$ and time $t = t_0$ can be found for which $C_{bf}^*(r, t_0)$ becomes identical to $C_b^*(r)$, the solution for Eq. (6) can be used and the intended closed-form solution can be obtained. This situation is met for $t = \frac{r_b^2}{8D}$, since it can be shown that

$$C_{bf}^*\left(r, \frac{r_b^2}{8D}\right) = C_b^*(r) \text{ for an initial concentration}$$

$$f^*(r) = C_0 \frac{\sigma q}{v\Delta y} \frac{\pi}{2} I_{b0} r_b^2 B(r) \quad (7)$$

as can be verified by substitution in Eq. (6). This initial concentration $f^*(r)$ in Eq. (7) can be considered to originate from bleaching a geometry $B(r)$ with a beam of infinite radial resolution $I_{b\delta}(r) = I_{b\delta 0} \delta(r)$, where $\delta(r)$ is the Delta-Dirac function. This can be seen as follows. Assuming the bleaching process to be linear, one obtains in analogy with Eq. (4):

$$\begin{aligned} C_{b\delta}^*(r) &= C_0 \frac{\sigma q}{v\Delta y} B(r) \otimes I_{b\delta}(r) \\ &= C_0 \frac{\sigma q}{v\Delta y} B(r) I_{b\delta 0} \end{aligned} \quad (8)$$

By comparing Eq. (7) and Eq. (8) we find that $I_{b\delta 0} = I_{b0} \frac{\pi}{2} r_b^2$ and Eq. (7) becomes

$$f^*(r) = C_0 \frac{\sigma q}{v\Delta y} I_{b\delta 0} B(r) = C_0 K_0 B(r) \quad (9)$$

where K_0 is the photobleaching parameter that determines the bleaching depth as used in the uniform disk model.²

In conclusion, the concentration distribution $C_b^*(r)$ as obtained by bleaching the geometry $B(r)$ with a laser beam of finite resolution characterized by r_b is the same distribution as obtained by diffusion from an initial concentration $C_0 K_0 B(r)$ after a time $t_0 = \frac{r_b^2}{8D}$ has elapsed. In other words, immediately after the bleaching phase, the bleached region with its

shallow slopes created by a scanning laser beam can be regarded as originating from the perfect uniform disk with sharp boundaries through a diffusion process which started a time

$t_0 = \frac{r_b^2}{8D}$ earlier. Introduction of the time shift $t \rightarrow t + \frac{r_b^2}{8D}$ in the UDM is sufficient to bring

the finite width of the bleaching scanning laser beam into account.

2.2 Effect of finite total detection resolution on the time evolution of FRAP

As the photobleached disk can now be of any size, the overall detection resolution r_d of the CLSM cannot be neglected anymore for very small or intermediate disks. This requires a revision of the uniform disk formula as described by Braeckmans *et al.*² The original UDM formula for 2-D diffusion is given here for the convenience of the reader:

$$\frac{F'(w, t)}{F'_0(w)} = 1 + (e^{-K_0} - 1) \left[1 - e^{-\xi} (I_0(\xi) + I_1(\xi)) \right] \quad (10)$$

where $\xi = \frac{w^2}{2Dt}$, $F'(w, t)$ is the integrated fluorescence over the photobleached disk with radius w at time t after the bleaching period as observed by the CLSM, $F'_0(w)$ is the initial integrated fluorescence level within the ROI before bleaching, and I_0 and I_1 are the modified Bessel functions of 0th and 1st order, respectively. It has to be noted that r_b is neglected in the UDM model which is valid for large ROIs.

Nevertheless, the UDM can be easily extended to take r_d into account (calculations not shown

here) leading to $\xi = \frac{w^2}{\frac{r_d^2}{4} + 2Dt}$ in Eq. (10). Note that for a large disk ($w \gg r_d$) this indeed

reduces to the familiar uniform disk formula.

