

# On the fundamental solution and integral formulae of a higher spin operator in several vector variables

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**Abstract.** The fundamental solution is constructed of a particular higher spin operator in  $\mathbb{R}^m$ , acting on functions taking values in an irreducible representation space for  $\text{Spin}(m)$  with highest weight  $(k + \frac{1}{2}, l + \frac{1}{2}, \frac{1}{2}, \dots, \frac{1}{2})$ , with  $k, l \in \mathbb{N}$  and  $k \geq l$ . Some basic integral formulae are stated.

**Keywords:** Clifford analysis, higher spin operator, fundamental solution.

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## INTRODUCTION AND DEFINITIONS

Let  $(e_1, \dots, e_m)$  be an orthonormal basis for the Euclidean space  $\mathbb{R}^m$  and denote by  $\mathbb{C}_m$  the complex  $2^m$ -dimensional Clifford algebra over  $\mathbb{R}^m$ , generated by the relations  $e_i e_j + e_j e_i = -2\delta_{i,j}$ ,  $i, j = 1, \dots, m$ . By identifying  $(x_1, \dots, x_m)$  with the real Clifford vector  $x = \sum_{j=1}^m e_j x_j$ , the space  $\mathbb{R}^m$  is embedded in  $\mathbb{C}_m$ . The multiplication of two vectors  $x$  and  $y$  is given by  $xy = -\langle x, y \rangle + x \wedge y$ , with

$$\langle x, y \rangle = \sum_{j=1}^m x_j y_j \quad \text{and} \quad x \wedge y = \sum_{1 \leq i < j \leq m} e_i e_j (x_i y_j - x_j y_i)$$

being the scalar-valued Euclidean inner product and the bivector-valued wedge product respectively.

The Pin group  $\text{Pin}(m)$  is the group consisting of products of unit vectors in  $\mathbb{R}^m$ . The Spin group  $\text{Spin}(m)$  is the subgroup of  $\text{Pin}(m)$  consisting of products of an even number of unit vectors in  $\mathbb{R}^m$ . For convenience we work in odd dimension  $m = 2n + 1$ , in which case there is a unique spinor space  $\mathbb{S}$ . This space is an irreducible  $\text{Spin}(m)$ -module and can be thought of as a minimal left ideal in  $\mathbb{C}_m$ .

The Dirac operator on  $\mathbb{R}^m$  is given by  $\partial_x = \sum_{j=1}^m e_j \partial_{x_j}$ . It is an elliptic  $\text{Spin}(m)$ -invariant first-order differential operator acting on  $\mathbb{S}$ -valued functions  $f(x)$  on  $\mathbb{R}^m$ . A spinor-valued function  $f$  is monogenic in an open region  $\Omega \subset \mathbb{R}^m$  if and only if it satisfies  $\partial_x f = 0$  in  $\Omega$ . For a detailed account of the theory of monogenic functions, so-called Euclidean Clifford analysis, we refer to e.g. [2, 10, 14]. More abstractly, on manifolds, the Dirac operator can be regarded as the simplest case of the unique (up to a multiple) conformally invariant elliptic first-order operator acting between appropriate sections of spin bundles, see e.g. [6, 11, 17]. In Clifford analysis these operators are studied from a function theoretical point of view, considering functions on  $\mathbb{R}^m$  instead of sections. Another example of such operators are the generalized Rarita-Schwinger operators  $\mathcal{R}_k$ , acting on functions taking values in irreducible  $\text{Spin}(m)$ -representation spaces with highest weight  $(k + \frac{1}{2}, \frac{1}{2}, \dots, \frac{1}{2})$  with  $k \in \mathbb{N}$ , see e.g. [7, 8]. This notion has been generalized in [3, 4], where the operator  $\mathcal{Q}_{k,l}$  is studied, acting on functions with values in the irreducible  $\text{Spin}(m)$ -module with highest weight  $(k + \frac{1}{2}, l + \frac{1}{2}, \frac{1}{2}, \dots, \frac{1}{2})$  with  $k, l \in \mathbb{N}$  and  $k \geq l$ .