Taken together, linearizing the photobleaching process, introducing the time shift as discussed above and taking the total imaging PSF into account, finally leads to the expression:

$$\frac{F(w, t)}{F_0(w)} = 1 - K_0 \left[1 - e^{-\xi} (I_0(\xi) + I_1(\xi)) \right] \quad (11)$$

where $\xi = \frac{w^2}{\frac{r_d^2}{4} + 2D \left(t + \frac{r_b^2}{8D} \right)} = \frac{w^2}{2Dt + R}$ with $R = \frac{r_d^2 + r_b^2}{4}$ and where $F(w, t)$ is the integrated

fluorescence over the photobleached disk with radius w at time t , $F_0(w)$ is the initial value of $F(w, t)$ before bleaching. Eq. (11) will be referred to as the GDM since it is the generalization of the classic UDM taking into account the effective resolution of the bleaching beam and the microscope imaging resolution.

Common assumptions between GDM and UDM are initially uniformly distributed fluorescence molecules, an isotropic diffusion process in an infinite medium, absence of flow, a sufficiently short bleaching phase so as to neglect diffusion during bleaching and 2-D diffusion.² The latter assumption is satisfied for low NA lenses (which cause a cylindrical bleach profile) or lenses of high NA when the sample thickness is small compared to the axial resolution of the lens, as is e.g. the case for biological cell membranes.

Finally we note that the presence of an immobile fraction of molecules M can be taken into account by substituting the expression in Eq. (11) into the right hand side of:

$$F(w, t) = M \frac{F(w, t)}{F_0(w)} + (1 - M) \frac{F(w, 0)}{F_0(w)} \quad (12)$$

3 Materials and Methods

3.1 FRAP equipment

FRAP experiments were performed on two independent CLSMs. The first microscope (*setup A*) was an MRC1024 UV (Bio-Rad, Hemel Hempstead, UK) equipped with a custom-built FRAP module and a 4 W Ar-ion laser (model Stabilite 2017; Spectra-Physics, Darmstadt, Germany).^{2,14} A 10× objective lens (CFI Plan Apochromat; Nikon, Badhoevedorp, The

Netherlands) with a numerical aperture (NA) of 0.45 was used. The back aperture of this lens was only partially filled, resulting in a lower effective NA of ~ 0.2 and an r_d value of $1.0\ \mu\text{m}$ as determined from subresolution beads.¹⁴

The second confocal setup (*setup B*) was an LSM 510 META (Carl Zeiss, Jena, Germany) installed on an Axiovert 200 M motorized frame (Carl Zeiss, Jena, Germany). It was equipped with a 30 mW Ar-ion laser and a $10\times$ objective lens (Plan-Neofluar; Zeiss, Jena, Germany) with a NA of 0.3. The overall detection resolution r_d was $0.9\ \mu\text{m}$, as determined using sub-resolution beads. For both setups, a laser power of 0.5 to 1 mW at the sample was used for photobleaching.

3.2 Sample preparation

FRAP measurements were performed on solutions of FITC-labeled dextran (FD) molecules (Sigma-Aldrich, Bornem, Belgium) with a molecular weight of 2000 kDa or 464 kDa. These compounds will be referred to as FD2000 and FD500, respectively. All stock solutions were prepared in HEPES buffered solution at pH 7.4. To increase the viscosity, solutions with 40% and 56% (w/w) sucrose (VWR Prolabo, Leuven, Belgium) were prepared from the FD stock solutions. The final fluorescence signal scaled linearly with the concentration as verified experimentally on each setup individually.

For FRAP experiments, $5\ \mu\text{l}$ of the solution was sandwiched between a microscope slide and cover glass with Secure-Seal stickers (Sigma, Bornem, Belgium) of $120\ \mu\text{m}$ thickness in between. This avoids any detectable flow inside the solutions while maintaining a 3D volume.

3.3 FRAP protocol

All measurements were performed at room temperature. A fresh homogeneous region of the sample of interest was brought into focus before the start of each FRAP experiment. The

execution of the experiments was controlled through the automated microscope bleach control software. The resulting stack of images represented a time-series recording with three consecutive phases [Figure 2]. The first phase was marked by a set of one to five images showing the sample before bleaching. In the second phase, the user-defined circular ROI was bleached. The duration of the bleach phase did not exceed $1/10^{\text{th}}$ of the characteristic recovery time τ_D ,¹⁶ where τ_D is defined as:

$$\tau_D = \frac{w^2}{4D} \quad (13)$$

Only one bleach iteration was allowed for all ROIs. Depending on the experiment, the nominal ROI radius as set in the control software ranged from $1.0 \mu\text{m}$ to $15.2 \mu\text{m}$.

Only the *setup A* registered fluorescence intensities during bleaching, resulting in one image showing the disk at the time of bleaching (hereafter called bleach image). All subsequent images showed the recovery of the fluorescence after the bleach-phase. This is the post-bleach phase. The time-interval between the images and the total acquisition time were selected so that when scaled to the recovery time τ_D a similar distribution of data points for all ROI sizes was obtained. Typically a time series of 50 images was recorded with a time of $\tau_D/3$ between the images as a trade-off between obtaining sufficient sampling of the recovery phase and limiting photobleaching due to imaging.

3.4 Recovery curve extraction and variance estimation

All data extraction was performed in MATLAB (The MathWorks BV, Eindhoven, The Netherlands) using custom written routines. First, the coordinates of the circular bleach ROI were determined. For *setup A*, the center of the bleach disk was determined using a center-of-mass algorithm using the bleach image. For *setup B* this information was obtained from the metadata of the image sequence using a home-written routine.

Secondly, all pixels within the ROI were integrated frame by frame. A background region was selected at a distance of at least 4 times the radius of the bleach ROI. The recovery curve was corrected for changes in fluorescence intensity (e.g. by photobleaching during imaging or laser fluctuations) by normalizing the ROI fluorescence intensity by the background intensity. Finally, the recovery values were normalized to the mean pre-bleach intensity according to Eq. (11).

To allow weighted least-squares fitting of the experimental data to the theoretical model, the weight of each data point of this normalized recovery curve was the inverse of the estimated variance, taking into account all applied mathematical operations. Key to this variance estimation is the variance of the individual pixels, which was determined using the observed linear relationship between the unnormalized average pixel intensity and variance of homogeneous regions as described by Dalal *et al.* [Figure 3(a)].¹⁷ The average and variance of selected homogeneous regions from the recorded time series were calculated and resulting average-variance pairs with identical experimental settings were pooled across the data set to determine this linear relationship. Application of the propagation of errors subsequently when integrating over the number of pixels inside the ROI and considering other correction and normalization steps results in distinct variances associated with large and small ROIs [Figure 3(b) and (c)].

3.5 Analysis of recovery curves

Each set of recovery curves was analyzed in two distinct ways. In the first approach, further referred to as single curve analysis, each curve was separately analyzed utilizing the calibrated resolution parameters r_b and r_d . In the second approach, two global analyses of all recovery data were performed. These global analyses respectively utilized and ignored the calibrated values of the resolution parameters to investigate the feasibility of recovering the

combined resolution parameter R (cfr. Eq. (11)) exclusively from the recovery data themselves.

In the single curve analysis, each experimental recovery curve was fit to the UDM [Eq. (10)] and to the new GDM [Eq. (11)], both adjusted for a mobile fraction M through Eq. (14), using a weighted non-linear least-squares optimization minimizing

$$\chi^2 = \sum_{i=1}^N \beta_i \left(y_{i_{fit}} - y_{i_{data}} \right)^2 \quad (14)$$

where N is the number of points in the fit and β_i is the inverse of the estimated variance of the i^{th} point. These weights were used to obtain statistically justified χ^2 values. Standard deviations on the recovered parameter values were obtained using a modification of the MATLAB routine *nlparci*.

For *setup A*, the photobleaching resolution r_b for both models was set to 2.5 μm : 2 μm as determined from lineFRAP experiments¹² increased with 0.5 μm to account for the 2 pixel rise time of the acousto-optical modulator (AOM). For *setup B*, r_b was set to 1.9 μm : 1.4 μm calibrated using lineFRAP (FD500 in 56% (w/w) sucrose, data not shown) increased with 0.5 μm analogously to *setup A*.