Both representations can be described by vector spaces of polynomials within the framework of Clifford analysis, see e.g. [9]. For example, the space  $\mathcal{H}_k$  of  $\mathbb{C}$ -valued harmonic homogeneous polynomials of degree  $k \in \mathbb{N}$  corresponds to the irreducible  $\text{Spin}(m)$ -module with highest weight  $(k, 0, \dots, 0)$ , and the space  $\mathcal{M}_k$  of  $\mathbb{S}$ -valued monogenic homogeneous polynomials of degree  $k$  forms an irreducible representation of  $\text{Spin}(m)$  with highest weight  $(k + \frac{1}{2}, \frac{1}{2}, \dots, \frac{1}{2})$ . Both notions are generalized in the following definitions, where  $N \in \mathbb{N}$  and  $\partial_i$  is short for the Dirac operator  $\partial_{u_i}$ .

**Definition 1.** A function  $f : \mathbb{R}^{Nm} \rightarrow \mathbb{C}$ ,  $(u_1, \dots, u_N) \mapsto f(u_1, \dots, u_N)$  is simplicial harmonic if the following conditions are satisfied:  $\langle \partial_i, \partial_j \rangle f = 0$  ( $i, j = 1, \dots, N$ ) and  $\langle u_i, \partial_j \rangle f = 0$  ( $1 \leq i < j \leq N$ ). The vector space of  $\mathbb{C}$ -valued simplicial harmonic polynomials,  $\lambda_i$ -homogeneous in the variable  $u_i$ , will be denoted by  $\mathcal{H}_{\lambda_1, \dots, \lambda_N}$  (with  $\lambda_1 \geq \dots \geq \lambda_N \geq 0$ ).

**Definition 2.** A function  $f: \mathbb{R}^{Nm} \rightarrow \mathbb{S}$ ,  $(u_1, \dots, u_N) \mapsto f(u_1, \dots, u_N)$  is simplicial monogenic if the following conditions are satisfied:  $\partial_i f = 0$  ( $i = 1, \dots, N$ ) and  $\langle u_i, \partial_j \rangle f = 0$  ( $1 \leq i < j \leq N$ ). The vector space of  $\mathbb{S}$ -valued simplicial monogenic polynomials,  $\lambda_i$ -homogeneous in the variable  $u_i$ , will be denoted by  $\mathcal{S}_{\lambda_1, \dots, \lambda_N}$  (with  $\lambda_1 \geq \dots \geq \lambda_N \geq 0$ ).

Take  $N = 2$  and  $(u_1, u_2) = (u, v)$  in both definitions. In [3], the operator  $\mathcal{Q}_{k,l}$  mentioned before was constructed by means of the following decomposition:

**Proposition 1.** For any pair of integers  $k \geq l \geq 0$  with  $k > 0$ , one has

$$\mathcal{H}_{k,l} \otimes \mathbb{S} = \mathcal{S}_{k,l} \oplus (1 - \delta_{l,0}) \nu_{k,l} \mathcal{S}_{k,l-1} \oplus (1 - \delta_{k,l}) \mu_{k,l} \mathcal{S}_{k-1,l} \oplus (1 - \delta_{l,0}) \kappa_{k,l} \mathcal{S}_{k-1,l-1},$$

with  $\nu_{k,l}$ ,  $\mu_{k,l}$  and  $\kappa_{k,l}$  the respective embedding maps of the spaces  $\mathcal{S}_{k,l-1}$ ,  $\mathcal{S}_{k-1,l}$  and  $\mathcal{S}_{k-1,l-1}$  in  $\mathcal{H}_{k,l} \otimes \mathbb{S}$ .