For the GDM, r_d was set to 1.0 μm and 0.9 μm for *setup A* and *B*, respectively. The mobile fraction M was set freely adjustable for consistency check: for the FITC-dextran solutions M should approximate one. Results were grouped per ROI radius by calculating the weighted average values of the obtained parameter values.

For simultaneous analysis of a set of recovery curves, home-written routines were used. These routines make use of a non-linear least-squares optimization algorithm. Parameters can be linked across recovery curves, i.e. made global, such that their values are equal for all selected recovery curves. Several sets of initial guesses were used to verify the true global minimum.

The obtained results were compared with the values obtained using the Stokes-Einstein equation

$$D = \frac{kT}{6\pi\eta r_H} \quad (15)$$

where r_H is the hydrodynamic radius of the diffusing molecules, k the Boltzmann constant, η the dynamic viscosity and T the absolute temperature. The hydrodynamic radius of the FDs was calculated according to the relationship reported by Braeckmans et al.² The dynamic viscosity η was obtained from the literature.¹⁸

4 Results

4.1 Comparison of uniform disk model with new generalized disk model: single curve analysis

FD2000 solutions were used for *setup A* and FD500 solutions for *setup B*. The viscosity of the solutions was increased using different amounts of sucrose (40% and 56% (w/w)) to cover a range of diffusion coefficients (0.205 to 2.22 $\mu\text{m}^2/\text{s}$). For a given viscosity FRAP experiments were carried out on the same sample using photobleaching disks of varying radii. Hence, it was expected that all experiments would yield identical results.

Each experiment was analyzed by single curve analysis using the classic UDM and the new GDM. Representative results are shown in Figure 4. For all FRAP experiments analyzed, the mobile fraction M was close to 1 for both models, as expected from the model system used in these experiments. For *setup A* with $r_b = 2.5 \mu\text{m}$, it could be expected that for the UDM the calculated D values are independent of the radius for radii $>10 \mu\text{m}$, i.e. 4-5 times the effective photobleaching resolution. For smaller radii, D values recovered using the UDM were expected to gradually decrease because of the underestimation of the effective ROI size. For the smallest disk size, the UDM underestimates the expected diffusion coefficient by a factor

of 2. On *setup B* an entirely similar trend was observed, demonstrating that the deviations are not instrument related [Figure 4(b)].

Analysis of the same data sets with the GDM nicely resulted in diffusion values that are independent of the size of the disk (within the experimental error). We note that small bleaching depths were used, $K_0 < 0.3$, consistent with the linearization in Eq. (4). Thus, the underestimation of the effective radius of the ROI in the UDM was appropriately compensated by the new model. For large ROIs, the D -values of the UDM asymptotically approach the values of the GDM, as can be expected.

4.2 Comparison of uniform disk model with new generalized disk model: global analysis

Since all recovery curves in a data set are derived from the same sample, they are all characterized by the same diffusion coefficient D . This is often implemented in the data analysis by averaging over all measurements, either at the level of the recovery curves themselves or at the level of the measured D values. In the current paper, a third implementation is used: a simultaneous analysis with the parameter D linked across all related recovery curves, i.e. a single, global parameter for D shared by all recovery curves. Other parameters like the bleaching depth and the mobile fraction can be kept local, i.e. these fitting parameters are specific for each individual recovery curve. This approach allows for differences in these parameters between recovery curves.

All data sets were analyzed using this global analysis applying both UDM and GDM. Similar as in the single curve analysis, the latter model made use of the calibrated resolution factor R by incorporating it as a fixed, i.e. non-adjustable parameter. For all data sets globally analyzed, D recovered using the UDM was consistently smaller as compared to the GDM that better approached the theoretical expected value [Table 1]. The uncertainties on these recovered values are similar for both models. The underestimation of D by the UDM,

however, was smaller than observed using solely the smallest ROI size in the single curve analysis. In other words, inclusion of large ROI sizes moderates the underestimation of D , as can be expected.

Single curve analysis of FRAP data using the GDM requires *a priori* knowledge of r_d and r_b , which are combined in the parameter R . This means that R is kept constant and identical in the individual analysis of all related recovery curves. However, in a simultaneous analysis of related recovery curves arising from different ROI sizes, the common value of the parameter R can be determined by linking R over the related curves, i.e. there is a single parameter R in the optimization routine. Best results are obtained when R is linked across as many ROI sizes as possible.

This approach does yield D values close to the expected values [Table 1]. Compared to the approach that applies a fixed R , the recovered values of D are identical within the experimental error and their uncertainties nearly double but remain of the same order of magnitude. In other words, even without making use of *a priori* information on R , the correct D can be obtained when a range of ROI sizes is considered.

Since r_d is known, the r_b values can be calculated from the fitted R values. For all data sets analyzed, a physically relevant r_b value approaching the calibrated value is obtained [Table 2].

4.3 Importance of R per ROI size

The correction factor R is more important with decreasing ROI radius (*vide supra*). To experimentally confirm this statement, the accuracy by which R can be determined per ROI size is investigated. R should be obtained with higher accuracy from small ROI radius as compared to their larger counterparts. This was investigated by linking D over all recovery

curves and by linking R within experiments of the same ROI size. The recovered values of R displayed a variability amongst the different ROI sizes without significant differences [Figure 5]. It was obvious that the uncertainty on R increased with increasing ROI radius, thereby corroborating the theory.

4.4 Effect of omitting the large ROIs in the global analysis

Analysis including recovery curves from all ROI sizes led to a reliable result. Experimental conditions, however, do not always allow for large ROIs in case of small samples. Therefore, it was necessary to investigate the performance of the new GDM with freely adjustable R when only these smaller sizes are considered. A global analysis was repeatedly performed starting with only the two smallest ROI sizes and progressively including larger ROI sizes. Figure 6 displays the results of two representative data sets obtained with FD2000 in 40% (a) and 56% (b) (w/w) sucrose using *setup A* and analyzed by the presented method.

Applying this strategy using UDM, it comes as no surprise that the obtained values for D using only small ROIs severely deviate from the expected value [Figure 6]. When also larger ROIs were included, D increased but never reached the expected value. With the GDM and a fixed value of R , on the other hand, the correct value of D could be found already with only 2 to 3 small ROI sizes at the expense of a larger uncertainty on the fitted parameters [Figure 6]. This conclusion was similar amongst all data sets, indicating that small ROIs are sufficient to accurately determine D .

A similar analysis was performed using GDM but with a linked R across the data as a free fitting parameter. This approach yielded better D values than the UDM. In comparison with GDM with a fixed R , somewhat more variation in the recovered D values could be observed. Nevertheless, even without any prior knowledge of R , the GDM clearly outperforms the UDM.

5 Discussion

The measured diffusion coefficient D of a system with free diffusion should not vary upon changing the radius of the bleach ROI. As shown by Braeckmans *et al.*² and also illustrated by the results in this study, this is not the case for small or intermediate sized ROIs when analyzed by the UDM. Since these small or intermediate sized ROIs are sometimes required for diffusion measurements in small samples,⁹⁻¹¹ a new model was developed. The resulting generalized disk model presented in this paper renders D essentially insensitive to the ROI radius, as is clearly demonstrated by the single curve analysis of recovery curves for each ROI size. This enables the use of these small and intermediate ROI sizes to obtain accurate D values on any CLSM.