If we denote by  $\pi_{k,l}$  the projection operator  $\pi_{k,l}: \mathcal{H}_{k,l} \otimes \mathbb{S} \rightarrow \mathcal{S}_{k,l}$ , we have

**Definition 3.** For all integers  $k \geq l \geq 0$  with  $k > 0$ , there are (up to a multiplicative constant) unique invariant first-order differential operators  $\mathcal{Q}_{k,l}$  defined by

$$\mathcal{Q}_{k,l}: \mathcal{C}^\infty(\mathbb{R}^m, \mathcal{S}_{k,l}) \rightarrow \mathcal{C}^\infty(\mathbb{R}^m, \mathcal{S}_{k,l}), f(x; u, v) \mapsto \pi_{k,l}(\partial_x f)(x; u, v).$$

Explicitly,

$$\mathcal{Q}_{k,l} f = \left( \mathbf{1} + \frac{u \partial_u}{m + 2k - 2} \right) \left( \mathbf{1} + \frac{v \partial_v}{m + 2l - 4} \right) \partial_x f.$$

**Remark 1.** In case  $l = 0$ , the operator  $\mathcal{Q}_{k,0}$  reduces to the Rarita-Schwinger operators  $\mathcal{R}_k$ , as defined in [7].

This contribution fits into the encompassing framework in which we are carrying out a thorough study of the higher spin Dirac operators  $\mathcal{Q}_{k,l}$ . It aims at constructing their fundamental solution which allows for establishing basic integral formulae such as Stokes' theorem, the Cauchy-Pompeiu representation theorem and the Cauchy integral formula for their null solutions, which are also stated. In their turn these integral formulae will prove to be essential to developing a function theory associated to these higher spin operators.

## CONFORMAL INVARIANCE IN CLIFFORD ANALYSIS

The higher spin operators  $\mathcal{Q}_{k,l}$  are not just  $\text{Spin}(m)$ -invariant, they are also conformally invariant, see e.g. [11, 16]. The conformal group can be described by the Vahlen group  $V(m)$  of  $2 \times 2$  matrices with entries in  $\mathbb{C}_m$ , defined as follows:

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} \in V(m) \Leftrightarrow \begin{cases} a, b, c, d \in \Gamma(m) \cup \{0\}, \\ ab^*, cd^*, d^*b, c^*a \in \mathbb{R}^m, \\ ad^* - bc^* = \pm 1, \end{cases}$$

with  $\Gamma(m)$  the Clifford group. See e.g. [1, 15] for a detailed account of conformal invariance in Clifford analysis. In order to construct the fundamental solution of  $\mathcal{Q}_{k,l}$ , we use a general theorem of [16] about conformally invariant differential operators, adapted to our case of operators acting on  $\mathcal{C}^\infty(\mathbb{R}^m, \mathcal{S}_{k,l})$ . Let  $L$  be the standard representation of  $\text{Pin}(m)$  acting on the values  $\mathcal{S}_{k,l}$ .

**Theorem 1.** If  $f$  is a null solution of  $\mathcal{Q}_{k,l}$ , so is the transformed function  $(g \cdot f)$ , defined by

$$(g \cdot f)(x) = |cx + d|^{-m+1} L \left( \frac{(cx + d)^*}{|cx + d|} \right) f((ax + b)(cx + d)^{-1}; u, v)$$

with  $g^{-1} = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$  an element of the conformal group  $V(m)$ .

## CONSTRUCTION OF THE FUNDAMENTAL SOLUTION

In this section it will become clear that the construction of the fundamental solution of the higher spin operators  $\mathcal{Q}_{k,l}$  needs results from a variety of disciplines such as Clifford analysis, distribution theory and group representation theory, and moreover involves a lot of technicalities. We present a concise but rather complete view on this construction. Denote by  $A_m$  the surface area of the unit sphere in  $\mathbb{R}^m$ . Recall that the Cauchy kernel  $E(x) = -(1/A_m)(\bar{x}/|x|^m)$  satisfies  $\partial_x E(x) = \delta(x)$  in distributional sense. Choose an arbitrary polynomial  $P$  in  $\mathcal{S}_{k,l}$  and let  $C_{k,l}$  be a constant, which will be specified later. Inspired by Theorem 1, we define:

$$E_{k,l}(x; u, v) = C_{k,l}|x|^{-m+1}L(x/|x|)P(u, v) = C_{k,l}|x|^{-(m+2k+2l)}P(xux, xv x) \quad (1)$$