Absence of recovery during photobleaching enables the full description of the resulting concentration profile immediately after bleaching by solely the radius of the ROI and the effective photobleaching resolution r_b . Together with the finite total detection resolution r_d , *a priori* knowledge of r_b is essential for the GDM using single curve analyses. While r_d can be determined by straightforward recording of sub-resolution beads, the calibration measurement required for determination of r_b is more laborious and difficult to obtain accurately since it depends on the type of fluorophore, photon flux and local chemical environment. Using a reference solution of known D , r_b can be estimated in a separate calibration measurement by means of lineFRAP.¹² However, this calibration does not include the rise time of the AOM and the calibrated r_b might still be underestimated. Therefore, it is suggested to use the GDM in combination with the global fitting procedure since this allows to extract the R value from FRAP experiments performed at various ROI sizes. Linking R across only a few small ROI sizes is sufficient to obtain a reliable result. For the experiments considered in this work about

three ROI sizes were sufficient. Besides the estimation of R from the optimization, global analysis offers a second important advantage. The lower signal-to-noise ratio associated with smaller ROI sizes, together with a limited photobleaching depth required for the GDM, might decrease the accuracy of single curve fit analyses. Global analysis of all recovery curves, in contrast, is still capable of returning relevant parameters, despite of a low signal-to-noise ratio.¹⁹

Conclusion

A new generalized disk model is introduced in a closed-form expression for analysis of FRAP recovery curves obtained at any ROI size. Using a simultaneous analysis of recovery curves obtained with a variety of ROI sizes, even in the absence of large ROIs, the diffusion coefficient D can reliably be obtained without prior calibration of the resolution parameters.

~~This enables the use~~ As a result, the range of ROI sizes that can be applied in FRAP experiments performed using a CLSM is successfully extended to small or intermediate circular ROI sizes which are only a few times the total imaging PSF of CLSMs in FRAP experiments using small or intermediate circular ROI sizes.

~~The exclusion of size dependent artifacts is essential when studying systems by varying the radius of the bleach ROI.~~ The possibility to use a large range of ROI sizes not only offers the possibility of more flexible diffusion measurements, but is also expected to be valuable for more complicated measurements such as detecting~~This approach is e latter is applied to probe the process of interest at different length scales to detect~~ spatial heterogeneities, as applied to the study of in the membrane organization or receptor distribution in the plasma membrane of live cells.^{9,10,20} ~~Another application that will benefit from this new approach includes and to the study investigate of the connectivity between different domains in heterogeneous heterogeneities in (bio)materials, through diffusion of a suitable fluorescent probe.~~²¹

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References

1. D. Axelrod, D. E. Koppel, J. Schlessinger, E. Elson and W. W. Webb, "Mobility measurement by analysis of fluorescence photobleaching recovery kinetics," *Biophys. J.* **16**(9), 1055-1069 (1976).
2. K. Braeckmans, L. Peeters, N. N. Sanders, S. C. De Smedt and J. Demeester, "Three-dimensional fluorescence recovery after photobleaching with the confocal scanning laser microscope," *Biophys. J.* **85**(4), 2240-2252 (2003).
3. A. Lopez, L. Dupou, A. Altibelli, J. Trotard and J. F. Tocanne, "Fluorescence recovery after photobleaching (FRAP) experiments under conditions of uniform disk illumination. Critical comparison of analytical solutions, and a new mathematical method for calculation of diffusion coefficient D," *Biophys. J.* **53**(6), 963-970 (1988).
4. D. M. Soumpasis, "Theoretical analysis of fluorescence photobleaching recovery experiments," *Biophys. J.* **41**(1), 95-97 (1983).
5. U. Kubitscheck, P. Wedekind and R. Peters, "Lateral diffusion measurement at high spatial resolution by scanning microphotolysis in a confocal microscope," *Biophys. J.* **67**(3), 948-956 (1994).
6. U. Kubitscheck, P. Wedekind and R. Peters, "Three-dimensional diffusion measurements by scanning microphotolysis," *J. Microsc.* **192**(2), 126-138 (1998).
7. R. Peters and U. Kubitscheck, "Scanning microphotolysis: three-dimensional diffusion measurement and optical single-transporter recording," *Methods* **18**(4), 508-517 (1999).

8. P. Wedekind, U. Kubitscheck and R. Peters, "Scanning microphotolysis: a new photobleaching technique based on fast intensity modulation of a scanned laser beam and confocal imaging," *J. Microsc.* **176**(Pt 1), 23-33 (1994).
9. A. M. Baker, A. Sauliere, G. Gaibelet, B. Lagane, S. Mazeres, M. Fourage, F. Bachelierie, L. Salome, A. Lopez and F. Dumas, "CD4 interacts constitutively with multiple CCR5 at the plasma membrane of living cells. A fluorescence recovery after photobleaching at variable radii approach," *J. Biol. Chem.* **282**(48), 35163-35168 (2007).
10. L. Salome, J. L. Cazeils, A. Lopez and J. F. Toccane, "Characterization of membrane domains by FRAP experiments at variable observation areas," *Eur. Biophys. J.* **27**(4), 391-402 (1998).
11. Y. I. Henis, B. Rotblat and Y. Kloog, "FRAP beam-size analysis to measure palmitoylation-dependent membrane association dynamics and microdomain partitioning of Ras proteins," *Methods* **40**(2), 183-190 (2006).
12. K. Braeckmans, K. Remaut, R. E. Vandenbroucke, B. Lucas, S. C. De Smedt and J. Demeester, "Line FRAP with the confocal laser scanning microscope for diffusion measurements in small regions of 3-D samples," *Biophys. J.* **92**(6), 2172-2183 (2007).
13. D. Mazza, F. Cella, G. Vicidomini, S. Krol and A. Diaspro, "Role of three-dimensional bleach distribution in confocal and two-photon fluorescence recovery after photobleaching experiments," *Appl. Opt.* **46**(30), 7401-7411 (2007).
14. K. Braeckmans, B. G. Stubbe, K. Remaut, J. Demeester and S. C. De Smedt, "Anomalous photobleaching in fluorescence recovery after photobleaching measurements due to excitation saturation--a case study for fluorescein," *J. Biomed. Opt.* **11**(4), 044013 (2006).
15. J. E. N. Jonkman and E. H. K. Stelzer, "Resolution and Contrast in Confocal and Two-Photon Microscopy," Chap. 5 in *Confocal and Two-Photon Microscopy: Foundations, Applications and Advances*, A. Diaspro, Ed., pp. 101-123, Wiley-Liss, Inc., New York (2002).
16. T. K. Meyvis, S. C. De Smedt, P. Van Oostveldt and J. Demeester, "Fluorescence recovery after photobleaching: a versatile tool for mobility and interaction measurements in pharmaceutical research," *Pharm. Res.* **16**(8), 1153-1162 (1999).
17. R. B. Dalal, M. A. Digman, A. F. Horwitz, V. Vetri and E. Gratton, "Determination of particle number and brightness using a laser scanning confocal microscope operating in the analog mode," *Microsc. Res. Tech.* **71**(1), 69-81 (2008).
18. A. V. Wolf, M. G. Brown and P. G. Prentiss, "Concentrative properties of aqueous solutions: conversion tables," in *CRC Handbook of Chemistry and Physics*, R. C. Weast, Ed., 60th edition, pp. 227-276, CRC Press, Inc., Boca Raton, Florida (1979-1980).

19. J. M. Beechem, E. Gratton, M. Ameloot, J. R. Knutson and L. Brand, "The global analysis of fluorescence intensity and anisotropy decay data: second generation theory and programs," in *Topics in Fluorescence Spectroscopy: Volume2: Principles*, J. R. Lakowicz, Ed., pp. 241-305, Plenum Press, New York (1991).
20. R. Saxena and A. Chattopadhyay, "Membrane organization and dynamics of the serotonin(1A) receptor in live cells," *J Neurochem* **116**(5), 726-733 (2011).
21. N. Loren, M. Nyden and A. M. Hermansson, "Determination of local diffusion properties in heterogeneous biomaterials," *Adv Colloid Interface Sci* **150**(1), 5-15 (2009).

Tables

Table 1: Global analysis of FITC-dextran experiments for UDM and GDM at room temperature. Uncertainties are reported as standard deviations.