We will prove that  $E_{k,l}(x; u, v)$  satisfies  $\mathcal{Q}_{k,l}E_{k,l}(x; u, v) = \delta(x)P(u, v)$  in distributional sense. Instead of taking an arbitrary  $P \in \mathcal{S}_{k,l}$ , it is convenient to work with the following generator of  $\mathcal{S}_{k,l}$ , see [9] for more information:

$$P_{k,l}(u, v) := \langle u, \tau_1 \rangle^{k-l} \left[ \det \begin{pmatrix} \langle u, \tau_1 \rangle & \langle u, \tau_2 \rangle \\ \langle v, \tau_1 \rangle & \langle v, \tau_2 \rangle \end{pmatrix} \right]^l I_{12} \in \mathcal{S}_{k,l} \quad (2)$$

with  $\tau_1$  and  $\tau_2$  null vectors in  $\mathbb{C}^m$ , i.e.  $\tau_1^2 = \tau_2^2 = 0$ , the idempotent  $I_{12} = \tau_1 \bar{\tau}_1 \tau_2 \bar{\tau}_2$ . After calculations involving Definition 3 and some well-known identities in Clifford analysis, we find for  $\alpha \in \mathbb{C}$ :

$$\mathcal{Q}_{k,l}(|x|^{\alpha-2(k+l)} x P_{k,l}(xux, xv x)) = -(\alpha + m)|x|^{\alpha-2(k+l)} \pi_{k,l}(P_{k,l}(xux, xv x)). \quad (3)$$

If  $\alpha = -m$ , then  $|x|^{-m-2(k+l)} x P_{k,l}(xux, xv x)$  belongs to the kernel of  $\mathcal{Q}_{k,l}$  and shows a pointwise singularity of homogeneity degree  $-m+1$  at the origin. In order to take the limit  $\alpha \rightarrow -m$ , it is necessary to recall some basic facts about distributions related to Riesz potentials in  $\mathbb{R}^m$ , see e.g. [12, 13]. We show in [5] that the relation (3) remains valid in distributional sense if  $\Re(\alpha) > -m-1$ . Furthermore,

$$\mathcal{Q}_{k,l}(|x|^{-m-2(k+l)} x P_{k,l}(xux, xv x)) = -\frac{\pi^{\frac{m}{2}} 2^{-2(k+l)+1}}{\Gamma(\frac{m}{2} + k + l)(k+l)!} \pi_{k,l}(\Delta_x^{k+l} P_{k,l}(xux, xv x)) \delta(x). \quad (4)$$

Denote by  $\mathcal{P}_h$  the space of  $h$ -homogeneous polynomials. To calculate the right-hand side of (4), define the operator  $\mathcal{R}_x : \mathcal{S}_{k,l} \rightarrow \mathcal{P}_{2(k+l)} \otimes \mathcal{H}_{k,l}$ ,  $P_{k,l}(u, v) \mapsto P_{k,l}(xux, xv x)$ . It can be proven that  $\Delta_x^{k+l} \mathcal{R}_x$  is a  $\text{Spin}(m)$ -invariant map acting on  $\mathcal{S}_{k,l}$ . Hence it follows from Schur's lemma that, for some constant  $C$ , holds  $\Delta_x^{k+l} P_{k,l}(xux, xv x) = C P_{k,l}(u, v)$ , which is valid for every  $P_{k,l}(u, v)$  in  $\mathcal{S}_{k,l}$  since  $\mathcal{S}_{k,l}$  is irreducible with respect to  $\text{Spin}(m)$ . To determine this constant  $C$ , we complexify  $u, v, \tau_1$  and  $\tau_2$ . If  $u = e_1 - ie_2, v = e_3 - ie_4, \tau_1 = e_1 + ie_2$  and  $\tau_2 = e_3 + ie_4$ , then we can show that

$$C = 2^{2(k+l)}(k+l)! \frac{\Gamma(\frac{m}{2} + k + l) \Gamma(\frac{m-2}{2} + k) \Gamma(\frac{m-4}{2} + l)}{\Gamma(\frac{m}{2} + k) \Gamma(\frac{m-2}{2} + l) \Gamma(\frac{m-4}{2})}.$$

The right-hand side of (4) can thus be simplified to

$$\mathcal{Q}_{k,l}(|x|^{-m-2(k+l)} x P_{k,l}(xux, xv x)) = -A_m \frac{(m-2)(m-4)}{(m+2k-2)(m+2l-4)} P_{k,l}(u, v) \delta(x) =: (C_{k,l})^{-1} P_{k,l}(u, v) \delta(x).$$

Hence,

$$E_{k,l}(x; u, v) := C_{k,l}|x|^{-m-2(k+l)} x P_{k,l}(xux, xv x) \in \mathcal{C}^\infty(\mathbb{R}^m \setminus \{0\}, \mathcal{S}_{k,l})$$

satisfies

$$\mathcal{Q}_{k,l}E_{k,l}(x; u, v) = \delta(x)P_{k,l}(u, v). \quad (5)$$

If we define

$$e_{k,l}(x) := C_{k,l}|x|^{-m-2(k+l)}L(x/|x|) \in \mathcal{C}^\infty(\mathbb{R}^m \setminus \{0\}, \text{End}(\mathcal{S}_{k,l})),$$

we have  $E_{k,l}(x; u, v) = e_{k,l}(x)P_{k,l}(u, v)$ . It follows from (5) that  $\mathcal{Q}_{k,l}e_{k,l}(x)P_{k,l}(u, v) = \delta(x)P_{k,l}(u, v)$ . This relation depends on  $P_{k,l} \in \mathcal{S}_{k,l}$ . For this reason, we introduce the reproducing kernel  $K_{k,l}$  satisfying

$$(K_{k,l}(u, v; u', v'), P_{k,l}(u, v))_{(u,v)} = P_{k,l}(u', v'),$$

with  $(\cdot, \cdot)_{(u,v)}$  the Fischer inner product on  $\mathcal{S}_{k,l}$  defined by  $(P, Q)_{(u,v)} = [\bar{P}(\frac{\partial}{\partial u}, \frac{\partial}{\partial v})Q]_0$ . Now, if we put

$$E_{k,l}(x; u, v; u', v') := e_{k,l}(x)K_{k,l}(u, v; u', v') \in \mathcal{C}^\infty(\mathbb{R}^m \setminus \{0\}, \text{End}(\mathcal{S}_{k,l})),$$

we finally find

$$\mathcal{Q}_{k,l}E_{k,l}(x; u, v; u', v') = \delta(x)K_{k,l}(u, v; u', v'). \quad (6)$$

## BASIC INTEGRAL FORMULAE

The basic integral formulae related to the operators  $\mathcal{Q}_{k,l}$  can be deduced from Stokes' theorem for the Dirac operator, see e.g. [10]. Put  $dx = dx_1 \wedge \dots \wedge dx_m$  and  $d\sigma_x = \sum_{j=1}^m (-1)^{j-1} e_j d\hat{x}_j$  with  $d\hat{x}_j = dx_1 \wedge \dots \wedge dx_{j-1} \wedge dx_{j+1} \wedge \dots \wedge dx_m$ . In [5] we proved:

**Theorem 2.** *Let  $\Omega' \subset \mathbb{R}^m$  and  $\bar{\Omega} \subset \Omega'$ . Then for  $f, g \in \mathcal{C}^1(\Omega', \mathcal{S}_{k,l})$  one has:*

(i) *Stokes' Theorem.*

$$\int_{\Omega} \left[ -(\mathcal{Q}_{k,l}g(x), f(x))_{(u,v)} + (g(x), \mathcal{Q}_{k,l}f(x))_{(u,v)} \right] dx = \int_{\partial\Omega} (g(x), \pi_{k,l}(d\sigma_x)f(x))_{(u,v)}.$$

(ii) *Cauchy-Pompeiu Theorem.* *Let  $y \in \Omega$ . Then*

$$f(y) + \int_{\Omega} e_{k,l}(x-y) \mathcal{Q}_{k,l}f(x) dx = \int_{\partial\Omega} e_{k,l}(x-y) \pi_{k,l}(d\sigma_x)f(x).$$

(iii) *Cauchy integral formula.* *If  $\mathcal{Q}_{k,l}f = 0$ , then for  $y \in \Omega$  one has:*

$$f(y) = \int_{\partial\Omega} e_{k,l}(x-y) \pi_{k,l}(d\sigma_x)f(x).$$

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