Setup	MW [kDa]	Sucrose / HEPES [% (w/w)]	Diffusion coefficient [$\mu\text{m}^2/\text{s}$]			
			Expected	UDM	GDM	
					R fixed*	R freely adjustable
A	2000	40	1.07	1.03 ± 0.01	1.08 ± 0.01	1.10 ± 0.02
A	2000	56	0.205	0.179 ± 0.003	0.216 ± 0.003	0.207 ± 0.005
B	500	40	2.22	2.11 ± 0.03	2.39 ± 0.03	2.26 ± 0.05
B	500	56	0.427	0.371 ± 0.008	0.423 ± 0.008	0.42 ± 0.02

* *Setup A*: $R = 1.8 \mu\text{m}^2$

Setup B: $R = 1.1 \mu\text{m}^2$

Table 2: Comparison of the calibrated bleach resolution r_b with the recovered value from the global analyses using the obtained R and the calibrated value of r_d .

Setup	MW FD [kDa]	sucrose [% (w/w)]	calibrated values			recovered values	
			r_d [μm]	r_b [μm]	R [μm^2]	R [μm^2]	r_b [μm]
A	2000	40	1.0	2.5	1.8	2.4 ± 0.3	2.9 ± 0.4
A	2000	56	1.0	2.5	1.8	1.4 ± 0.2	2.1 ± 0.3
B	500	40	0.9	1.9	1.1	0.8 ± 0.2	1.6 ± 0.4
B	500	56	0.9	1.9	1.1	1.1 ± 0.3	1.9 ± 0.5

Figures

Figure 1: (a) Three-dimensional illustration of $B(r)$ representing the ideal circular bleach geometry with sharp edges. In reality, this geometry is bleached by a scanning, focused laser beam with an effective bleaching intensity distribution $I_b(r)$ shown in (b). The white circle indicates the e^{-2} intensity level at which its radius equals r_b . The resulting bleach intensity distribution $K(r)$ is consequently the convolution product of $B(r)$ and $I_b(r)$. One half of $K(r)$ is shown for a radius w of 5 times r_b (c) and 1 time r_b (d). The gray shaded cylinder with dashed lines indicates the corresponding $B(r)$.

Figure 2: Three typical phases in a time series of a FRAP experiment. (a) The fluorescence intensity is first recorded before bleaching. (b) The selected region of interest (indicated in white) is bleached. Only *setup A* registers the fluorescence intensities during this phase. (c) and (d) are respectively the first (directly after bleaching) and the last image of a series of post-bleach images monitoring the recovery of the bleached region.

Figure 3: The experimental relationship between the average pixel intensity (in digital levels) and the associated variance is shown in (a) for a given set of experimental parameters of *setup B* and a solution of FD500 in 40% (w/w) sucrose. Each average-variance pair is obtained from a homogeneous region from FRAP time series. In (b) and (c), two representative experimental, normalized recovery curves are shown. These curves were obtained using a ROI with a radius of 1.9 μm (b) and 15.2 μm (c). Error bars indicate the estimated standard deviations. Time is expressed in units of τ_D to display both curves on an identical scale.

Figure 4: Comparison of the UDM (empty squares) and the GDM (filled bullets). Representative data sets are shown for *setup A* with FD2000 (a) and for *setup B* with FD500 (b) dissolved in respectively 56% and 40% (w/w) sucrose. Results are grouped per ROI radius using weighted averages of the obtained diffusion coefficient D . Number of recovery curves per ROI radius varied between 5 and 10. The UDM returned a diffusion coefficient that decreased with decreasing ROI radius, reaching a maximum underestimation at the smallest ROI size. Error bars are shown as standard deviations.

Figure 5: Resolution factor R as a function of ROI radius, recorded with *setup A* and FD2000 dissolved in 56% (w/w) sucrose in HEPES buffer. Error bars represent standard deviations.

Figure 6: Related data sets were repeatedly analyzed using global analysis while progressively including larger ROIs. D was linked across all included recovery curves. Representative data sets are shown obtained with FD2000 in 40% (a) and 56% (b) w/w sucrose/HEPES buffer using *Setup A*. Error bars representing the standard deviations are shown if larger than the symbol size.